# Quantification of Boundary Layer Mixing over the Southern Ocean Using In-Situ and Remotely-Sensed Measurements

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#### Abstract

We demonstrate that the relationship between the abundance of particulate surface area observed at sea-level and measurements of backscattered light by a ceilometer can be used to classify the mixing state of the atmospheric layer beneath the lowest observed cloud, where the relationship is defined by the Spearman Rank correlation. The accuracy of this correlation-based method was compared to two methods of detecting boundary layer decoupling based on radiosonde measurements. An optimized version of the new methodology correctly determined the mixing state of the below-cloud layer for  $76 \pm 4\%$  of the radiosondes available for comparison. Further, it was more accurate than an alternative ground-based metric used to determine the below-cloud mixing state. For the majority of the time series in which the correlation analysis could be applied, the below-cloud boundary layer was well-mixed (54%), or else fog was present (27%), which indicated that aerosol particles observed at sea-level often have a direct pathway into low-cloud (81%). In the remaining analysis period, the near-surface atmospheric layer was stable and the atmospheric layer near the ocean surface was decoupled from the overlying cloud (19%). Forecasts from the Antarctic Mesoscale Prediction System also support our findings, showing that conditions that mix aerosol particles from the ocean surface to the low-cloud occur 84% of the time over the open Southern Ocean. As a result, aerosol particles measured near sea-level are often tightly coupled to low-cloud formation over the Southern Ocean, highlighting the utility of shipborne aerosol observations in the region.

## Classification of the Below-Cloud Mixing State Over the Southern Ocean Using In-Situ and Remotely-Sensed Measurements

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## Key Points:

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8	•	A new, correlation-based method determines the mixing state of the below-
9		cloud layer with an accuracy of $76\%$ .
10	•	The below-cloud layer is only ever poorly-mixed when winds are below 8 m
11		$s^{-1}$ and the near-surface atmosphere is neutrally stable.
12	•	Sea spray particles are available to low-level cloud over the Southern Ocean
13		more than $80\%$ of the time throughout austral summer.

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#### 14 Abstract

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#### <sup>36</sup> Plain Language Summary

Particles suspended in the atmosphere (aerosol) act as seeds for cloud droplet 37 formation. The abundance of such particles directly influences the opacity of clouds, 38 while their physical and chemical characteristics govern if and when those cloud 39 droplets freeze. As a result, both the amount of solar radiation a cloud can reflect 40 and the temperature of waters below are sensitive to the quantity and type of par-41 ticles available to the cloud. We present a new methodology for understanding the 42 conditions in which low-level clouds have direct access to the large and diverse reser-43 voir of particles in the surface layer. We find that meteorological conditions which 44 transfer particles from sea-level to low-level cloud are satisfied up to 81% of the 45 time over the Southern Ocean. This suggests that the particles we observe near the 46 surface almost always play a significant role in the formation of low-level cloud. 47

#### 48 1 Introduction

The balance of incoming and outgoing shortwave radiation is mediated by the 49 presence of clouds, and to a lesser extent, aerosol particles. While aerosol particles 50 are significantly smaller than clouds, they are the seeds of cloud droplet forma-51 tion (Pruppacher et al., 1998). The chemical and physical nature of these particles 52 strongly determines both the ultimate phase of cloud droplets (e.g. liquid or ice) 53 and the resulting size distribution of cloud droplets (Twomey, 1977). It therefore fol-54 lows that errors in how these particles are represented within global climate models 55 can cause significant climatological biases in the radiative balance. In particular, 56 it has been observed that uncertainties in predicting cloud phase leads to sub-57 stantial biases in the radiation balance within the cold sector of Southern Ocean 58 cyclones (Bodas-Salcedo et al., 2014). While the abundant cyclones of the South-59 ern Ocean occur solely as a function of favourable synoptic conditions (Irving et 60 al., 2010), a global climate model's predictions of cloud phase in the cold sector of 61 Southern Ocean cyclones (Vergara-Temprado et al., 2018), and in the wider South-62 ern Ocean (Schuddeboom et al., 2019), is extremely sensitive to the properties of 63

particles in the underlying boundary layer. In addition, recent observations directly
show that the large biases between modeled and observed outgoing shortwave radiation (Trenberth & Fasullo, 2010) are related to errors in how low-level cloud are
represented within global climate models (Kuma et al., 2020). Improving our understanding of the conditions in which particles can reach, and form, low-level cloud is
therefore crucial to understanding the radiation balance over the Southern Ocean.

As wind speeds have increased over the Southern Ocean (Young et al., 2011; 70 Hande et al., 2012a), there is significant interest in how naturally-produced particles 71 72 impact cloud formation and the optical properties of the resultant clouds (McCoy et al., 2015), and whether this interaction represents a substantial climate feed-73 back (Korhonen et al., 2010). It is well-known that increasing the population of 74 cloud condensation nuclei (CCN) directly increases the opacity of the overlying 75 cloud (Twomey, 1977). Increases in wind speed over the open ocean will enhance 76 the flux of sea spray particles (SSPs) from breaking waves (Hartery et al., 2020). In 77 most regions of the Southern Ocean SSPs are the only local source of ice-nucleating 78 particles (INPs) (DeMott et al., 2016), a region almost entirely devoid of such par-79 ticles (Bigg & Hopwood, 1963). While other, more potent, INPs like dust particles 80 may be entrained into the boundary layer in specific seas (e.g. coastal seas near 81 Patagonia), ice nucleating particles collected on Southern Ocean voyages have a 82 much weaker surface activity than dust particles, which reflects the predominant 83 abundance of sea spray (McCluskey et al., 2018). These particles can have a sub-84 stantial influence on the radiative and physical properties of the resulting cloud. Not 85 only are ice clouds much less opaque (Hu et al., 2010), they are much more likely to 86 precipitate (Borys et al., 2003). Thus, changes in the abundance of SSPs may have 87 significant impacts on cloud radiative properties. 88

One of the challenges in unravelling aerosol-cloud interactions over the South-89 ern Ocean is that the region is frequently covered in cloud (80%) of the time; Haynes 90 et al. (2011)), which leads to difficulties in monitoring low-level clouds (McErlich et 91 al., 2021) and the structure of boundary layer below (Hande et al., 2015). While in 92 situ observational records of radiosondes from Macquarie Island provide rich data 93 on the thermodynamic structure of the Southern Ocean boundary layer (Hande et 94 al., 2012b), a lack of accompanying observations of CCN, INPs, and in situ micro-95 physical properties of low-level cloud leaves a gap in our understanding of how these 96 particles interact with cloud over the Southern Ocean. Previous research, such as 97 the dedicated ACE-1 (Russell et al., 1998), SOCEX (Boers et al., 1998), HIPPO 98 (Wofsy, 2011) and more recently SOCRATES (McFarquhar et al., 2020) campaigns have used aircraft observations to bridge this knowledge gap. However, aircraft can 100 only fly in a limited range of conditions, as the strong vertical wind shear and icing 101 conditions present within boundary layer cloud poses a significant threat. By con-102 trast, ship-based measurements can be made in nearly all conditions. Here, we use 103 measurements on the R/V Tangaroa during a voyage to the Ross Sea in the austral 104 summer of 2018 to establish conditions in which particles near the surface are tur-105 bulently mixed to cloud base (Kremser et al., 2020). Establishing conditions when 106 sea-level measurements are relevant to cloud will enable future research to better ex-107 ploit sea-level measurements in aerosol-cloud interaction studies, and adds value to 108 the growing catalogue of near-surface measurements available from recent voyages. 109

#### 110 2 Measurements

Over the course of a voyage between New Zealand and the Ross Sea, air was drawn from the mast of the R/V *Tangaroa* (~20 m above sea level "a.s.l.") to a shipping container laboratory (~2 m a.s.l.) via 40 m of conductive hose. Within the laboratory, a passive cavity aerosol spectrometer probe (PCASP-100X; Droplet Measurement Technologies) and a differential mobility particle sizer (DMPS, TSI) mea-

sured the ambient concentration of particles suspended in the atmosphere (Kremser 116 et al., 2020). The PCASP measured the number concentration size spectra of par-117 ticles suspended in the boundary layer in 30 size bins  $(0.1-3.0 \ \mu m)$  every minute. 118 The DMPS measured the number concentration size spectra in the size range 0.02-119  $0.3 \ \mu m$  every 10 minutes. Following Modini et al. (2015) and Quinn et al. (2017), we 120 fit three lognormal size distributions to estimate the average diameter and number 121 concentration of Aitken, accumulation and coarse mode particulate. With rare ex-122 ception, coarse mode particulate is almost entirely composed of sea spray particles 123 (SSPs) in the marine environment (Modini et al., 2015; Quinn et al., 2017); hence, 124 we will refer to the coarse mode as the SSP mode throughout the remainder of this 125 work. The PCASP was used exclusively to estimate the average size and abun-126 dance of SSPs, while the DMPS was used for the Aitken and accumulation mode 127 particles. When data from the DMPS were not available, measurements from the 128 PCASP were used to constrain the abundance and size of accumulation mode parti-129 cles. Further details on sampling set-up and analysis, including correction factors for 130 losses through the sampling line and methods for handling contamination from ship 131 exhaust, are described in Hartery et al. (2020) and Kremser et al. (2020). In par-132 allel to the size-resolved particle concentration spectra generated by the SMPS and 133 PCASP, the total number of cloud condensation nuclei (CCN) was measured using 134 a CCN counter (CCNC-100; Droplet Measurement Technologies). The CCN counter 135 sampled from the same sampling conduit that drew ambient air to the PCASP and 136 DMPS. A measurement of the average number of ambient CCN was made twice an 137 hour at intervals of 0.1% supersaturation between 0.2-1.0%. 138

A ceilometer (CHM-15K; Lufft) transmitted pulses of laser light at a wave-139 length of 1064 nm and recorded the total power of light scattered back to the 140 ceilometer per laser pulse from different levels of the atmosphere. Measurements 141 were recorded at a temporal resolution of one record per minute and a vertical reso-142 lution of 15 m. For each record, the instrument also estimated the cloud base height, 143  $z_{CBH}$ . A raw quality control flag provided by the instrument was used to screen for 144 field-of-view contamination from fog or residual precipitation on the outer optical 145 window. A micro-rain radar (MRR-2; Metek) operated in proximity was also used to 146 detect and screen for precipitation events. 147

An Automated Weather Station (AWS) provided by New Zealand MetService 148 was positioned above the bridge of the R/V Tangaroa at 22.5 m. Relevant mea-149 surements included ambient pressure, air temperature, relative humidity, long and 150 shortwave radiation fluxes, wind speed, and wind direction. Measurements from the 151 AWS were corrected to a height of 10 m according to the COARE 3.5 bulk-flux al-152 gorithms (Edson et al., 2013) as detailed in Hartery et al. (2020). The bulk seawater 153 temperature was measured at a depth of 5.5 m below sea level with a thermistor 154 (SBE38; Sea-Bird Scientific). We also used the COARE 3.5 bulk-flux algorithms 155 (Edson et al., 2013) to calculate the sea skin temperature from the bulk tempera-156 ture, accounting for long and shortwave fluxes (Edson et al., 2013). 157

Fifty-seven meteorological balloons were launched during the voyage. The radiosondes (iMet-ABx; InterMet) recorded pressure, relative humidity, temperature and wind speed. In quality control, two of the radiosondes were found to have a faulty relative humidity sensor and one had more than one faulty sensor, leaving 54 useful profiles of the boundary layer. The radiosondes were launched approximately twice daily once the ship was further south than 60° S.

Regional meteorological forecasts were downloaded from the Antarctic Mesoscale Prediction System (AMPS). AMPS initializes a new forecast every twelve hours, with subsequent output provided every three hours. AMPS provides forecasts within several nested spatial grids. However, only forecasts for the outermost spatial grid, "domain 1," were used as it was the only grid which fully contained the <sup>169</sup> ship track. Domain 1 has a horizontal resolution of 24 km and is a  $544 \times 412$  grid <sup>170</sup> centred on 90°S. AMPS uses the Mellor-Yamada-Janjić (MYJ) scheme, a 2.5-level <sup>171</sup> closure model of turbulence, to predict the behavior of the planetary boundary layer <sup>172</sup> (PBL). AMPS calculates the height of the PBL to be the height at which the turbu-<sup>173</sup> lent kinetic energy falls below a pre-determined threshold (Janjic, 2001). To allow <sup>174</sup> for a brief model spin-up, only forecasts between 3–12 hours were used (Jolly et al., <sup>175</sup> 2016).

#### $_{176}$ 3 Methods

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#### 3.1 Classification of the Below-Cloud Layer

The suspended particle cross-sectional surface area, A, was calculated from the number concentration size spectra measured by the PCASP:

$$A(t) = \int \frac{\mathrm{d}n(t, D_p)}{\mathrm{d}\log D_p} \pi \left(\frac{D_p}{2}\right)^2 \mathrm{d}\log D_p \tag{1}$$

Where  $D_p$  is the particle diameter, n is the partial concentration of particles, and 181 t is time. Note that as this is a correlation-based study, a more exact treatment of 182 the interaction of particulate with light which accounts for both Mie and Rayleigh 183 scattering (e.g. Bohren and Huffman (1983)) is not strictly necessary. In addition, 184 such calculations would necessitate a priori information about particle composition 185 and morphology which were not available for this study. The geometric surface area 186 is dominated by the sea spray and accumulation mode particles (97%), on average; 187 Fig. 1d), which the PCASP can readily measure. 188

To classify the below-cloud layer mixing state, we calculated rolling Spearman 189 Rank correlation coefficients centred on each hour of observation between the sea-190 level concentration of aerosol surface area, A(t), and the background-corrected total 191 power of backscattered light received by the ceilometer,  $P_c$ . The Spearman Rank 192 correlation coefficient was used as non-linearities related to the two-way transmis-193 sion of light through an atmospheric layer are likely; however, the Pearson moment 194 correlation coefficient produced qualitatively similar results. Before calculating the 195 correlation coefficients, four quality control measures were implemented to ensure 196 that the calculated correlation coefficients would be meaningful. First, the obser-197 vations were screened based on the ceilometer's quality control flag and the ship 198 contamination flag described in Kremser et al. (2020). Second, only backscattered 199 light retrieved from heights below the 10<sup>th</sup>-percentile of CBH were studied. This 200 step ensured that correlations between backscattered light from a given altitude and 201 sea-level particulate surface area resulted from co-variations of the abundance of 202 aerosol particles at the surface and the selected height and not from variations in the 203 presence of cloud droplets. Next, we performed a signal-to-noise analysis, where the 204 signal-to-noise ratio (SNR) is defined as follows: 205

$$SNR = \frac{P(t,z)}{P_{bg}(t)}$$
(2)

 $P_{bg}$  is the ceilometer's background signal, which the instrument measures at the end of its laser pulse cycle, P is the raw laser power received by the instrument and zis the altitude a.s.l. from which the backscattered light was retrieved. We removed any data points from profiles which had an SNR less than two. Following the SNR analysis, the total power of backscattered light detected by the ceilometer, P(t, z), was corrected for the background signal:

$$P_c(t,z) = P(t,z) - P_{bg}(t)$$
 (3)

Following the initial quality control, rolling correlation coefficients were calculated between A(t) and  $P_c(t, z)$ . This was completed in a two-step process. First, a sub-set of the time-series, T, was defined:

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$$T = \left\{ t_{i-\Delta t/2}, t_{i-\Delta t/2+1}, ..., t_{i+\Delta t/2} \right\}$$
(4)

where  $t_i$  is a specific time in the observation period and  $\Delta t$  defines the temporal width of the sub-set around  $t_i$ . In this work, temporal widths between 1 and 20 hours were studied. As observations were recorded every minute, the sub-set Tcontained at least 60 data points and at most 1200. Spearman Rank Correlation coefficients,  $r_s$ , were then calculated as follows:

$$r_s(x,y) = \frac{cov(rank(x), rank(y))}{\sigma_{rank(x)}\sigma_{rank(y)}}$$
(5)

where rank(x) is a function which assigns an integer ranking to each value of a set 224 x; cov(x, y) is the covariance of two sets of data, x and y; and  $\sigma_x$  is the standard 225 deviation of the set x (Spearman, 1904). Here, x and y are the sub-sets of A(t) and 226  $P_c(t,z)$  defined by T. Applying equations 4 & 5 to the entire time-series forms a 227 matrix,  $R_{\Delta t}(t,z)$ . A detailed justification of this range of time-scales is provided 228 in Section 4.3. Two additional post-processing procedures were implemented after 229 the correlations were calculated. If a subset, T, contained less than 20 valid data 230 points, then the correlation coefficient was labelled as not a number. For the remain-231 ing data, a significance test was performed for each correlation value to ensure that 232 the value was significantly larger than zero (p < 0.05). If the calculated correlation 233 coefficient failed the significance test, it was re-assigned a value of zero. 234

Once the fully quality-controlled correlation analysis had been completed, we 235 developed a simple metric to classify the mixing state of the atmospheric layer below 236 the lowest observed cloud. First, the average below-cloud correlation coefficient, 237  $\overline{r_{bc}}$ , was calculated. When  $\overline{r_{bc}} > 0$  (p < 0.05), the below-cloud layer was classified: 238 "well-mixed". In such cases, particles observed at sea-level were considered to be 239 well-mixed into the overlying cloud. However, if the average below-cloud correlation 240 coefficient didn't exceed zero, then the surface layer below-cloud layer was classified: 241 "poorly-mixed". In these cases, the atmospheric layer at the surface was assumed 242 to be decoupled from the overlying cloud. Finally, if there were insufficient data 243 points in any sub-set of the time series, T, then the correlation analysis was unable 244 to classify the mixing state of the below-cloud layer for that period. 245

#### <sup>246</sup> 3.2 Validation

To validate the proposed methodology and classification metric, we compared results to four separate methods of determining the mixing state of the below-cloud layer. The first two methods were variations on a conventional radiosonde analysis, one was a surface-based method and the final method was a model-based method.

We compared the classification of the below-cloud mixing state according to 251 the correlation metric to two methods for detecting boundary layer decoupling based 252 on radiosonde profiles. The first method searched for maxima in the virtual poten-253 tial temperature gradient  $(\partial \theta_v \ \partial z^{-1})$  (Hande et al., 2012b). If a local maxima in 254 the virtual potential temperature gradient was detected and found to exceed 10 K 255 km<sup>-1</sup>, then the height at which this occurred was labelled as the main inversion, 256 or the boundary layer height. The method then searched for secondary maxima 257 larger than 5 K  $\rm km^{-1}$  below the main inversion. If secondary inversions exist, then 258 the boundary layer is decoupled (Hande et al., 2012b). To be consistent with our 259 methodology, which can only classify the atmospheric layer below the lowest ob-260 served cloud, the below-cloud layer was only labelled as decoupled if a secondary 261 inversion was located between the surface and the cloud. 262

A second method for detecting below-cloud decoupling was adapted from Truong et al. (2020). In this method, a main inversion was only identified if a local maximum in the virtual potential temperature gradient exceeded 14 K km<sup>-1</sup>. To detect decoupling, the decoupling parameter  $\mu$  was studied (Truong et al., 2020; Yin & Albrecht, 2000). The decoupling parameter,  $\mu$ , is defined as follows:

$$\mu = -\left(\frac{\partial\theta}{\partial z} - \frac{0.608\theta}{1 + 0.608r}\frac{\partial r}{\partial z}\right) \tag{6}$$

Where r is the water vapour mixing ratio and  $\theta$  is the potential temperature. Yin 269 and Albrecht (2000) devised  $\mu$  for their study of "transition layers" in the bound-270 ary layer, as it is more sensitive to changes in the water vapour mixing ratio than 271 the vertical gradient of virtual potential temperature and is therefore more likely 272 to detect subtle boundary layer features like decoupling. Decoupling of the bound-273 ary layer over the Southern Ocean was only detected when a value of  $\mu$  exceeded 274 2.5 times its average value throughout the boundary layer (Truong et al., 2020). 275 To be consistent with our method, we adapted this method to only classify the 276 below-cloud layer as decoupled if the threshold for  $\mu$  was exceeded in the below-277 cloud layer (Truong et al., 2020). We used a simple optimization methodology to 278 determine which combination of time-scale,  $\Delta t$ , and correlation threshold,  $r_t$ , best 279 predicted the state of coupling between the surface and cloud layers as compared to 280 the reference methods (Hande et al., 2012b; Truong et al., 2020). 281

To provide a benchmark for our methodology, we compared the optimized 282 performance of the correlation-based method against another surface-based method-283 ology for defining the mixing state of the below-cloud layer (Jones et al., 2011). 284 Briefly, if the difference in height between the observed cloud base height (CBH) 285 and lifted condensation level (LCL) exceeded 150 m, then the below-cloud layer was 286 considered to be decoupled from the cloud (and well-mixed otherwise). For these 287 calculations, the LCL was calculated from the AWS measurements, where the LCL 288 represents the height at which a cloud is expected to form based on a parcel of air 289 adiabatically ascending through a well-mixed boundary layer (Romps, 2017). Here, 290 we used the 1-hour averaged LCL and for consistency, the 10<sup>th</sup>-percentile of CBH 291 within each hour. 292

<sup>293</sup> While the radiosonde profiles collected throughout the voyage provided a ro-<sup>294</sup> bust benchmark for the new methodology, radiosonde data were available at most <sup>295</sup> twice-a-day. To increase our confidence in the methodology, we compared its classifi-<sup>296</sup> cation of the below-cloud layer to near-surface measures of atmospheric stability. We <sup>297</sup> used two measures of near-surface atmospheric stability: the square of the Brunt-<sup>298</sup> Väisälä Frequency, N, and the 10-m wind speed. Values of N<sup>2</sup> were calculated from <sup>299</sup> the AWS measurements and the COARE 3.5 bulk-flux algorithms:

$$N^2 = \frac{g}{\theta_v} \frac{\partial \theta_v}{\partial z} \tag{7}$$

where q is the gravitational acceleration,  $\theta_v$  is the virtual potential temperature, 302 and z the height above sea level. This stability analysis was combined with fore-303 casts from AMPS to quantitatively define conditions in which aerosol-cloud coupling 304 was expected. As a coarse proxy for aerosol-cloud coupling, we investigated the 305 difference in the LCL and the predicted planetary boundary layer height (PBL) in 306 the AMPS forecasts. If the planetary boundary layer exceeded the lifted conden-307 sation level, then aerosol particles measured at the ocean surface were considered 308 well-mixed to the minimum height where cloud could have occurred. 309

#### 310 4 Results

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#### 4.1 Time Series Analysis

Throughout the voyage to and from the Ross Sea (voyage track shown in Fig. 1a), the number-size distribution of particulate was predominantly trimodal as

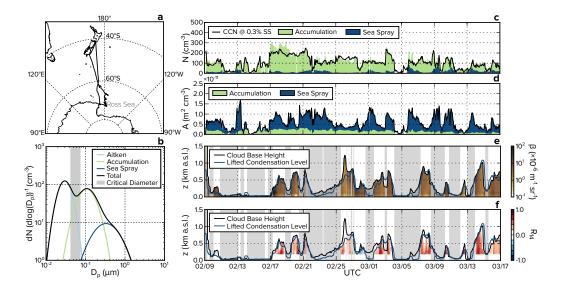


Figure 1. (a) The track of the R/V Tangaroa during the Marine Environment and Ecosystem Voyage. (b) A typical size distribution for particles in the Southern Ocean. The expected range of cloud activation diameters for marine stratus is shown in grey. (c) The sea-level abundance of sea spray particles (SSPs; blue filled region) and accumulation mode particles (green filled region) is compared to the abundance of cloud condensation nuclei (CCN) at a supersaturation of 0.3% (black line). (d) The abundance of suspended surface area was calculated from the measured particle size distributions (Eq. 1). (e) A contour plot of the attenuated backscatter coefficient measured by the CHM-15K ceilometer. The lifted condensation level ( $z_{LCL}$ ) and cloud base height ( $z_{CBH}$ ) are also shown for reference. (f) Spearman Rank correlation coefficients between the sea-level abundance of particulate surface area and ceilometer backscatter are shown. Time periods when the ceilometer optical window was obscured, the cloud base was below 200 m, fog was present, or the aerosol sampling system was contaminated by ship exhaust are shaded.

seen in Fig. 1b. The representative number–size distribution shown in Fig. 1b was 314 constructed by taking voyage wide averages of the total number, width and median 315 size of the individual modes that were fit to the observations. The appearance of 316 these modes is consistent with previous observations in marine settings (Bates et 317 al., 1998; Quinn et al., 2017). A large majority of the particles in the smallest two 318 modes, the Aitken (30 nm,  $\sigma = 1.4$ ) and accumulation modes (100 nm,  $\sigma = 1.6$ ), are 319 thought to be produced as a single mode from homogeneous nucleation of volatile 320 sulfate species, with mode separation occurring as a result of cloud-processing 321 (Hoppel et al., 1986). These particles are nucleated in-situ from the condensation 322 of oxidized marine gasses and grow via self-coagulation and condensation. In con-323 trast, sea spray particles (400 nm,  $\sigma = 2$ ) are directly generated from breaking ocean 324 waves, and tend to be much larger than particles in the Aitken and accumulation 325 mode (Prather et al., 2013). Note that size statistics presented in this section have 326 been corrected to a relative humidity of 80%. For sulfate and sea spray particles, a 327 particle at a relative humidity of 80% is approximately twice as large compared to 328 when it is dry (Gerber, 1985). 329

A representative size distribution of particles observed in the Southern Ocean marine boundary layer at a relative humidity of 80% is shown in Fig. 1b. The bifurcation of the Aitken and accumulation modes occurs when these particles pass through non-precipitating cloud, since only the largest particles will be acti-

vated (Hoppel et al., 1986). Previous research has shown that the supersaturation 334 of water vapour within nascent marine stratus is relatively modest (<0.3%; Hegg 335 et al. (2009)). An estimation of the activation diameter based on a supersaturation 336 of 0.3%, and a range of particle hygroscopicity parameters is also shown in Fig. 1b. 337 The estimation of the range of activation diameter is based on the  $\kappa$ -Köhler model 338 for a range of expected hygroscopicity values (Petters & Kreidenweis, 2007). This 339 coincides well with the local minimum between the Aitken and accumulation mode, 340 supporting the cloud-processing hypothesis of Hoppel et al. (1986). 341

342 Fig. 1c displays the number of particles in both the accumulation and sea spray modes, as these are the only particles relevant to cloud formation. This is 343 compared to the number concentration of CCN measured at a fixed supersaturation 344 of 0.3%. As expected, these two measurements are highly correlated. Across the 345 entire voyage, SSPs did not comprise a substantial fraction of CCN (14%). However, 346 in the latter half of the voyage we encountered several low pressure systems. These 347 cyclones were accompanied by high winds, resulting in substantial wave-breaking 348 and subsequent SSP generation in the region. This led to an enhanced relevance of 349 SSPs to the total CCN population (20%). 350

Fig. 1d shows the abundance of suspended particle surface area. Despite the 351 relatively low abundance of SSPs by number, the total amount of particulate surface 352 area is strongly dominated by variations in their abundance. In Fig. 1e, the time se-353 ries of attenuated backscatter profiles measured by a coincident ceilometer is shown, 354 along with rolling averages of cloud base height and the lifted condensation level. 355 As demonstrated both empirically and theoretically, if the difference between cloud 356 base height and lifted condensation level is less than 150 m, the below-cloud layer 357 can be considered well-mixed (Jones et al., 2011). As a result, it is clear that there 358 was significant coupling between the surface layer and overlying cloud for much of 359 the time-series. 360

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### 4.2 New Classification Methodology

We used the Spearman Rank correlation analysis between suspended particle 362 surface area at sea-level (Fig. 1d) and ceilometer backscatter from particles overhead 363 (Fig. 1e) to assess whether our measurements at the surface were representative of 364 the below-cloud population of CCN. Fig. 1f displays strong correlations between 365 these two quantities over time-scales of 14 hours when fog, precipitation, or con-366 tamination from ship exhaust did not inhibit the analysis. This suggests that the 367 Southern Ocean boundary layer was consistently well-mixed throughout this mea-368 surement campaign. We note that correlation coefficients could not be calculated 369 below 200 m, as these data typically failed the SNR analysis. While one would nor-370 mally expect a large backscattered signal close to a lidar, and thus a high SNR, the 371 returning backscatter is not well-aligned with the FOV of the receiving optics in the 372 near-range, resulting in a low SNR. As an additional control, the significance of the 373 calculated correlation coefficients was assessed with a two-way t-test. 374

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## 4.3 Comparison to Radiosonde Analysis

To validate the correlation analysis and establish the most accurate time-376 scale for calculating correlation coefficients, we analyzed the 57 radiosonde profiles 377 recorded throughout the voyage. For each radiosonde, we used two gradient methods 378 to detect whether the surface layer was decoupled from the cloud layer (Hande et 379 al., 2012b). Out of the 57 radiosondes, three could not be used for analysis due to 380 faulty sensors, and 21 were launched when the cloud base height was below 200 m. 381 In such cases, there was insufficient ceilometer data to perform the correlation anal-382 ysis, as the power of the returning backscatter was on the same order of magnitude 383

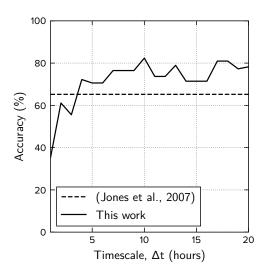


Figure 2. The accuracy with which the new correlation metric correctly classified the mixing state of the below-cloud layer for different correlation time-scales,  $\Delta t$ . The accuracy of the method was calculated in reference to two radiosonde analyses which classified the mixing state of the below-cloud layer based on thermodynamic gradients (N = 26; non-precipitating conditions, no fog, CBH>200 m) (Hande et al., 2012b; Truong et al., 2020). The accuracy of the proposed method can be compared to the accuracy of another ground-based methodology of determining the state of below-cloud mixing (Jones et al., 2011).

as the instrument noise due to the FOV effects described earlier. In three additional cases, the gradient methods did not detect a boundary layer. As a result, there were only 30 radiosondes available for which the correlation analysis was valid. In these remaining 30 cases, the radiosonde-based methods of classifying the below-cloud layer differed only slightly. Overall, the below-cloud layer appeared well-mixed in 83% of profiles according to the criteria of Truong et al. (2020), and 90% of profiles according to Hande et al. (2012b).

The classifications of the below-cloud mixing state defined by the radiosonde 391 analyses were used to define the optimal time-scale,  $\Delta t$ , and threshold,  $r_t$ , for the 392 correlation analysis. For consistency, only the profiles for which both radiosonde 393 methods agreed on the mixing state of the below-cloud layer were used as a refer-394 ence when calculating the accuracy (N = 26). In Fig. 2, the accuracy with which the 395 correlation analysis determined the mixing state of the below-cloud layer is shown 396 as a function of time-scale. Across all time-scales, the threshold for detecting a well-397 mixed below-cloud layer was  $\overline{r_{bc}} > 0$ , where  $\overline{r_{bc}}$  is the average correlation coefficient 398 between sea-level and the  $10^{\text{th}}$ -percentile of cloud base height. As a benchmark, we 399 have also shown the accuracy of another ground-based method for determining the 400 below-cloud mixing state (Jones et al., 2011). 401

Fig. 2 demonstrates that the accuracy of the correlation-based method increased from 35% to 76% as the time-scale increased, until time-scales of 7 hours or longer were reached. Differences in accuracy at time-scales beyond 7 hours were negligible considering the sample size (N = 26). The increase in accuracy with increasing time-scale is a direct result of increasing the number of samples in the subset T (defined in Eq. 4) used for calculating the correlation coefficient. While shorter time-scales are likely more representative of the time-scale of turbulence,

Statistic	P (hPa)	T(K)	$T_d$ (K)	U (m s <sup>-1</sup> )
RMSD Bias	$2.0 \\ 0.3$	1.2	3.2 -0.8	2.7 0.7
	1	0.96	0.87	0.74

**Table 1.** This table summarizes statistics comparing measurements from radiosondes launchedthroughout the voyage and predictions from AMPS below 3 km (a.s.l.).

RMSD: Root Mean Squared Deviation

 $\mathbf{R}^2:$  Pearson Correlation Coefficient

there is also a higher likelihood that the remaining noise in the ceilometer obser-409 vations will result in weaker correlations which fail the two-way t-test (p > 0.05). 410 Increasing the time-scale resulted in more consistent correlation coefficients across 411 time-scales and more statistically significant results overall. The accuracy of this 412 method also suggests that despite there being longer time-scale phenomena which 413 could also correlate particulate surface area and backscatter (e.g. frontal systems, 414 convective forcing at cloud top, precipitation, turbulent perturbations of relative 415 humidity, air mass history, etc.), these phenomena are not likely to result in sub-416 stantial misclassification of the below-cloud layer. However, considering that such 417 long time-scale phenomena do exist and may be more prevalent in other regions or 418 observation periods, correlation coefficients calculated over time-scales beyond those 419 presented here should be avoided as false positives and false negatives are likely to 420 become more abundant. 421

Finally, we compared the accuracy of our new methodology to another method 422 of remotely classifying the below-cloud mixing state (Jones et al., 2011). In this 423 case, the referenced method was only 65% accurate at determining the mixing state 424 of the below-cloud layer, whereas the proposed method was 76% accurate when 425 correlation time-scales greater than 7 hours were considered. While the set of ra-426 diosondes for which we could compare both methods was quite limited (N = 26), 427 these results suggest that the proposed method more accurately classified the mixing 428 state of the below-cloud layer than the referenced method (p < 0.05). 429

Overall, the correlation analysis found that the below-cloud layer was well-430 mixed for 14% of the entire time series and poorly-mixed just 5%. Fog was found to 431 occur 7% of the time, where fog was diagnosed when the relative humidity was mea-432 sured to be 100% and cloud base was less than 50 m. The remaining portion of the 433 time series could not be analyzed (74%), as one or more of the following occurred: 434 the ceilometer's quality control flag was raised; the ship exhaust contaminated the 435 aerosol sample; or, the cloud base was below 200 m but greater than 50 m, such 436 that the entire profile of below-cloud backscatter failed the SNR analysis due to a 437 lack of overlap between the FOV of the ceilometer's optical system and the return-438 ing backscatter. While this may seem like a large loss of the time-series, if a given 439 radiosonde was only representative of conditions for the hour of measurements in 440 which it was operating, then the radiosonde analysis provided data for just 7% of 441 the time series. With the proposed correlation analysis, we were able to classify the 442 boundary layer for 26% of the time series, a marked improvement. 443

444

#### 4.4 Comparison to Stability Analysis

While the comparison to conventional radiosonde analyses provided evidence that the correlation analysis accurately classified the below-cloud mixing state, it still seemed prudent to evaluate the analysis against other metrics of atmospheric mixing. Here, we examine expected rates of occurrence for aerosol-cloud coupling
 based on AMPS forecasts for the period of study.

First, the forecasts were compared to our observations from all available radiosondes, except those with non-functioning RH sensors. Table 1 provides a summary of various comparison statistics between forecasts and measurements. The statistics were only calculated below 3 km to restrict the comparison to relevant planetary boundary layer (PBL) and lifted condensation level (LCL) heights. These are presented in Table 1.

Overall, observed and measured values of the selected variables were rea-456 sonably well correlated. However, there were minor biases worth mentioning. In 457 Table 1, statistically significant biases between modelled and measured values of 458 pressure, dew point temperature, and wind speed were observed (p < 0.001). Within 459 AMPS, the height of the PBL is determined according to the turbulent kinetic en-460 ergy profile (Janjic, 2001). This implies that the height of the PBL may have been 461 under-estimated by AMPS. The dew point temperature was also negatively biased 462 as a result of the over-abundance of water vapour in the AMPS boundary layer rel-463 ative to observations. This implies that the LCL was also under-predicted. Still, 464 considering the spatial and temporal scale of the AMPS forecasts, the agreement be-465 tween model and measured values was quite good, and highly statistically significant 466 (p < 0.001).467

In Fig. 3a, the frequency of occurrence with which the depth of the planetary 468 boundary layer exceeded the lifted condensation level is shown, based on forecasts 469 from AMPS between 40 - 70 S (excluding areas less than 100 km from a coast). 470 The frequency of occurrence is shown as a function of two variables which are often 471 used to describe the stability of the near-surface atmospheric layer: the square of the 472 Brunt-Väisälä frequency,  $N^2$ , and the 10-m wind speed,  $U_{10}$ . The results in Fig. 3a 473 demonstrate that in near-neutral stability (N<sup>2</sup>  $\sim$  0) and weak winds (U<sub>10</sub>  $\sim$  0), the 474 layer below the LCL was less-likely to be well-mixed, as the PBL was too shallow. 475 However, in all other cases, the boundary layer was likely well-mixed as the PBL 476 exceeded the LCL. Note that there was still a small percentage of the time when the 477 layer below the LCL was well-mixed despite the surface layer being near-neutrally 478 stable. While a well-mixed boundary layer would not be expected in such cases, the 479 near-surface layer is typically much shallower than the LCL, and is therefore not 480 always a perfect determinant of the mixing state of the entire layer below the LCL. 481 However, it is clear from Fig. 3a that in most other conditions, the layer below the 482 LCL is almost guaranteed to be well-mixed. Overall, The AMPS analysis in Fig. 3a 483 provides a general rule of thumb: if the 10 m wind speed exceeds 8 m  $\rm s^{-1}$  then the 484 boundary layer will be well-mixed to the LCL, regardless of the near-surface stabil-485 ity. 486

In Fig. 3c and d, the classification of the below-cloud layer according to the 487 correlation metric is shown for time-scales of 7 and 14 hours, respectively. We can 488 see that despite the accuracy with which the correlation analysis at 7 hour and 14 489 hour time-scales classified the mixing state of the boundary layer (Fig. 2), the cor-490 relation metric calculated over a 14 hour time-scale provided a more qualitatively 491 consistent result with the AMPS analysis. In comparing Figs. 3c & d, it is clear 492 that the correlation metric calculated over a time-scale of 7 hours misclassified the 493 boundary layer more frequently, as a poorly mixed boundary layer is not expected 494 to occur at all if  $N^2 < -5 \times 10^{-3} \text{ s}^{-2}$  or  $U_{10} > 8 \text{ m s}^{-1}$  (Fig. 3a). Barring a few 495 exceptions, Fig. 3d shows that the correlation metric at a time-scale of 14 hours 496 typically only classified the below-cloud layer as decoupled only when the stability of 497 the near-surface layer was near-neutral and winds were less than 8 m s<sup>-1</sup>, consistent 498 with the AMPS analysis. While only two time-scales are presented here, analysis at 499 all time-scales longer than 14 hours produced qualitatively similar results. Finally, in 500

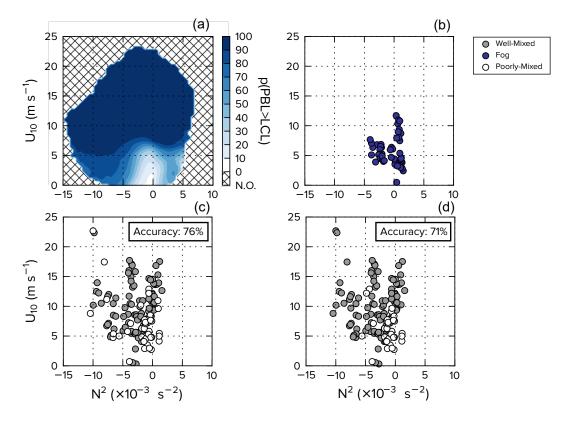


Figure 3. (a) The frequency with which the height of the planetary boundary layer (PBL) predicted by AMPS was higher than the lifted condensation level (LCL) over the open Southern Ocean (40–70 S, >100 km from coastline) in February and March 2018 (N.O. = Conditions occurred less than frequently than 0.001%). (b) The occurrence of fog (CBH < 50 m, RH = 100%). (c) The classification of the below-cloud layer based on correlation coefficients calculated over 7-hour timescales (non-precipitating conditions; CBH > 200 m). The measure of accuracy is in reference to the radiosonde analyses (Fig. 2). (d) As in (c), but for a time-scale of 14-hours.

Fig. 3b, it is clear that fog tended to occur only in both near-neutral stability (N<sup>2</sup> >  $-5 \times 10^{-3} \text{ s}^{-2}$ ) and low winds, or stable conditions, consistent with advection fog.

One limitation of this analysis is that a cloud is not necessarily guaranteed to 503 occur at the LCL. As such, a direct quantitative comparison between Figs. 3a, c & 504 d is not possible, as cloud was always occurring in the subset of data we were able 505 to analyze but may not have been occurring in the AMPS forecasts. Still, we found 506 this figure to be a useful qualitative reference for our methodology. In addition, it 507 demonstrates that even though the fraction of the time-series available for analysis 508 via the correlation metric is low  $(26\%; \sim 10 \text{ days of observations})$ , the conditions 509 encountered within this subset of the data are representative of the wide set of con-510 ditions forecast by AMPS. As a result, statistics presented in the previous section 511 can be used to conclude that in non-precipitating conditions, the below-cloud layer 512 over the Ross Sea was likely well-mixed 54% of the time, poorly-mixed 19% of the 513 time, and contained fog 27% of the time. 514

#### 515 5 Discussion

In this work, we were interested in understanding how often aerosol particles 516 measured near the surface of the ocean were relevant to low cloud formation over 517 the Southern Ocean. We proposed a new methodology, based on the correlation 518 of particle surface area and ceilometer backscatter, which identified when aerosol 519 particles observed at the surface were available to the lowest observed cloud. To 520 validate the proposed methodology, we needed an accurate reference classification of 521 the boundary layer against which we could compare our results. Here, we modified 522 523 two radiosonde-based methodologies which determined the mixing state of the entire boundary layer based on gradients of thermodynamic variables (Hande et al., 2012b; 524 Truong et al., 2020). Radiosonde-based methodologies were selected as the preferred 525 reference methodology, as the observations were sensitive to fine thermodynamic 526 changes in the boundary layer. These reference methodologies were modified to sim-527 ply determine the mixing state of the atmosphere between the surface and the base 528 of the lowest observed cloud. This allowed us to optimize the parameters of our pro-529 posed methodology (correlation time-scale, threshold of correlation strength) such 530 that the predicted mixing state of the below-cloud layer best matched the referenced 531 radiosonde methodologies. In the comparison (Fig. 2), the proposed correlation-532 based method correctly classified the mixing state of the below-cloud layer  $76 \pm 4\%$ 533 of the time for correlation time-scales greater than 7 hours. The accuracy of our 534 method was then compared to a more simple metric for classifying the mixing state 535 of the below cloud layer, which was only accurate 65% of the time (Jones et al., 536 2011). 537

In a more qualitative comparison (Fig. 3), the classification of the below-cloud 538 mixing state by the proposed methodology was also shown to be consistent with 539 surface-based measurements of atmospheric stability and model predictions of tur-540 bulence in a wide range of conditions. The high accuracy of the new methodology's 541 predictions in comparison to radiosonde-based methods, in situ observations of near-542 surface atmospheric stability, and model forecasts of boundary layer turbulence gives 543 us high confidence that the proposed method is accurate even when reference data is 544 not available. 545

With the accuracy of our proposed methodology validated against multiple 546 methods of determining the below-cloud mixing state, we can compare statistics to 547 previous observations in the Southern Ocean. We find that while the below-cloud 548 layer was often well-mixed, this was not always guaranteed. It is well-known that 549 the marine boundary layer can stratify into a near-surface boundary layer and a 550 sub-cloud layer (Garratt, 1994). In fact, radiosondes launched from Macquarie Is-551 land (54.62°S, 158.85°E) over the past two decades found that the boundary layer 552 was well-mixed just 17.8% of the time (Hande et al., 2012b). In contrast, our time 553 series analysis showed that in non-precipitating conditions, the below-cloud layer 554 was well-mixed 54% of the time. This seems to be in stark contrast to the accuracy 555 data presented in Fig. 2. However, the difference in frequency of occurrence comes 556 primarily from a difference in the definition of decoupling. The method presented in 557 this work was only designed to detect whether the boundary layer was well-mixed up 558 to the lowest cloud. In contrast, the method used to analyze radiosondes launched 559 from Macquarie Island was designed to detect decoupling throughout the entire 560 boundary layer (Hande et al., 2012b). However, multi-layer clouds are frequently 561 observed over the Southern Ocean (Hande et al., 2012b). In such settings, the inver-562 sion atop the lowest cloud will tend to decouple the atmospheric layer beneath the 563 cloud from the rest of the boundary layer. Despite this decoupling, near-surface air 564 is typically still well-mixed up to the lowest cloud, as cloud was often present in the 565 atmospheric layer beneath the decoupling height (Hande et al., 2012b). As a result 566 of this inconsistency, statistics retrieved from the radiosonde analysis at Macquarie 567

Island cannot be easily used to infer information about aerosol-cloud coupling and
 are not directly comparable to this study.

Overall, the results from the correlation analysis highlight that particles are al-570 most always available to the lowest cloud (Fig. 1f). The percentage of time in which 571 aerosol-cloud coupling occurred within the valid section of our time series is simply 572 the sum total of the rates of occurrence of fog and a well-mixed below-cloud layer: 573 81%. Forecasts from AMPS tend to agree, as the layer of the atmosphere below the 574 LCL was found to be well-mixed 84% of the time over the Southern Ocean through-575 576 out February and March 2018. Based on the good agreement between these methods of defining the below-cloud mixing state, we are confident in concluding that sea 577 spray particles are available to low-level cloud over the Southern Ocean more than 578 80% of the time in austral summer. As Kuma et al. (2020) noted, the ability to 579 correctly predict the occurrence of low cloud is a critical necessity for improving 580 the Southern Ocean shortwave radiation bias. The proposed method increases our 581 understanding of these low clouds and the particles which help form them. 582

For instance, we found that the number of CCN at a supersaturation of 0.3%583 was consistent with the number of particles in the accumulation and sea spray mode. 584 As a supersaturation of 0.3% is the expected water vapor supersaturation within 585 marine stratocumulus (Hegg et al., 2009), this suggests that sea-level observations 586 may provide a good constraint on the number of cloud droplets in a wide variety 587 of conditions. We found that despite being readily-available to nascent clouds, sea 588 spray particles were typically outnumbered by smaller, cloud-processed accumulation 589 mode particles (Fig. 1c), consistent with previous studies (Quinn et al., 2017). How-590 ever, in addition to abundance, the ice-nucleating ability of particles is known to be 591 a strong determinant of cloud phase and albedo: a climate model which determined 592 the primary nucleation of ice within low-level clouds according to the abundance and 593 type of boundary layer ice-nucleating particles found that predictions of cloud opac-594 ity were significantly more accurate in the cold sector of Southern Ocean cyclones 595 relative to simpler glaciation schemes (Vergara-Temprado et al., 2018). Though less 596 numerous than accumulation mode particles, sea spray particles are thought to be 597 the only local source of ice-nucleating particles (DeMott et al., 2016) in a region 598 that is often devoid of more potent ice nuclei (e.g. dust) (McCluskey et al., 2018). 599 This study highlights that sea spray particles are available to many more cloud 600 systems than just within the cold sector of cyclones. As a result, climate models 601 which implement glaciation schemes that connect the primary nucleation of ice to 602 the microphysical properties of aerosol particles will likely see more widespread im-603 provement to the Southern Ocean shortwave radiation bias. It also highlights that 604 should models adopt more complex models of cloud glaciation, then they must also 605 more carefully parameterize the flux of sea spray particles (Hartery et al., 2020). 606

The new method does not come without limitations, however. Depending on 607 the FOV of the ceilometer's optical receiver, the ability of the analysis to analyze 608 below-cloud coupling in low cloud settings (CBH < 200 m) can be severely impaired 609 due to an incomplete overlap of the returning laser beam and the receiver's FOV. 610 In addition, though we have provided a reasonably comprehensive validation of the 611 appropriate time-scale for calculating correlation coefficients and the threshold for 612 classification of the below-cloud layer, there are potentially instances where the cor-613 relation analysis could trigger false positives and false negatives in other synoptic 614 settings. These include, but are not limited to, frontal systems, convective forcing 615 at cloud top, precipitation, turbulent perturbations of relative humidity, air mass 616 history, etc. Still, given the accuracy of the methodology as quantitatively compared 617 to the radiosonde analyses, and qualitatively to a forecast analysis, we are confident 618 in the results presented as they pertain to this specific region and period of study. 619 As an added benefit, the proposed method also uses instruments which function 620

nearly autonomously, with little need for oversight or on-site personnel. In contrast, 621 radiosonde programs require highly-trained personnel and can only be launched is a 622 limited set of meteorological conditions. It becomes exceedingly difficult to success-623 fully launch a radiosonde once winds surpass 15 m s<sup>-1</sup>, and potentially dangerous 624 when aboard a research vessel in unfavorable wave conditions. As a result, statis-625 tics of boundary layer mixing collected from radiosonde programs are likely skewed 626 towards calm conditions. Finally, it is worth mentioning that there are likely un-627 intended, negative environmental consequences of leaving irretrievable radiosonde 628 packages in the Southern Ocean. 629

#### 630 6 Conclusions

In this work we presented a new technique for determining the state of bound-631 ary layer mixing based on the value of the Spearman Rank correlation coefficient 632 calculated between sea-level observations of suspended particle surface area and 633 ceilometer backscatter. When data was available, these correlations were often high, 634 implying that particles measured at sea-level were well-mixed throughout the bound-635 ary layer and were therefore readily-available to nascent, low-level cloud. From this 636 analysis, a simple metric was created to diagnose whether coupling occurred or not. 637 This revealed that in non-precipitating conditions the boundary layer was well-638 mixed 54% of the time, contained fog 27% of the time, and was poorly-mixed just 639 19% of the time. This simple metric based on the correlation analysis was compared 640 to two conventional radiosonde analyses. The correlation-based metric accurately 641 classified the mixing state of the boundary layer 76% of the time when correlation 642 coefficients were calculated over periods longer than 7 hours. This is a noticeable im-643 provement over the accuracy of a simpler ground-based method (65%). In addition, 644 the frequency of occurrence of below-cloud mixing estimated by the correlation-645 based metric was qualitatively consistent with an analysis of mixing based on the 646 near-surface stability within regional forecasts. We estimate that aerosol will have 647 a direct pathway into low cloud either through a well-mixed below cloud layer or 648 surface-level fog, 81% of the time when clouds are present. Thus, in situ sea-level 649 observations of particulate offer substantial insight into cloud formation over the 650 Southern Ocean in a wide set of conditions. 651

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667 amps.html.

## 668 References

669 Bates, T. S., Huebert, B. J., Gras, J. L., Griffiths, F. B., & Durkee, P. A. (1998).

670 671	International global atmospheric chemistry (igac) project's first aerosol char- acterization experiment (ace 1): Overview. Journal of Geophysical Research:
672	Atmospheres, $103(D13)$ , $16297-16318$ .
673	Bigg, E. K., & Hopwood, S. C. (1963). Ice nuclei in the antarctic. Journal of the
674	Atmospheric Sciences, $20(3)$ , 185-188. doi: $10.1175/1520-0469(1963)020<0185$ :
675	INITA>2.0.CO;2
676	Bodas-Salcedo, A., Williams, K. D., Ringer, M. A., Beau, I., Cole, J. N., Dufresne,
677	JL., Yokohata, T. (2014). Origins of the solar radiation biases over the
678	Southern Ocean in CFMIP2 models. Journal of Climate, 27(1), 41–56.
679	Boers, R., Jensen, J., & Krummel, P. (1998). Microphysical and short-wave ra-
680	diative structure of stratocumulus clouds over the southern ocean: Summer
681	results and seasonal differences. Quarterly Journal of the Royal Meteorological
682	Society, 124(545), 151–168.
683	Bohren, C. F., & Huffman, D. R. (1983). Absorption and scattering by a sphere. Ab-
684	sorption and scattering of light by small particles.
685	Borys, R. D., Lowenthal, D. H., Cohn, S. A., & Brown, W. O. J. (2003). Moun-
686	taintop and radar measurements of anthropogenic aerosol effects on snow
687	growth and snowfall rate. $Geophysical Research Letters, 30(10).$ doi:
688	10.1029/2002GL016855
689	DeMott, P. J., Hill, T. C. J., McCluskey, C. S., Prather, K. A., Collins, D. B., Sulli-
690	van, R. C., Franc, G. D. (2016). Sea spray aerosol as a unique source of ice
691	nucleating particles. Proceedings of the National Academy of Sciences, 113(21),
692	5797–5803. doi: $10.1073/\text{pnas}.1514034112$
693	Edson, J. B., Jampana, V., Weller, R. A., Bigorre, S. P., Plueddemann, A. J.,
694	Fairall, C. W., Hersbach, H. (2013). On the exchange of momentum over
695	the open ocean. Journal of Physical Oceanography, $43(8)$ , 1589–1610. doi:
696	10.1175/JPO-D-12-0173.1
697	Garratt, J. (1994). Review: the atmospheric boundary layer. Earth-Science Reviews,
698	37(1), 89 - 134. Retrieved from http://www.sciencedirect.com/science/
699	article/pii/0012825294900264 doi: https://doi.org/10.1016/0012-8252(94)
700	90026-4
701	Gerber, H. E. (1985). Relative-humidity parameterization of the Navy Aerosol Model
702	(NAM) (Tech. Rep.). Washington, DC: Naval Research Lab.
703	Hande, L., Siems, S., & Manton, M. (2012a). Observed trends in wind speed over
704	the Southern Ocean. Geophysical Research Letters, $39(11)$ . doi: 10.1029/
705	2012GL051734
706	Hande, L., Siems, S., Manton, M., & Belusic, D. (2012b). Observations of wind
707	shear over the southern ocean. Journal of Geophysical Research: Atmospheres,
708	117(D12). doi: $10.1029/2012JD017488$
709	Hande, L. B., Siems, S. T., Manton, M. J., & Lenschow, D. H. (2015). An eval-
710	uation of cosmic radio occultation data in the lower atmosphere over the
711	southern ocean. Atmospheric Measurement Techniques, $8(1)$ , 97–107. Re-
712	trieved from https://amt.copernicus.org/articles/8/97/2015/ doi:
713	10.5194/amt-8-97-2015
714	Hartery, S., Toohey, D., Revell, L., Sellegri, K., Kuma, P., Harvey, M., & McDon-
715	ald, A. (2020). Constraining the surface flux of sea spray particles from the
716	southern ocean. Journal of Geophysical Research: Atmospheres, 125(4). doi:
717	10.1029/2019JD032026
718	Haynes, J. M., Jakob, C., Rossow, W. B., Tselioudis, G., & Brown, J. (2011). Major
719	characteristics of southern ocean cloud regimes and their effects on the energy
720	budget. Journal of Climate, 24(19), 5061-5080. doi: 10.1175/2011JCLI4052.1
721	Hegg, D. A., Covert, D. S., Jonsson, H. H., & Woods, R. (2009). Differentiating
722	natural and anthropogenic cloud condensation nuclei in the california coastal
723	zone. Tellus B: Chemical and Physical Meteorology, 61(4), 669-676. doi:
724	10.1111/j.1600-0889.2009.00435.x

725	Hoppel, W. A., Frick, G. M., & Larson, R. E. (1986). Effect of nonprecipitating
726	clouds on the aerosol size distribution in the marine boundary layer. Geophysi-
727	cal Research Letters, 13(2), 125-128. doi: 10.1029/GL013i002p00125
728	Hu, Y., Rodier, S., Xu, KM., Sun, W., Huang, J., Lin, B., Josset, D. (2010). Oc-
729	currence, liquid water content, and fraction of supercooled water clouds from
730	combined caliop/iir/modis measurements. Journal of Geophysical Research:
731	Atmospheres, 115(D4).
732	Irving, D., Simmonds, I., & Keay, K. (2010). Mesoscale cyclone activity over the ice-
733	free southern ocean: 1999–2008. Journal of Climate, 23(20), 5404-5420. doi: 10
734	.1175/2010 JCLI3628.1
735	Janjic, Z. I. (2001). Nonsingular implementation of the mellor-yamada level 2.5
736	scheme in the ncep meso model (Office Note No. 437). National Center for En-
737	vironmental Prediction (NCEP).
738	Jolly, B., McDonald, A. J., Coggins, J. H. J., Zawar-Reza, P., Cassano, J., Laz-
739	zara, M., Dale, E. (2016). A validation of the Antarctic mesoscale pre-
740	diction system using self-organizing maps and high-density observations
741	from SNOWWEB. Monthly Weather Review, 144(9), 3181-3200. doi:
742	10.1175/MWR-D-15-0447.1
743	Jones, C. R., Bretherton, C. S., & Leon, D. (2011). Coupled vs. decoupled boundary
744	layers in vocals-rex. Atmospheric Chemistry and Physics, 11(14), 7143–7153.
745	Retrieved from https://acp.copernicus.org/articles/11/7143/2011/
746	doi: 10.5194/acp-11-7143-2011
747	Korhonen, H., Carslaw, K. S., Forster, P. M., Mikkonen, S., Gordon, N. D., &
748	Kokkola, H. (2010). Aerosol climate feedback due to decadal increases in
749	southern hemisphere wind speeds. Geophysical Research Letters, $37(2)$ . doi:
750	10.1029/2009GL041320
751	Kremser, S., Harvey, M., Kuma, P., Hartery, S., Saint-Macary, A., McGregor, J.,
752	Parsons, S. (2020). Southern ocean cloud and aerosol data: a compilation of
753	measurements from the 2018 southern ocean ross sea marine ecosystems and
754	environment voyage. Earth System Science Data(submitted).
755	Kuma, P., McDonald, A. J., Morgenstern, O., Alexander, S. P., Cassano, J. J., Gar-
756	rett, S., Williams, J. (2020). Evaluation of southern ocean cloud in the
757	hadgem3 general circulation model and merra-2 reanalysis using ship-based
758	observations. Atmospheric Chemistry and Physics, 20(11), 6607–6630. doi:
759	10.5194/acp-20-6607-2020
760	McCluskey, C. S., Hill, T. C. J., Humphries, R. S., Rauker, A. M., Moreau, S.,
761	Strutton, P. G., DeMott, P. J. (2018). Observations of ice nucleating
762	particles over Southern Ocean waters. $Geophysical Research Letters, 45(21),$
763	11,989-11,997. doi: 10.1029/2018GL079981
764	McCoy, D. T., Burrows, S. M., Wood, R., Grosvenor, D. P., Elliott, S. M., Ma,
765	PL., Hartmann, D. L. (2015). Natural aerosols explain seasonal and spa-
766	tial patterns of Southern Ocean cloud albedo. Science Advances, $1(6)$ . doi:
767	10.1126/sciadv.1500157
768	McErlich, C., McDonald, A., Schuddeboom, A., & Silber, I. (2021). Comparing
769	satellite- and ground-based observations of cloud occurrence over high south-
770	ern latitudes. Journal of Geophysical Research: Atmospheres, 126(6). doi:
771	https://doi.org/10.1029/2020JD033607
772	McFarquhar, G. M., Bretherton, C., Marchand, R., Protat, A., DeMott, P. J.,
773	Alexander, S. P., McDonald, A. (2020). Observations of clouds, aerosols,
774	precipitation, and surface radiation over the southern ocean: An overview of
775	capricorn, marcus, micre and socrates. Bulletin of the American Meteorological
776	Society, 1 - 92. doi: 10.1175/BAMS-D-20-0132.1
777	Modini, R. L., Frossard, A. A., Ahlm, L., Russell, L. M., Corrigan, C. E., Roberts,
778	G. C., Leaitch, W. R. (2015). Primary marine aerosol-cloud interactions off
779	the coast of California. Journal of Geophysical Research: Atmospheres, 120(9),

780	4282-4303. doi: 10.1002/2014JD022963
781	Petters, M. D., & Kreidenweis, S. M. (2007). A single parameter representation
782	of hygroscopic growth and cloud condensation nucleus activity. Atmospheric
783	Chemistry and Physics, 7(8), 1961–1971. Retrieved from http://www.atmos
784	-chem-phys.net/7/1961/2007/ doi: 10.5194/acp-7-1961-2007
785	Prather, K. A., Bertram, T. H., Grassian, V. H., Deane, G. B., Stokes, M. D., De-
786	Mott, P. J., Zhao, D. (2013). Bringing the ocean into the laboratory to
787	probe the chemical complexity of sea spray aerosol. Proceedings of the National
788	Academy of Sciences, $110(19)$ , 7550–7555. doi: 10.1073/pnas.1300262110
789	Pruppacher, H. R., Klett, J. D., & Wang, P. K. (1998). Microphysics of clouds and
790	precipitation. Taylor & Francis.
791	Quinn, P. K., Coffman, D. J., Johnson, J. E., Upchurch, L. M., & Bates, T. S.
792	(2017). Small fraction of marine cloud condensation nuclei made up of sea
793	spray aerosol. Nature Geoscience, 10, 674–679. doi: 10.1038/ngeo3003
794	Romps, D. M. (2017). Exact expression for the lifting condensation level. <i>Journal of</i>
795	the Atmospheric Sciences, 74(12), 3891-3900. doi: 10.1175/JAS-D-17-0102.1
	Russell, L. M., Lenschow, D. H., Laursen, K. K., Krummel, P. B., Siems, S. T.,
796	Bandy, A. R., Bates, T. S. (1998). Bidirectional mixing in an ace 1 marine
797	boundary layer overlain by a second turbulent layer. Journal of Geophysical
798	Research: Atmospheres, 103(D13), 16411–16432.
799	Schuddeboom, A., Varma, V., McDonald, A. J., Morgenstern, O., Harvey, M., Par-
800	sons, S., Furtado, K. (2019). Cluster-based evaluation of model compen-
801	sating errors: A case study of cloud radiative effect in the Southern Ocean.
802	Geophysical Research Letters, 46(6), 3446-3453. doi: 10.1029/2018GL081686
803	Spearman, C. (1904). The proof and measurement of association between two
804	things. The American Journal of Psychology, 15(1), 72–101. Retrieved from
805	http://www.jstor.org/stable/1412159
806	Trenberth, K. E., & Fasullo, J. T. (2010). Simulation of present-day and twenty-
807	first-century energy budgets of the Southern Oceans. Journal of Climate,
808	23(2), 440-454.
809	Truong, S. C. H., Huang, Y., Lang, F., Messmer, M., Simmonds, I., Siems, S. T., &
810	Manton, M. J. (2020). A climatology of the marine atmospheric boundary
811	layer over the southern ocean from four field campaigns during 2016–2018.
812	Journal of Geophysical Research: Atmospheres, 125(20), e2020JD033214. doi:
813	https://doi.org/10.1029/2020JD033214
814 815	Twomey, S. (1977). The influence of pollution on the shortwave albedo of
	clouds. Journal of the Atmospheric Sciences, 34(7), 1149-1152. doi:
816	10.1175/1520-0469(1977)034<1149:TIOPOT>2.0.CO;2
817	Vergara-Temprado, J., Miltenberger, A. K., Furtado, K., Grosvenor, D. P., Ship-
818	way, B. J., Hill, A. A., Carslaw, K. S. (2018). Strong control of southern
819	ocean cloud reflectivity by ice-nucleating particles. <i>Proceedings of the National</i>
820	Academy of Sciences, 115(11), 2687–2692. doi: 10.1073/pnas.1721627115
821	Wofsy, S. C. (2011). Hiaper pole-to-pole observations (hippo): fine-grained, global-
822	scale measurements of climatically important atmospheric gases and aerosols.
823	Philosophical Transactions of the Royal Society A: Mathematical, Physical and
824	Engineering Sciences, 369(1943), 2073-2086. doi: 10.1098/rsta.2010.0313
825	Yin, B., & Albrecht, B. A. (2000). Spatial variability of atmospheric boundary layer
826	structure over the eastern equatorial pacific. Journal of Climate, 13(9), 1574 -
827	1592. doi: 10.1175/1520-0442(2000)013<1574:SVOABL>2.0.CO;2
828	Young, I. R., Zieger, S., & Babanin, A. V. (2011). Global trends in wind speed and
829	wave height. Science, 332(6028), 451–455. doi: 10.1126/science.1197219
830	wave neight. Detente, 002(0020), 401-400. 001. 10.1120/SCIENCE.1197219