

The Mysterious Green Streaks Below STEVE

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Abstract

STEVE (Strong Thermal Emission Velocity Enhancement) is an optical phenomenon of the sub-auroral ionosphere arising from extreme ion drift speeds. STEVE consists of two distinct components in true-color imagery: a mauve or whitish arc extended in the magnetic east-west direction, and a region of green emission adjacent to the arc, often structured into quasi-periodic columns aligned with the geomagnetic field (the “picket fence”). This work employs high-resolution imagery by citizen scientists in a critical examination of fine scale features within the green emission region. Of particular interest are narrow “streaks” of emission forming underneath field-aligned picket fence elements in the 100–110-km altitude range. The streaks propagate in curved trajectories with dominant direction toward STEVE from the poleward side. The elongation is along the direction of motion, suggesting a drifting point-like excitation source, with the apparent elongation due to a combination of motion blur and radiative lifetime effects. The cross-sectional dimension is <1 km, and the cases observed have a duration of ~10–30 s. The uniform coloration of all STEVE green features in these events suggests a common optical spectrum dominated by the oxygen 557.7-nm emission line. The source is most likely direct excitation of ambient oxygen by superthermal electrons generated by ionospheric turbulence induced by the extreme electric fields driving STEVE. Some conjectures about causal connections with overlying field-aligned structures are presented, based on coupling of thermal and gradient-drift instabilities, with analogues to similar dynamics observed from chemical release and ionospheric heating experiments.

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The Mysterious Green Streaks Below STEVE

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Key Points:

- Extremely small point-like features are observed within the ‘picket fence’ STEVE phenomenon.
- Such features point to the role of local ionospheric production.
- The features are among the smallest optical features found in the natural airglow or aurora.

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Abstract

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1 Introduction

STEVE (Strong Thermal Emission Velocity Enhancement) is a recently identified optical feature in the sub-auroral ionosphere appearing within a narrow channel of extreme westward ion drifts (MacDonald et al., 2018). The phenomenon was identified by citizen scientists using consumer camera equipment. In true-color photography STEVE appears as a diffuse arc extended in the East-West direction with color ranging from mauve to gray-white, which is often, but not always, accompanied by ephemeral green features nicknamed “the picket fence”. Multi-point triangulation has placed the mauve component at an altitude range of 130–270 km, with picket fence features extending below to as low as ~ 95 km (Archer et al., 2019). Conjugate measurements by the Swarm satellites at ~ 400 km altitude have detected B_{\perp} ion velocities approaching 6 km/s in the STEVE

48 channel, with evidence for electron temperatures approaching ~ 1 eV near the edges (Archer
49 et al., 2019). Initial spectroscopy of STEVE has revealed the mauve color to arise from
50 the oxygen 630-nm red line emission superimposed on a continuum spectrum from ~ 400
51 to ~ 700 nm (Gillies et al., 2019). The whitish color is attributed to lower altitude events,
52 where the metastable $O(^1D)$ state responsible for the red line component is collisionally
53 quenched (Liang et al., 2019). The companion picket fence region has been found to be
54 predominantly oxygen 557.7-nm green line, with a trace contribution from N_2 first pos-
55 itive emissions (Mende, Harding, & Turner, 2019). The lack of emissions from higher en-
56 ergy N_2^+ transitions has argued against precipitating magnetospheric electrons as a source
57 of the green features (Mende et al., 2019), although the role of precipitation in STEVE
58 remains a subject of debate (Nishimura et al., 2019).

59 This paper considers STEVE’s green companion from a morphological point of view.
60 Particular attention is placed on ephemeral “streaks” of green emission observed below,
61 and conjugate to, field-aligned structures comprising the picket fence. Several examples
62 of this feature are shown in Figure 1, as captured in true-color imagery by citizen sci-
63 entists. In all cases we have examined thus far, the streaks form on the poleward side
64 of STEVE, exhibit an approximate north-south elongation, and propagate equatorward
65 toward the main STEVE channel. Similar streaks to those shown in Figure 1 are also
66 found embedded in more dynamic displays (e.g., Figure 5), where their motion is some-
67 what more complicated and their relationship to the magnetic field topology is less clear.
68 Notably, the lifetime of the streaks is sufficiently long (~ 10 s) to enable tracking through
69 multiple images, allowing for deconvolution of finite exposure-time effects (motion blur)
70 and spatiotemporal correlation with other features in the field.

71 In the remainder of this work, we provide an initial examination of the altitude,
72 trajectory, orientation, and dimensions of this feature, followed by some conjectures about
73 its origin, and connections to broader questions of the picket fence source and the modes
74 of energy dissipation represented by the STEVE phenomenon. This work also suggests
75 new opportunities for the use of photometric imaging as a diagnostic of ionospheric tur-
76 bulence under extreme conditions.

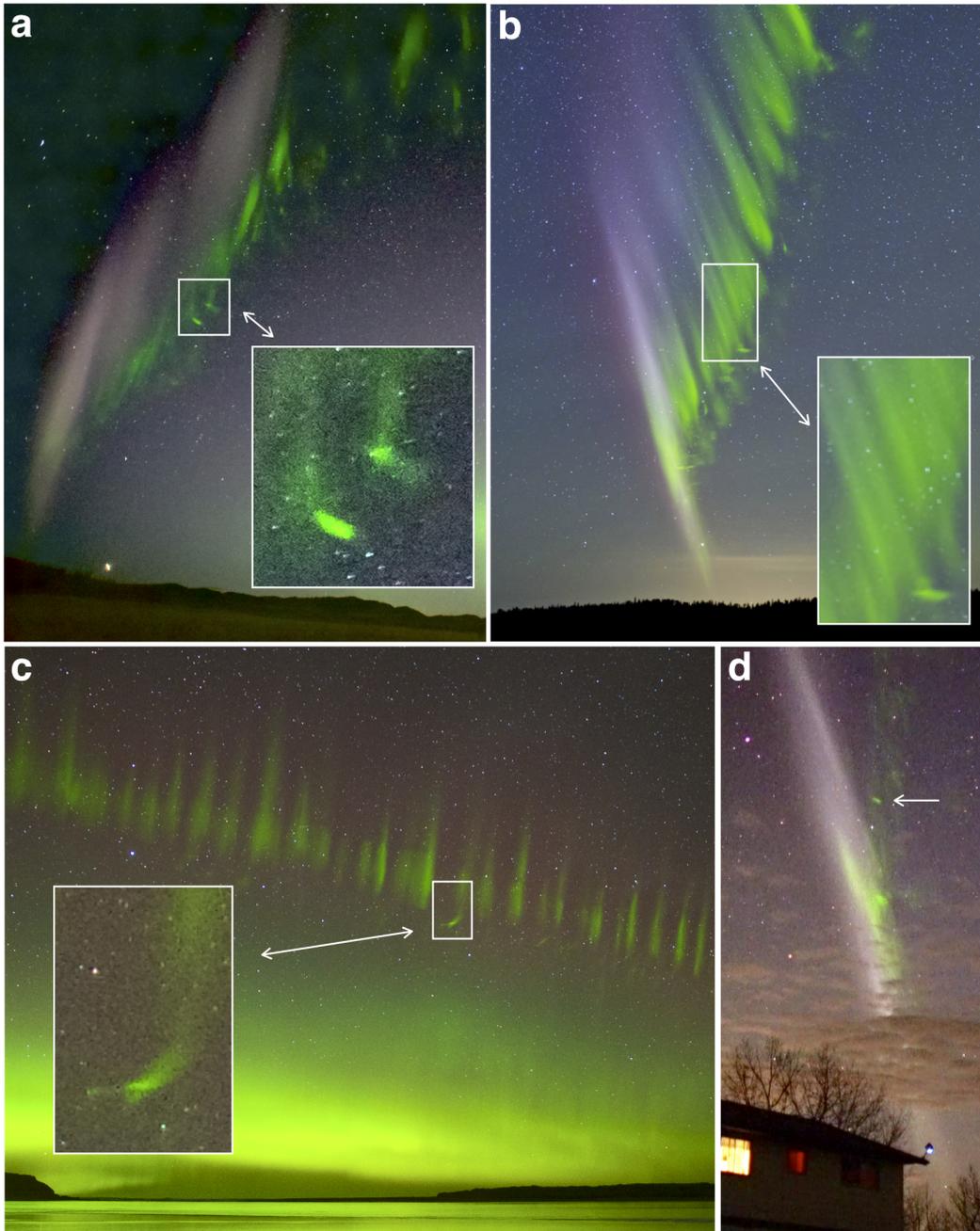


Figure 1. Sample images highlighting STEVE's green component and the mysterious green streaks appearing below the "picket fence" a) 6 May 2018, 11:21:25 LT, 4-s exposure, 51.255°N, 114.701°W (credit: Alexei Chernenkov). b) 13 September 2018, Isle Royale National Park (credit: Shawn Malone). c) 28 March 2017, 15-s exposure (credit: Stephen Voss). d) 6 May 2018, ~11:19 LT, 33-ms exposure (NTSC video) (credit: Alan Dyer).

2 Analysis results

The experimental results in this work were all derived from imagery recorded by citizen scientists using commercially available equipment. The unfamiliar nature of this phenomenon and the unusual views obtained by the photographers make interpretation with respect to physical hypotheses challenging. Observer perspective, camera exposure times, magnetic field topology, and radiative lifetime effects must all be considered carefully in drawing physical conclusions from these observations. We endeavor to discuss these issues in the context of our analysis. The features of interest are often faint, and the images displayed in this work have been adjusted to enhance contrast. Our findings do not rely on absolute photometric calibration, although that will undoubtedly become an important consideration in future studies.

2.1 Altitude determination

Knowledge of the height of the features provides an important constraint for assessing physical mechanisms. Direct triangulation on the basis of time synchronized images from two locations allows the calculation of the height of aurora, STEVE or airglow structures. Archer et al. (2019) and Palmroth et al. (2019) present algorithms for this purpose in their works. In standard optical triangulation, the position of an object (latitude, longitude, height) is calculated from observer locations and object bearing using standard trigonometric relations. The bearing (elevation and azimuth) of a point in the night sky may be determined with high precision via analysis of the background star field (Lang, Hogg, Mierle, Blanton, & Roweis, 2010). Limiting the analysis to a neighborhood of stars surrounding the features minimizes potential image distortion effects.

Figures 2a and 2b show the images used in our triangulation analysis, recorded from locations (51.255°N , 114.701°W), and (51.267°N , 114.328°W), respectively (about 50 km WNW of Calgary, Alberta). For the current study, we focus our analysis on the natural triangulation points defined by the leading edge and trailing edge of the streak, labeled R1 and R2. The triangulation geometry is shown to scale in Figure 2c. This triangulation problem exemplifies the opportunities and challenges in using citizen science imagery for precision analyses. The method developed for this study is described in detail in a supporting document, and presented in overview here.

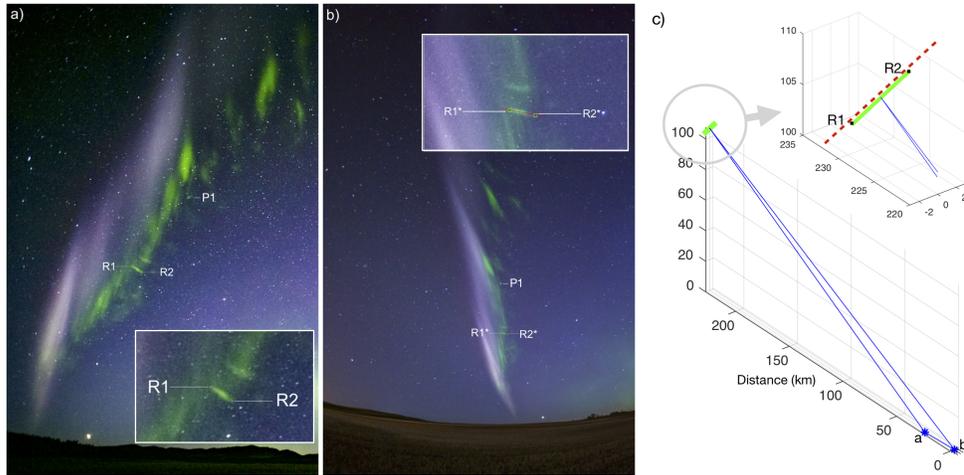


Figure 2. a) Image from the location (51.255°N , 114.701°W) during the same event as Figure 1a. b) Time-synchronized image of the same features recorded from location (51.267°N , 114.328°W) (courtesy Jun Wang). This pair of images was used for triangulation of the low altitude streak with the end points R1 and R2. The altitudes were determined to be 104.6 ± 1.6 km for R1, 106.4 ± 1.6 km for R2.

107 The images were recorded near 23:20:14 MDT, but the camera time information
 108 is unreliable. By analyzing the positions of prominent landmarks on the horizon rela-
 109 tive to reference stars in the background, we were able to reconstruct the observation
 110 time to within ± 2 minutes. The star field is a slowly varying background, and this level
 111 of uncertainty translates to a height uncertainty of ~ 1 km for objects in the lower iono-
 112 sphere such as STEVE.

113 The more significant source of error is the relative time offset between the images.
114 The fine-scale features were changing rapidly in these images. The triangulation result
115 may thus be highly sensitive to small timing errors between images. This is especially
116 true for this geometry (Figure 2c), where the observer separation was small (25.9 km)
117 relative to the target distance (~ 230 km). Fortunately, the observer at the location of
118 Figure 2a captured a contiguous time sequence of images of this event. The sequence is
119 shown in Figure 3, highlighting the two prominent streaks in Figure 2. Comparing fine-
120 scale features in the image sequence with features in single frame observation of Figure 2b
121 yielded a near perfect match to the frame displayed in Figure 2a. We conclude that the
122 time offset between the images is < 2 s.

123 Image synchronization is further improved through the use of a modified triangu-
124 lation method that exploits the unique one-dimensional geometry of these features. The
125 streaks in Figure 3 are seen to move along a line corresponding to their direction of elon-
126 gation (the significance of this is further discussed in Section 2.2). We therefore assume
127 the streaks were moving in a straight line through three-dimensional space during the
128 exposure, and that the images represent two-dimensional projections of this line (i.e.,
129 an affine transformation). We next let R1 and R2 in Figure 2a be fixed reference points.
130 The projections of these points in Figure 2b are unknown, but they are assumed to lie
131 along a line defined by the streak (depicted in red in Figures 2b (inset) and Figure 2c).
132 The geophysical coordinates of points R1 and R2 were found by testing points along the
133 red line for consistency with both projections. Specifically, for each test point the po-
134 sition of the ground intersection was computed from the elevation and azimuth informa-
135 tion, and the corresponding altitude was computed independently for each observer. The
136 solution is the position at which the altitudes agree. The logic of the method is that there
137 is only one line in 3D space that is consistent with a line in 2D space observed from two
138 locations.

139 The optimal positions of R1 and R2 in Figure 2b are labeled R1* and R2*. It is
140 interesting to note that R1 and R1* both lie on the leading edge of the streak in their
141 respective images, where the edge is well defined, while R2* lies a bit behind the trail-
142 ing edge, which is somewhat broader and more diffuse. For R1 the optimized mean height
143 is $104.6 \text{ m} \pm 1.6 \text{ km}$, and for R2 it is $106.4 \text{ km} \pm 1.6 \text{ km}$. Uncertainties were estimated us-
144 ing a Monte Carlo method. Histograms of triangulated heights were constructed based
145 on 10,000 trials distributed uniformly over the range of pointing uncertainties defined

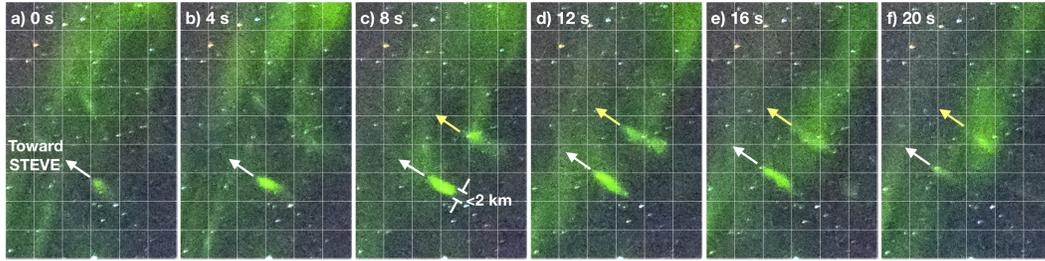


Figure 3. Sequence of 4-second image samples around the time of Figure 1a, showing the general propagation of the streaks toward the main STEVE arc, and their general elongation in the direction of propagation.

146 by uncertainties in timing, feature dimensions, and observer location (minor). The un-
 147 certainty ranges reported above are the interdecile ranges of the resulting histograms.
 148 Further detail is presented in the supplemental document.

149 We conclude that the streak R1–R2 resides in the lower ionospheric E-region in the
 150 altitude range of 100–108 km, with some evidence for a slight downward orientation of
 151 the streak toward the main STEVE arc. This is consistent with observations of trajec-
 152 tories in other events, as discussed in Section 2.2 and 3.3.

153 2.2 Trajectory

154 Figure 3 shows a sequence of cropped images at 4-second cadence, documenting
 155 the formation and evolution of the two streaks highlighted in Figure 1a and 2a. The streaks
 156 initially appeared as point-like features (white arrow in panel a, yellow arrow in panel
 157 c), which subsequently elongate along their direction of motion. The trajectory has a dom-
 158 inant component toward the the main STEVE channel in these observations. The elon-
 159 gation of the features is influenced by at least three effects. The first is simple motion
 160 blur caused by the 4-s exposure. The second is emission afterglow. As suggested in prior
 161 work (Gillies et al., 2019; Mende et al., 2019), the green color is predominantly due to
 162 the 557.7 nm line, produced by the metastable $O(^1S-^3P)$ transmission of atomic oxy-
 163 gen, with radiative lifetime 0.74 s. A moving source of $O(^1S)$ will produce a luminous
 164 tail in the 557.7-nm emission due to the finite radiative lifetime. The third effect is spa-
 165 tiotemporal variability in the excitation source itself.

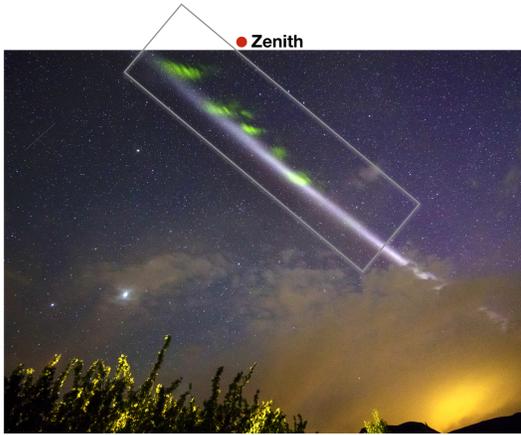


Figure 4. STEVE observed with 15-second exposure on 00:33:22 LT on 20 March 2017 from location 51.66 N, 112.91 W. The rectangle is the region detailed at 3.5-second exposure in the image sequence of Figure 5 (courtesy Neil Zeller Photography).

166 The streaks appear with with a wide range of contrast in the image sequences ex-
 167 amined thus far (e.g., Figure 1a is a 4-s exposure, while Figure 2d is a 40-ms exposure).
 168 From the evidence acquired thus far, it is difficult to assign differences in apparent streak
 169 brightness to equipment disparities versus differences in source brightness. In order to
 170 better quantify the effects described above, we turn to an event that was far more dy-
 171 namic while also exhibiting greater dynamic range in the image sensor.

172 Figure 4 was recorded on 20 May 2017 at 00:32 Local Time (LT) from location 51.66
 173 N, 112.91 W at 15-second exposure. In this display, a series of coherent green structures
 174 are observed to extend away from the main STEVE channel. The two structures toward
 175 the top of the image are composed of periodically spaced bands, each aligned approx-
 176 imately parallel to the main STEVE channel. The local vertical zenith direction lies just
 177 outside the field of view, as indicated by the red circle. The mauve-white arc is stretched
 178 out along the magnetic east-west direction. The orientation of the green features in Fig-
 179 ure 4 is difficult to establish in this projection. Accounting for the geometric point of view
 180 is particularly important in the under-studied sub-auroral region, where the field lines
 181 are not vertical, and the features are not easily reconciled with intuition developed from
 182 auroral observations.

183 This event was captured by a second co-located camera with narrower field of view
 184 and higher image cadence (3.5-second). Figure 5 shows four selected images from this

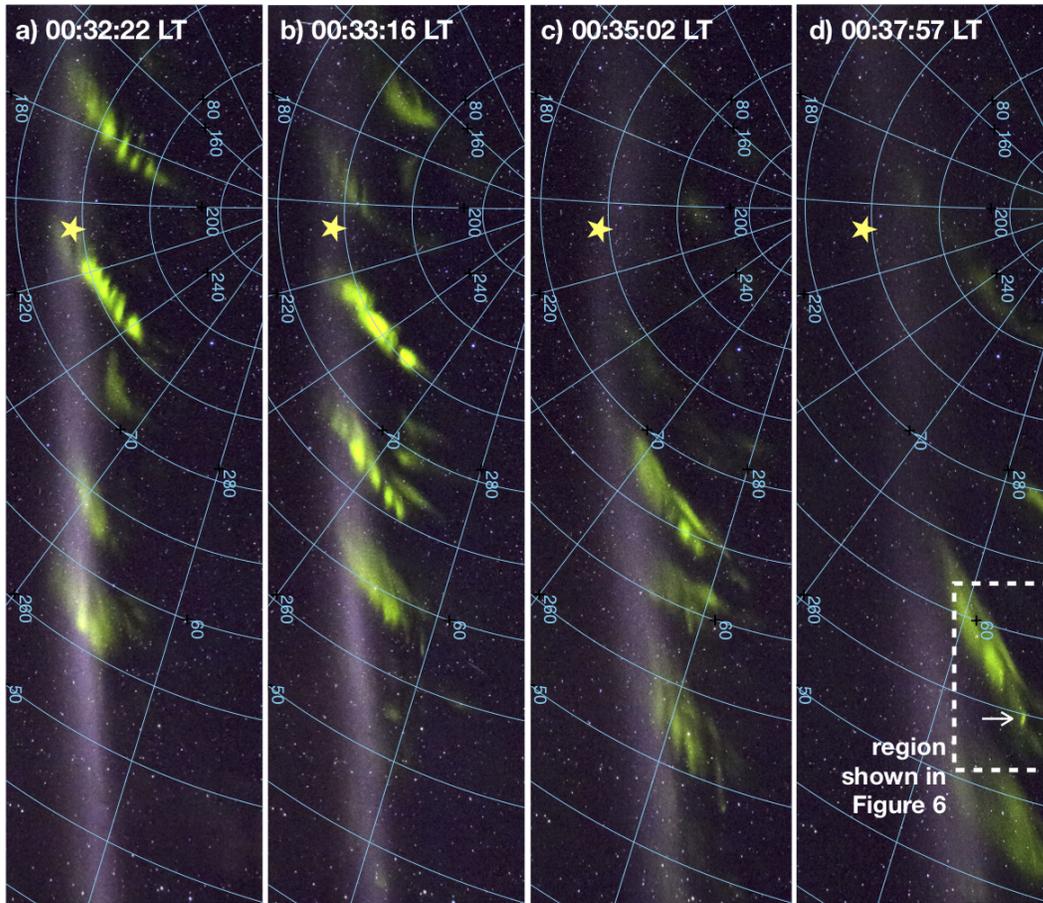


Figure 5. Sample images of dynamic green emissions observed within the rectangular region of Figure 4. Azimuth and elevation contours are shown in blue. The magnetic zenith direction is indicated by the yellow star.

185 camera, with local time as indicated. The field-of-view corresponds to the rectangular
 186 region in Figure 4. Blue contours indicate geographic azimuth and elevation as deter-
 187 mined from star field fitting using the Astrometry.net package (Lang et al., 2010). The
 188 yellow star in each frame indicates the magnetic zenith direction (inclination 73.5° , dec-
 189 lination 14°), calculated using the International Geomagnetic Reference Field (IGRF)
 190 model (Thébault et al., 2015))

191 The image samples in Figure 5 were selected to give a sense of how these features
 192 varied as they moved westward from near zenith (panel a) to lower elevation (panel d).
 193 The individual features changed substantially from frame to frame, indicating that this
 194 phenomenon was not fully resolved at 3.5-seconds. Dynamic features in the high-latitude

195 aurora exhibit orientations and motion that are clearly organized with respect to the mag-
196 netic field, with the local magnetic zenith serving as the point of convergence (e.g., Dahlgren,
197 Semeter, Marshall, & Zettergren, 2013). The features in Figure 5 exhibit no definitive
198 or stable orientation with respect to the magnetic field. The reader is encouraged to view
199 the full image sequence included as supporting information in time-lapse format (movie
200 MS3).

201 One feature that remained coherent across multiple frames is the small streak within
202 the dashed box of Figure 5d. This streak has characteristics similar to the streaks in Fig-
203 ures 1 and 3 – i.e., it is the smallest object within the field, it appears below the other
204 features (i.e., lower elevation), and, unlike other features in this sequence, it persisted
205 as a coherent drifting object across several frames.

206 Figure 6a-e shows the evolution of this feature through five consecutive 3.5-second
207 frames. The field-of-view corresponds to the dashed box in Figure 5d. Figure 6f dupli-
208 cates the image of Figure 6e with fiducial marks inserted: the white arrows show the lo-
209 cation, length, and direction-of-motion of the streak as extracted from each of panels a-
210 e. The streak is seen to be contiguous from frame to frame (i.e., tip of one arrow lines
211 up with tail of the next). This suggests that it is produced by a drifting point-like source,
212 and that the observed elongation is primarily caused by motion blur and afterglow ef-
213 fects previously discussed. The trajectory is also seen to be slightly curved in this per-
214 spective. The curvature is consistent with bending toward the main STEVE channel. This
215 trajectory may have similarities to Figure 1c, where the trace emission behind the streak
216 suggests a drift path that bent into the horizontal plane. The red line inserted in panel
217 f indicates the magnetic field-aligned direction. Its significance will be discussed in Sec-
218 tion 2.3.

219 **2.3 Dimensions and Velocity**

The high fidelity of Figure 6 allows for a quantitative examination of dimensions and velocities of the streak. Figure 7 shows the relative brightness of the camera’s green channel versus distance along its trajectory for each panel (a–e) of Figure 6. The distance scale was computed in the following manner. First, the plate scale, $p = 0.0257$ radians/pixel, was determined for the region of interest (Figure 5d, inset) using the star field calibration. Next we select as the origin a pixel (x_0, y_0) corresponding to the tail

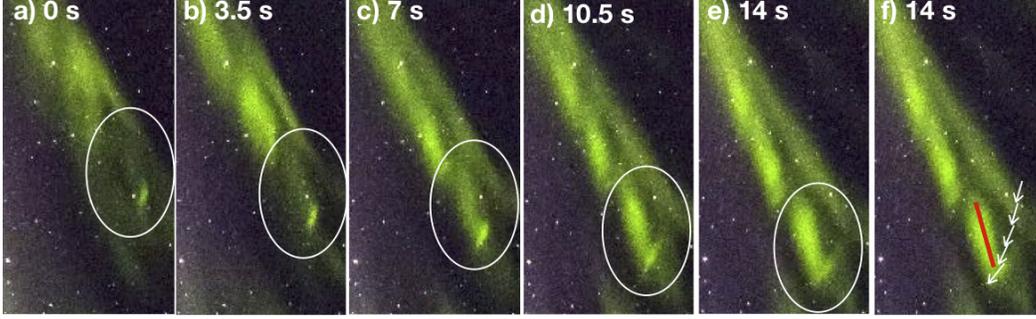


Figure 6. Image sequence at 3.5-second cadence corresponding to the rectangular region in Figure 5d, showing the evolution of an emission streak that persisted in 5 contiguous frames (feature within the white oval). In panel f, the length, orientation, and direction of motion of the streak is shown as a series of vectors. The red line shows the magnetic field-aligned direction superimposed on a features that formed above the streak.

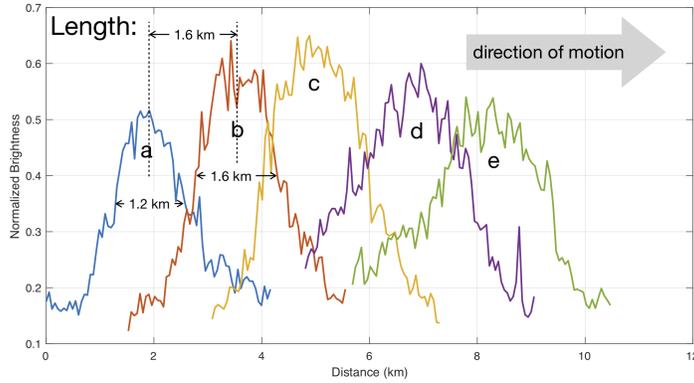


Figure 7. Normalized brightness as a function of distance along the trajectory of the streak feature in Figure 6a-e. The behavior is consistent with a drifting point source. An asymmetry develops as the object moves, consistent with an “afterglow” tail due to the 0.74-s radiative lifetime of the $O(^1S)$ state.

of the first arrow in Figure 6f. The selection of this point is somewhat arbitrary, as it is the relative motion between frames that is of interest. The range to this point is given by $R_0 = z_0 / \sin(\theta_0)$, where θ_0 is the elevation and z_0 is an assumed feature altitude, taken to be 100 km. The pixel coordinates (x_p, y_p) of a cut through the feature are then converted to physical distance d using the small angle formula,

$$d = R_0 p \sqrt{(x_p - x_0)^2 + (y_p - y_0)^2} \quad (\text{km}). \quad (1)$$

220 The streak trajectory through three-dimensional space is unknown. The distance
 221 scale so derived corresponds to a projection of the actual distance scale into the image
 222 plane under the stated assumptions. If we assume a point-like object that elongates in
 223 the direction of motion, this scale approximates the projected distance in km along the
 224 trajectory traced by the white arrows in Figure 6f. At its initial appearance, the full-
 225 width-at-half-maximum (FWHM) is similar to the separation between peaks. This re-
 226 sult is consistent with a moving point source subject to motion blur. The streak can also
 227 be seen to broaden and develop an asymmetric “tail” behind its trajectory. This is qual-
 228 itatively consistent with the afterglow effect due to the finite radiative lifetime of the $O(^1S)$
 229 state. If the actual trajectory has a component orthogonal to the image plane, the widths
 230 and peak positions extracted from Figure 7 would be compressed by a common scale fac-
 231 tor.

232 Figure 8 shows the evolution of this feature in the direction *transverse* to its prop-
 233 agation. The curves have been manually shifted to align the lower altitude edge of the
 234 streaks, in order to compare relative changes in width during its lifetime. When first ob-
 235 served (curve a) the streak has a cross-sectional width of ~ 350 m. This streak is thus
 236 among the smallest optical aeronomical features observed at any latitude in the aurora
 237 or airglow. As it evolves through subsequent frames (curves b–e), an extended region of
 238 emission is seen to develop to the left (i.e., at higher elevation angle). Some context for
 239 this can be obtained by returning to Figure 6. The new region of emission corresponds
 240 to a developing magnetic field-aligned feature in the upper part of the encircled region,
 241 above the streak. In Figure 6f, a red bar has been inserted to show the magnetic field-
 242 line direction projected into the image plane. Magnetic conjugacy of low-altitude streaks
 243 and field-aligned features is also observed in the wide-field image samples of Figure 1.
 244 The analysis of Figure 8 provides possible evidence for the contemporaneous develop-
 245 ment of the low altitude streaks of emission and magnetic field-aligned features above
 246 it.

247 **3 Discussion**

248 True-color images of selected STEVE events obtained by citizen scientists have been
 249 used in a critical examination of small-scale features in the green “picket fence” region.
 250 Image sequences acquired at 3.5- and 4-second cadence have revealed dynamic sub-kilometer
 251 features with varying orientations, dimensions, and motions (Figures 5 and 6). Unlike

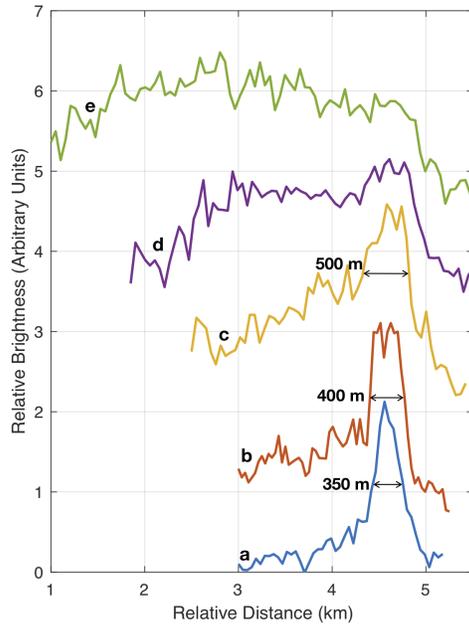


Figure 8. Brightness cuts extracted through the highlighted feature in panels a-e of Figure 6 in a direction aligned with the magnetic field, showing as a line plot the apparent field-aligned elongation of the structure during its lifetime.

252 the quasi-periodically spaced green columns that inspired the “picket fence” designation,
 253 these features are not extended along magnetic field lines, and are thus inconsistent with
 254 production via energetic particle precipitation. Readers are encouraged to view the three
 255 time-lapse videos included as supplemental material in order to develop their own im-
 256 pression of these unusual features, and the unique perspectives obtained by citizen sci-
 257 entists.

258 This work has focused on a particular repeatable feature: a narrow “streak” of emis-
 259 sion appearing below the picket fence that propagates toward STEVE from the poleward
 260 side. This feature is noteworthy for several reasons: 1) it is the lowest-altitude and smallest-
 261 scale optical feature associated with STEVE, 2) it has been observed in many STEVE
 262 events (Figure 1), 3) in events where contiguous image sequences are available, the fea-
 263 ture has been observed to persist for >10 s as a coherent propagating object (Figure 3,
 264 Figure 6, and supporting videos), 4) it is magnetically conjugate to, and sometimes op-
 265 tically connected with, overlying field-aligned structures. The following sections present
 266 some conjectures based on the initial analysis reported herein.

267 **3.1 Source of green line excitation**

268 The periodic spacing and magnetic field elongation often observed in the green fea-
 269 tures adjacent to STEVE (e.g., Figure 1) have naturally led many to assume production
 270 via usual auroral mechanisms – i.e., penetration of magnetospheric electrons with kinetic
 271 energy >1 keV (e.g., Gillies et al., 2019; Mishin & Streltsov, 2019; Nishimura et al., 2019).
 272 The initial spectroscopic measurements acquired by Gillies et al. (2019) are irreconcil-
 273 able with this hypothesis. A careful analysis by (Mende et al., 2019) found the spectrum
 274 to be dominated by the metastable oxygen 557.7-nm line (4.19 eV excitation energy, 0.74-
 275 s radiative lifetime) but with a trace contribution from prompt N_2 first positive (1P) emis-
 276 sions (7.35 eV excitation energy). Entirely absent, however, were contributions from higher
 277 energy emissions of N_2^+ , often represented in auroral studies by the band-head of the the
 278 N_2^+ first negative (1N) group at 427.8 nm (18.75 eV excitation energy). This emission,
 279 produced by collisional ionization and excitation of ambient N_2 , must be present for par-
 280 ticle penetration to these altitudes. The presence of N_2 1P without N_2^+ 1N has argued
 281 for a lack of primary electrons with the requisite >1 -keV energy range, rather than a de-
 282 pletion of ambient N_2 (Mende et al., 2019). This finding supported earlier conjectures
 283 based on color comparisons (Mende et al., 2019) that the source of the green compan-
 284 ion to STEVE is likely direct excitation of oxygen $O(^1S)$ by superthermal electrons en-
 285 ergized locally in the ionosphere.

286 For the features examined in this work, the particle precipitation hypothesis is ex-
 287 cluded based on more direct morphological considerations. The small scales and varie-
 288 gated orientations of the lower ionospheric features highlighted in Figures 1–6 cannot
 289 be accounted for by energetic electron penetration. Analysis of the streak in Figures 6
 290 and 8 are consistent with a drifting point-like source, with cross-sectional size as small
 291 as ~ 350 meters. Triangulation of a similar feature in Figure 2 has placed this source in
 292 the lower ionospheric E -region in the 100–110 km range, and below other picket fence
 293 features in the field.

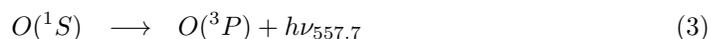
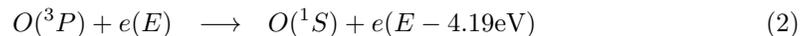
294 **3.2 Superthermal electron production**

295 In seeking a local source of free energy able to excite oxygen green line features at
 296 sub-kilometer scales, it should be noted that the ion drift speeds within STEVE exceed
 297 6 km/s (Archer et al., 2019; MacDonald et al., 2018). The patterns and intensities of tur-

298 bulent heating caused by such supersonic plasma jets in the outer atmosphere are not
 299 well known. Additional complexity arises from the entanglement of chemistry and elec-
 300 trodynamics: ion velocities in this range are known to be associated with rapid conversion
 301 from atomic (O^+) to molecular (NO^+) ions (Anderson, Heelis, & Hanson, 1991), which
 302 would impact momentum balance in the channel in a highly nonlinear manner. Under
 303 these conditions it is not surprising to find pockets of extreme electron heating.

304 Some evidence for this has been found in conjugate satellite measurements. Nishimura
 305 et al. (2019) and Archer et al. (2019) have reported measurements from the SWARM satel-
 306 lites of a narrow channel of electron heating conjugate to STEVE, with T_e exceeding 8,000
 307 K, as well as a single-point measurement approaching 12,000 K (>1 eV). These measure-
 308 ments occurred in a region of depleted plasma densities and large upward ion velocities,
 309 consistent with expected signatures of low-altitude heating. The measured electrons are
 310 certainly non-thermal, but a distribution with average energy ~ 1 -eV would be expected
 311 to include a significant population at the requisite 4.19 eV energy for green-line excita-
 312 tion. Candidate mechanisms for superthermal electron production in the low-altitude
 313 ionosphere include the modified two-stream (Farley-Buneman) instability (Farley, 1963;
 314 Oppenheim & Dimant, 2013), and the electron and ion thermal instabilities (Dimant &
 315 Sudan, 1997). These instabilities have the lowest threshold in the 100 to 120 km alti-
 316 tude range (Dimant & Oppenheim, 2004), which is consistent with the altitude range
 317 found via triangulation in Section 2.1. The kilometer cross-sectional dimensions found
 318 in Section 2.3 are consistent with scale sizes of irregularities produced in simulations (Op-
 319 penheim & Dimant, 2013). Electron temperatures as high as 6000 K have been observed
 320 in association with plasma heating by Farley-Buneman waves (Bahcivan & Cosgrove, 2010).
 321 Using the linear relationship of (Foster & Erickson, 2000), this would correspond to an
 322 electric field of ~ 300 mV/m which, in turn, corresponds to ion drifts of ~ 6 km/s. It is
 323 quite likely that the STEVE parameter set is even more extreme.

324 The manifestation of these turbulent processes in airglow or auroral signatures has
 325 not been fully considered. The hypothesized photochemical model for the production of
 326 the 557.7-nm emission is



327 (Itikawa & Ichimura, 1990). (We note that the $O(^1D)$ state, responsible for the oxygen
 328 630-nm redline, is excited through the same collisional reaction but at lower energy (1.96 eV).
 329 However, it is quenched at lower ionospheric altitudes due to its long (~ 120 -s) radiative
 330 lifetime). Extracting spatiotemporal information about electron heating from images of
 331 this emission requires careful consideration of source dynamics and radiative lifetime ef-
 332 fects, as represented in the space-time perspective of the sensor. It must also be borne
 333 in mind that the images only provide information about electrons with energy > 4.19 eV
 334 (45,000 K). Lower energy superthermal populations are important and presumably present,
 335 but invisible in green-line imagery. A full treatment of these effects is beyond the scope
 336 of this work. But a qualitative appreciation of these effects is useful in the evaluation
 337 of evidence from citizen science imagery.

338 3.3 Relation to magnetic field-aligned features

339 An initial examination of image sequences has provided evidence for connections
 340 between the low altitude emission streaks, and magnetic field-aligned features compris-
 341 ing the “picket fence.” The evidence is summarized in Figure 9 for the examples of this
 342 study. In each panel, the grayscale image depicts the green channel of the camera, dis-
 343 played as a negative for ease of annotation. The red arrows highlight features aligned
 344 with the magnetic field; the blue arrows indicate the orientation and trajectory of the
 345 streaks. Figure 9a corresponds to Figure 3c. At this point in time, the upper streak has
 346 developed a visible tail extending behind the trailing edge, which is qualitatively con-
 347 sistent with the afterglow effects discussed in Section 2.2. This streak also exhibits a faint
 348 emission column extending above it in the magnetic field-aligned direction. The com-
 349 bination of these effects form an intriguing “L” shape in the image.

350 Figure 9b is from Figure 1b. Multiple horizontal streaks can be seen conjugate to
 351 field-aligned aurora-like structures. Panels c and d provide some supporting evidence for
 352 this conjecture. Panel c (from Figure 1c) shows evidence of a curved tail behind the fea-
 353 ture, which bends from the field-parallel direction into the horizontal direction at lower
 354 altitudes. If the green line is excited by a drifting source of hot electrons, then the drift
 355 would be confined to the B_{\parallel} direction at higher altitudes, but develop an increasing B_{\perp}
 356 component at lower altitudes, where increasing electron-neutral collisions lead to increas-
 357 ing cross-field mobility, allowing electrons to respond directly to the poleward directed
 358 electric field of the sub-auroral ion drift (SAID) channel.

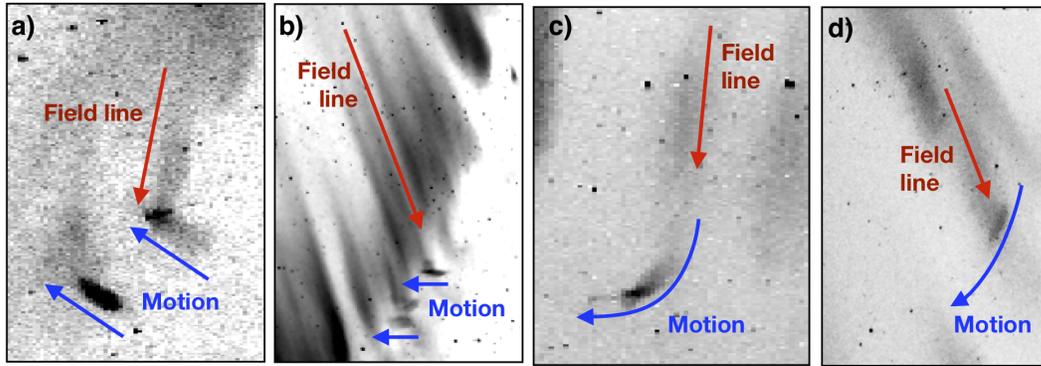


Figure 9. Relationship between magnetic field-aligned features (red), streak orientation, and streak motion (blue) for examples used in this study. The images show the green channel of the cameras, displayed as a negative (emissions are dark) a) From Figures 1a and 3c. Blue arrows show the horizontal equatorward direction of motion, as observed in Figures 3a-f. b) From Figure 1b. Blue arrows show the orientation of the streaks in relation to magnetic field direction. c) From Figure 1c. The blue curve highlights the curved shape of the airglow feature which appears to bend into the field-perpendicular equatorward direction. d) From Figure 5g (rotated 90°). Blue curve shows the complete trajectory as extracted from the image sequence in Figures 5e-j.

359 This interpretation is speculative based on a single still image. However, a simi-
 360 larly curved trajectory was observed in Figure 9d. The blue curve corresponds to the col-
 361 lection of arrows in Figure 6f, which track the trajectory through Figures 6a-e. The tem-
 362 poral development of a magnetic field-aligned feature above the streak was similarly ob-
 363 served in these frames. This process was summarized in Figure 8, where the cross-sectional
 364 width of the streak began with a value of ~ 350 m (curve a). As the streak drifted, the
 365 emission was observed to extend upward in the magnetic field-aligned direction.

366 Our triangulation analysis has placed the altitude of the streak in the lower iono-
 367 spheric *E*-region (100–110 km), where neutral collisions become an important consid-
 368 eration in plasma transport. The aggregate evidence in Figure 9 suggest a propagation
 369 path that bends into the horizontal plane as altitude decreases. This observation is qual-
 370 itatively consistent with expected effects in this altitude range.

3.4 Formation of field-aligned features from point sources

One question that arises naturally from the evidence assembled thus far is whether field-aligned “aurora-like” optical features can evolve from point-like sources created through turbulent heating. Preferential expansion of an isolated plasma population along the magnetic field direction is expected considering the difference in field-parallel versus field-perpendicular mobility (Rishbeth & Garriott, 1969). A possible structuring mechanism lies in the coupling between Farley-Buneman and gradient-drift instabilities, which arise from a common dispersion relation (Fejer & Kelley, 1980). The coupling effects can be amplified in regions of extreme density gradients (Haldoupis, Ogawa, Schlegel, Koehler, & Ono, 2005), such as those found on the poleward edge of an extreme SAID channel.

The evolution of cloud-like features into aurora-like features has been well documented in artificial plasma release experiments (Haerendel & Lüst, 1968), where a cloud of ionized Barium (Ba^+) is observed to rapidly striate across magnetic field lines and elongate along magnetic field lines, forming structures reminiscent of rayed aurora within a few seconds (Simons, Pongratz, & Gary, 1980). The cross-field striation is thought to be caused by the gradient drift instability (Linson & Workman, 1970; Simons et al., 1980). This mechanism is also plausible here due to the presence of extreme electric fields and extreme density gradients on the edges of the STEVE channel (Nishimura et al., 2019).

In chemical release experiments, the injected plasma was a cold long-lived Barium ion (Ba^+) illuminated by sunlight fluorescence. In the present situation, the illuminating agent would be superthermal electrons with energy >4.19 eV exciting the oxygen $O(^1S)$ state. In the image sequences of Figures 3 and 6, the excitation source persisted for >10 s, which is long compared with time scales for field-line elongation observed in Barium releases (Simons et al., 1980).

However, based on the limited evidence thus far obtained, it is also possible that the streaks and the picket fence are not causally connected but, rather, represent two distinct responses to the same free energy source. Mende et al. (2019) conjectured that the magnetic field-aligned features could be produced by wave heating induced by the extreme electric fields in the SAID region (Streltsov & Mishin, 2003). Fine-scale green emissions at lower conjugate altitudes could be excited by these same extreme fields.

4 Conclusions

Thus far, nothing about the optical phenomenology associated with STEVE is conveniently explained in terms of our historical understanding of auroras and airglow (Gillies et al., 2019; Harding, Mende, Triplett, & Wu, 2020; Mende et al., 2019). The reason for this may be generally traced to the extreme nature of the driving electric fields and attendant supersonic ion drifts. This regime of ion-neutral interaction has not been systematically treated in aeronomical models. The mauve-white emission of the main STEVE channel is likely related to the continuum nightglow emission that has been known for decades Bates (1993). Our understanding of its origin remains incomplete, but STEVE has attracted renewed attention to the topic (Harding et al., 2020). The green features accompanying STEVE often exhibit elongation in the magnetic field direction, suggesting they are produced by precipitation of magnetospheric electrons (Nishimura et al., 2010). But initial spectroscopic measurements are inconsistent with this hypothesis (Mende et al., 2019). The green phenomena accompanying STEVE also include a variety of blobs, streaks, and curved bands of emission (Figure 5), that are irreconcilable with the precipitation hypothesis on morphological grounds alone.

This work has focused on a particular repeatable feature of the STEVE green-line phenomenology: a narrow streak of emission observed below, and connected with, the field-aligned features comprising the picket fence. An initial analysis has been presented using time-lapse imagery and multi-station observation, revealing the following characteristics.

1. Formation is poleward of STEVE and in the lower ionosphere (100–110 km).
2. Movement follows a curved trajectory, with dominant component toward the main STEVE arc.
3. Elongation is along the direction of motion, suggestive of an unresolved moving point source.
4. Lifetime of an individual streak is ~ 10 –30 s,
5. Location is conjugate to, and sometimes optically connected with, field-aligned features above it.
6. Cross-sectional dimension is < 1 km and thus irreconcilable with production via magnetospheric precipitation.

432 The origin and significance of these unusual features is currently subject to spec-
433 ulation. The similarity in coloration with, and magnetic conjugacy to, other overlying
434 green features suggests a common optical spectrum dominated by the oxygen 557.7-nm
435 emission. The point-like nature of the emission suggests excitation via turbulent elec-
436 tron heating, rather than variations in neutral oxygen density. Prior work on radar backscat-
437 tered from ionospheric regions impinged by extreme electric fields has implicated an in-
438 terplay among Farley-Buneman, ion-thermal, electron-thermal, and gradient-drift insta-
439 bilities in creating small-scale irregularities in density and temperature in the lower iono-
440 spheric E-region (Dimant & Oppenheim, 2011). The evidence in these studies has been
441 derived from radar backscatter and in situ plasma measurements. The extreme electric
442 fields, and attendant extreme velocities, observed within STEVE have not been fully treated
443 in theoretical or simulation studies. The question is whether under extreme conditions
444 such mechanisms could produce an electron population exceeding the 4.19-eV thresh-
445 old for green-line excitation, creating an observable optical diagnostic while also provid-
446 ing a mechanism to explain the point-source nature of these objects.

447 5 Next Steps

448 STEVE has elevated the role of citizen scientists in primary research. In their quest
449 to obtain beautiful imagery of the natural world, photographers have serendipitously dis-
450 covered a phenomenon overlooked by professional scientists. References to the optical
451 phenomena of STEVE have since been uncovered in historical literature (Hunnekuhl &
452 MacDonald, 2020), but without its current appreciation as a phenomenon distinct from
453 the typical aurora and airglow (Gallardo-Lacourt, Liang, Nishimura, & Donovan, 2018).

454 Specific gaps in our understanding of STEVE can be filled through enhanced part-
455 nerships with citizen science community. First, we cannot yet state with certainty that
456 the optical spectra of all green features are the same. Needed are collaborative obser-
457 vations from high frame-rate broad-band cameras of the type employed herein, and high
458 resolution imaging spectrographs of the type employed by (Gillies et al., 2019). Second,
459 there is a need for higher-cadence observations; the 3.5-second cadence of the image se-
460 quences used herein (figure 6) have not fully resolved STEVE phenomena. Third, care-
461 fully coordinated multi-site measurements are needed for better triangulation of orien-
462 tations and trajectories of these features. This requires that precise geolocation and tim-
463 ing information be recorded into image metadata. And fourth, continued documenta-

464 tion and sharing of observations with the global community is needed – e.g., through the
 465 Aurorasaurus project (MacDonald et al., 2015)

466 Progress will also requires new collaborations within the Heliophysics community.
 467 If the fine-scale features described herein are a consequence of extreme electric fields, then
 468 they present a new diagnostic for understanding ionospheric instabilities and turbulence.
 469 Developing this capability requires that we fill the gap between theoretical predictions
 470 of heating rates to structured excitation of the 4.19-eV $O(^1S)$ state producing the vis-
 471 ible features. This will require collaborative efforts that conjoin regional transport mod-
 472 eling (e.g., Zettergren & Semeter, 2012), kinetic plasma simulation (e.g., Oppenheim &
 473 Dimant, 2004), and aeronomical modeling of optical emissions (e.g., Solomon, 2017).

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