

Weakening of the Gulf Stream at Florida Straits over the past century inferred from coastal sea-level data

Christopher Piecuch¹

¹Woods Hole Oceanographic Institution

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Abstract

The Florida Current marks the beginning of the Gulf Stream at Florida Straits, and plays an important role in climate. Nearly continuous measurements of Florida Current transport have been made at ~27N since 1982, but these data are too short to allow an assessment of possible centennial changes. Here I reconstruct Florida Current transport during 1909-2018 using probabilistic methods and principles of ocean dynamics applied to available transport measurements and longer coastal sea-level data. The Florida Current transport very likely (probability $P=0.93$) has weakened since the 1920s, such that modern measurements made within Florida Straits since 1982 likely ($P=0.87$) portray the transport in a reduced state. The weakest decadally averaged transport during the last 110 y probably ($P=0.74$) took place sometime in the last two decades. Weakening of Florida Current transport is consistent with a hypothesized steady reduction of the deep Atlantic meridional overturning circulation during the past century.

1 Coversheet for “Weakening of the Gulf Stream at Florida
2 **Straits over the past century inferred from coastal sea-**
3 **level data”**

4 Christopher G. Piecuch^{1,†}

5 ¹*Department of Physical Oceanography, Woods Hole Oceanographic Institution, Woods Hole,*
6 *Massachusetts, USA*

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9 † cpiecuch@whoi.edu

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³ Christopher G. Piecuch¹

⁴ ¹ Woods Hole Oceanographic Institution, Woods Hole, Massachusetts, USA

⁵ **The Florida Current marks the beginning of the Gulf Stream at Florida Straits, and plays**
⁶ **an important role in climate. Nearly continuous measurements of Florida Current transport**
⁷ **have been made at $\sim 27^{\circ}\text{N}$ since 1982, but these data are too short to allow an assessment of**
⁸ **possible centennial changes. Here I reconstruct Florida Current transport during 1909–2018**
⁹ **using probabilistic methods and principles of ocean dynamics applied to available transport**
¹⁰ **measurements and longer coastal sea-level data. The Florida Current transport very likely**
¹¹ **(probability $P = 0.93$) has weakened since the 1920s, such that modern measurements made**
¹² **within Florida Straits since 1982 likely ($P = 0.87$) portray the transport in a reduced state.**
¹³ **The weakest decadally averaged transport during the last 110 y probably ($P = 0.74$) took**
¹⁴ **place sometime in the last two decades. Weakening of Florida Current transport is consistent**
¹⁵ **with a hypothesized steady reduction of the deep Atlantic meridional overturning circulation**
¹⁶ **during the past century.**

¹⁷ Swiftly flowing north through the narrow, shallow Florida Straits, the Florida Current marks
¹⁸ the headwaters of the Gulf Stream^{1–4} (Figure 1). Together with the weaker Antilles Current⁵, the
¹⁹ Florida Current forms the major western boundary current in the subtropical North Atlantic Ocean
²⁰ at 27°N , providing closure to the wind-driven interior gyre circulation^{6,7}, and acting as a vital limb

²¹ of the Atlantic meridional overturning circulation⁸. Due to its transport of heat and other tracers,
²² the Florida Current plays an important role in climate^{9,10}.

²³ The integrated volume transport of the Florida Current (hereafter Florida Current transport)
²⁴ has been monitored nearly continuously at $\sim 27^{\circ}\text{N}$ since 1982 by means of abandoned submarine
²⁵ telephone cables between West Palm Beach and Grand Bahama Island^{1–4} (Figure 1). Before then,
²⁶ observations were made occasionally as part of short hydrographic cruises or brief field campaigns,
²⁷ each measuring a different component of the flow at a different location. Earlier observations^{11,12}
²⁸ only measured the near-surface transports, but missed transports at depth. Later full-depth transport
²⁹ measurements^{13–16} were made variously between Florida and Havana, Cay Sal Bank, the Cat Cays,
³⁰ or Bimini, capturing flow through Yucatán Channel, but omitting flows through Nicholas, Santaren,
³¹ or Northwest Providence Channels, all of which contribute to the transport at 27°N (Figure 1). Such
³² disparities make it difficult to produce a stable instrumental estimate of Florida Current transport
³³ through time. Without such a coherent, longterm estimate, it has been unclear whether the Florida
³⁴ Current has undergone multidecadal- or longer-timescale change. Meinen et al.² concluded that
³⁵ the extant data, “provide no evidence for a longterm trend in the Florida Current transport,” during
³⁶ 1964–2009. However, it remains unclear whether a trend would emerge in a longer, more complete
³⁷ transport history.

³⁸ Questions of possible longterm changes in Florida Current transport bear on hypotheses that
³⁹ the Atlantic meridional overturning circulation is weakening or has weakened. Proxy indicators,
⁴⁰ including surface and subsurface ocean temperatures at subpolar latitudes and sortable silts from

41 sediment cores off Cape Hatteras, suggest that the deep return flow of the meridional overturning
42 circulation weakened either continuously during the twentieth century or earlier at the end of the
43 Little Ice Age^{17–19}. Yet, uncertainties in the proxies and their relationship to overturning render the
44 robustness of these suggestions unclear. Models simulate that, under climate change, a slowing of
45 the deep overturning circulation is balanced by a weakening surface western boundary current^{20,21}.
46 A determination of whether the Florida Current transport changed over the past century would thus
47 serve as a test of both model simulations and hypotheses of a reduced deep overturning.

48 Previous authors reasoned that sea level from coastal tide gauges is informative of changes
49 in Florida Current transport^{12,14,22,23}. These arguments are predicated on the notion of geostrophic
50 balance—on timescales longer than a day, the northward flow through Florida Straits imparts an
51 eastward acceleration owing to the Coriolis force that is counteracted by a pressure gradient across
52 the Florida Straits, which manifests as a sea-level difference that can be observed by tide gauges on
53 opposite sides of the Florida Current. However, circulation inferences based on tide gauges need to
54 be made cautiously. Tide gauges measure the distance between the sea surface and Earth’s crust at
55 the coast. So, they observe not only large-scale ocean dynamics, but also coastally trapped signals²⁴
56 and isostatic geophysical phenomena, including changes in the planet’s gravity field and rotation
57 vector, and viscoelastic deformation of the solid Earth²⁵. Tide-gauge data are also heterogeneously
58 distributed in space and time. Long, continuous records are available at some southeastern USA
59 and Caribbean sites far afield of the submarine cable at 27°N, but extant records from tide gauges
60 close to the cable’s endpoints near West Palm Beach and Grand Bahama are short, incomplete, and
61 largely not overlapping with one another²⁶.

62 To overcome these challenges, I use probabilistic data assimilation^{27,28} to estimate annual
63 Florida Current transport at 27°N over the past 110 y (see Methods). The estimate is based on 1,390
64 y of annual coastal sea level from 46 tide gauges²⁶ in the southeastern USA and Caribbean during
65 1909–2018 (Figure 1a) and 37 y of annual Florida Current transport from cable measurements^{1–4}
66 since 1982 (Figure 2a). Sea level is represented as a process with spatial correlation and temporal
67 memory. The Florida Current transport is related to the difference in sea level across Florida Straits
68 through geostrophy, but account is also taken of non-oceanographic and ageostrophic impacts on
69 sea level and transport. The data are cast as corrupt, imperfect versions of the processes. Bayes'
70 rule is used to invert the model equations, and solutions are generated using numerical methods.
71 The model equations are coupled, sharing information across space, time, and processes, allowing
72 data gaps to be filled and unobserved processes to be estimated. The solution is fully probabilistic,
73 and comprises thousands of ensemble members, each an equally likely history of transport that is
74 consistent with the data and model equations. This allows the calculation of subtle spatiotemporal
75 statistics, for example, the probability density function of the magnitude or timing of the minimum
76 or maximum decadally averaged transport value during the study period (see Methods). Residual
77 analyses and synthetic data experiments demonstrate the appropriateness of the algorithm and show
78 that it accurately estimates the quantities of interest given the data (cf. Supplementary Information).

79 Weakening of the Florida Current

80 The probabilistic Florida Current transport reconstruction is summarized in Figure 2a. The 110-y
81 mean transport is 32.6 ± 1.4 Sv (Supplementary Figure 1a), which is likely (probability $P = 0.87$)

82 larger than the mean over 1982–2018 (31.8 ± 0.1 Sv). This implies that the cable data^{1–4} probably
83 represent the Florida Current in a reduced state of transport. Unless otherwise stated, \pm values are
84 95% posterior credible intervals estimated from the Bayesian model. Estimated uncertainties since
85 1982 are comparatively small, and essentially reflect instrumental errors on the cable data, which
86 place strong observational constraints on the process. Before then, cable data are unavailable, and
87 the inference is largely constrained by tide gauge data, which have a more uncertain relationship
88 to the transport and are sparser earlier in time, resulting in larger errors that grow into the past.

89 Superimposed on the mean are interannual-to-decadal fluctuations in transport (Figure 2a).
90 The standard deviation of annual transports is 1.3 Sv (posterior median estimate). A 3.3 ± 1.1 Sv
91 weakening from 1997–1998 to 1999–2000, when there was a gap in cable measurements and low
92 transports were seen upstream in Yucatán Channel²⁹, was followed by a 2.5 ± 1.1 Sv strengthening
93 from 1999–2000 to 2001–2002 (Supplementary Figure 1c). Decadal-average transport was likely
94 ($P \geq 0.79$) greater than the longterm average during 1922–1932 (33.6 ± 2.8 Sv) and 1956–1966
95 (33.0 ± 1.7 Sv), but less than average in 1946–1956 (32.2 ± 2.0 Sv) and 1986–1996 (31.7 ± 0.2 Sv)
96 (Supplementary Figure 1d). A wavelet coherence analysis demonstrates that transport fluctuations
97 can be related to major modes of surface climate variation (Supplementary Figure 2). The transport
98 is probably ($P > 68\%$) coherent with the North Atlantic Oscillation³⁰ at 2–8-y periods centered
99 between the late 1970s and early 2000s, consistent with past studies of cable data^{1,31}; coherence is
100 also found at 2–4-y periods around 1960 and 8-y periods between the late 1930s and early 1950s,
101 which have not been previously reported, and possibly result from changes in subtropical wind
102 curl mediated by planetary waves³¹. Transport is also likely ($P > 0.68$) coherent with Atlantic

¹⁰³ Multidecadal Variability³² at 2–16-y periods centered on the mid 1990s and 16-y periods from the
¹⁰⁴ late 1940s to early 2000s. The weaker coherence earlier in time could reflect nonstationarity in the
¹⁰⁵ relationship between transport and climate, or the growth in transport uncertainties into the past.

¹⁰⁶ Changes are also apparent on the longest timescales. The transport trend during 1909–2018 is
¹⁰⁷ $-1.7 \pm 3.7 \text{ Sv century}^{-1}$, which overlaps zero, but implies that transport likely ($P = 0.82$) declined
¹⁰⁸ (Supplementary Figure 1b). This inference of a longterm weakening is qualitatively insensitive to
¹⁰⁹ the selection of time period. Computing differences between all pairs of decadal averages, I find
¹¹⁰ most (67%) instances are such that transport probably ($P > 0.68$) declined from one decade to
¹¹¹ another (Figure 3). For example, it is very likely ($P = 0.93$) transport weakened from 1920–1930
¹¹² ($2.1 \pm 2.9 \text{ Sv}$), and extremely likely ($P = 0.96$) that it declined from 1970–1980 ($1.2 \pm 1.2 \text{ Sv}$)
¹¹³ to the present more than expected from a stationary red-noise process. Indeed, if the transport was
¹¹⁴ stationary, extrema would be uniformly likely to occur at any point over a given time period, while
¹¹⁵ in the presence of a longterm decline, the maximum transport would be more likely to occur at the
¹¹⁶ beginning and the minimum transport at the end of the period. Consistent with the latter case, the
¹¹⁷ minimum decadal-average transport ($31.1 \pm 1.0 \text{ Sv}$) likely ($P = 0.74$) started sometime after 2002,
¹¹⁸ and the maximum decadal average ($34.1 \pm 2.5 \text{ Sv}$) likely ($P = 0.70$) ended some year before 1936
¹¹⁹ (Figure 2b). Timing of the extrema cannot be explained in terms of fluctuations about a stationary
¹²⁰ mean. After subtracting the longterm trend (Supplementary Figure 1b), I find that it would have
¹²¹ been unlikely ($P = 0.18$) that the minimum transport would have started after 2002, and chances
¹²² would have been lower ($P = 0.38$) that the maximum would have ended before 1936 (Figure 2c).

123 **Relation Between Florida Current Transport and Sea-Level Difference Across Florida Straits**

124 In addition to transport (Figure 2a), the Bayesian algorithm also solves for the regression coefficient
125 between the transport and sea-level difference across Florida Straits (see Methods). The estimated
126 change in transport per unit change in sea-level difference is $0.21 \pm 0.11 \text{ Sv cm}^{-1}$ (Supplementary
127 Figure 3a). Geostrophy allows interpretation of this value in terms of an effective depth describing
128 the vertical scale over which velocity variations decay in amplitude from the surface to the bottom
129 within Florida Straits^{33,34}. Following Little et al.³⁴, I multiply by the ratio of the Coriolis parameter
130 over gravity ($\sim 7 \times 10^{-6} \text{ s m}^{-1}$ at 27°N), obtaining an effective depth of $144 \pm 74 \text{ m}$. This estimate is
131 consistent with the vertical structure of northward currents observed *via* shipboard acoustic doppler
132 current profiler aboard the R/V Walton Smith during 70 cruises across Florida Straits at 27°N over
133 2001–2018. At the longitude of the core of the current, the average meridional velocity taken over
134 all cruises decays almost linearly in the vertical from $\sim 1.2 \text{ m s}^{-1}$ near the surface to $\sim 0.9 \text{ m s}^{-1}$
135 and $\sim 0.6 \text{ m s}^{-1}$ at 200- and 400-m depth, respectively (Figure 4a). Computing standard deviations
136 in meridional velocity over cruises, I find that flow-variation amplitudes decay more exponentially
137 with depth, decreasing rapidly from $\sim 0.6 \text{ m s}^{-1}$ near the surface to $\sim 0.3 \text{ m s}^{-1}$ and $\sim 0.2 \text{ m s}^{-1}$
138 at 200- and 400-m depth, respectively (Figure 4b). Similar vertical structures of mean and variable
139 meridional currents were reported based on earlier observations made during 1982–1984 as part of
140 the Subtropical Atlantic Climate Studies Program³⁵.

141 I have assumed that the regression coefficient between sea level and transport is time invariant
142 (see Methods). To test whether this assumption is reasonable, I compute coherence and admittance

¹⁴³ between sea level and transport output from an ocean reanalysis product³⁶ spanning 1871–2010.

¹⁴⁴ Considering interannual to multidecadal periods, I find that transports and sea-level differences are

¹⁴⁵ coherent across all accessible timescales, such that the admittance amplitude (transfer function) is

¹⁴⁶ qualitatively insensitive to frequency band, and that the change in transport per a unit change in the

¹⁴⁷ sea-level difference is similar for interannual and multidecadal periods (Supplementary Figure 4).

¹⁴⁸ Importantly, I also find that the Bayesian algorithm successfully estimates the correct regression

¹⁴⁹ coefficient between the two quantities in a synthetic data experiment based on this ocean reanalysis

¹⁵⁰ product (see Supplementary Information). These findings suggest that assuming a constant-in-time

¹⁵¹ relationship between transport and sea-level difference is reasonable, and that my model correctly

¹⁵² estimates the relationship between the two quantities given the available data.

¹⁵³ **Distinguishing Dynamic and Static Sea-Level Differences Across Florida Straits**

¹⁵⁴ The meaningfulness of the transport estimate hinges on the model's ability to identify and separate

¹⁵⁵ dynamic and static components of the sea-level difference across the Florida Straits. The posterior

¹⁵⁶ solution for the 110-y trend in sea-level difference across the Florida Straits (Grand Bahama minus

¹⁵⁷ West Palm Beach) is $-0.2 \pm 1.0 \text{ mm y}^{-1}$ (Supplementary Figure 3b). This trend results from the

¹⁵⁸ competing influences of a dynamic trend in sea-level difference of $-0.9 \pm 2.2 \text{ mm y}^{-1}$ and a static

¹⁵⁹ trend of $0.7 \pm 2.3 \text{ mm y}^{-1}$ (Supplementary Figure 3b), which I interpret respectively as indicating

¹⁶⁰ differential trends in sea-surface height and vertical land motion across Florida Straits.

¹⁶¹ Several lines of independent observational evidence corroborate these model inferences based

on data from tide gauges and submarine cables. The Global Positioning System (GPS) provides an instrumental means for measuring vertical land motion. Version 6a of the dataset from Université de la Rochelle³⁷ provides continuous GPS records from three locations in southeastern Florida and two Bahamas locations (Supplementary Figure 5; Supplementary Table 1). Computing the average vertical velocity for the two Bahamas sites, and doing the same over the three sites in southeastern Florida, I determine after taking the difference that sea level is statically rising $1.0 \pm 1.3 \text{ mm y}^{-1}$ faster in the Bahamas than along southeastern Florida owing to differential land subsidence, where the \pm value is twice the estimated standard error, assuming that the standard errors provided with the data are independent (Supplementary Table 1). This rate is consistent with the static trend in the sea-level difference across Florida Straits determined by the Bayesian model.

Proxy records of sea level are informative of background rates of change unrelated to ocean dynamics. I consider recent standardized compilations of Holocene sea-level index points from the Caribbean and southeastern USA derived from coral reefs, mangrove peats, and other indicators^{38,39}. To estimate present-day rates of background change unrelated to circulation and climate, I consider only the locations in the databases that have at least three sea-level index points with best-estimate ages between 2,000 and 150 y before present. This criterion is satisfied by two sites in southeastern Florida and one site in the Bahamas (Supplementary Figure 5; Supplementary Table 2). Taking the difference between the linear trend fit to the index points from the Bahamas site and the average of the trends fit to the data at the two southeastern Florida locations, I estimate that sea level rose $0.6 \pm 0.6 \text{ mm y}^{-1}$ more rapidly in the Bahamas relative to southeastern Florida in the pre-industrial Common Era (Supplementary Table 2), where the \pm value is twice the standard error furnished by

183 ordinary least squares applied to the best estimates of proxy age and sea level. Interpreted in terms
184 of differential land motion, this sea-level trend difference revealed by proxy data suggests that the
185 difference in rates of vertical land motion between the Bahamas and southeastern Florida observed
186 by GPS is, at least partly, due to background geological effects (e.g., glacial isostatic adjustment).

187 Modern radar altimeters have observed sea-surface height over nearly the global ocean since
188 1993. Once adjusted for static effects, altimeter data can be interpreted in terms of surface currents.
189 I consider time series of along-track sea-surface height processed by the Centre of Topography of
190 the Oceans and the Hydrosphere⁴⁰ at the altimeter data points closest to Settlement Point on Grand
191 Bahama Island and Virginia Key in southeastern Florida (Supplementary Figure 5). Differencing
192 the two altimetric time series and fitting a linear trend, I determine that the average rate of change
193 in the sea-surface-height difference across Florida Straits over 1993–2017 was $-2.2 \pm 3.0 \text{ mm y}^{-1}$
194 (Supplementary Figure 6), where the \pm value is twice the standard error estimated accounting for
195 residual autocorrelation using repeated simulations with surrogate data⁴¹. This rate from altimetry,
196 while reflecting a relatively short period, basically agrees in sign and magnitude with the dynamic
197 trend in sea-level difference across Florida Straits inferred by the Bayesian model. Note that, while
198 it is closer to the western end of the submarine cable than Virginia Key, the West Palm Beach gauge
199 is not considered in this exercise based on altimetry data; given the geometry of the satellite tracks,
200 the closest altimeter data point to the latter gauge is $\sim 50 \text{ km}$ offshore, east of the Florida Current
201 core (cf. Figure 4; Supplementary Figure 5), and does not reflect sea level at the western boundary.

202 **Relation to Wind Stress and the Interior Gyre**

203 Assuming no changes in Bering Straits throughflow or evaporation and precipitation over the basin,
204 weakening of the Florida Current transport must have been balanced by changes in the interior gyre
205 or overturning transports at 27°N. To explore possible changes in the gyre, I calculate geostrophic
206 Sverdrup streamfunction⁶ using wind-stress curl from two reanalyses of the twentieth century^{42,43}.
207 Both yield a climatological southward transport of ~ -23 Sv at 27°N over 1900–2010 (Figure 5a),
208 consistent with basic expectations⁴⁴. However, the two reanalyses give conflicting trend estimates,
209 with one⁴² yielding a weaker northward trend of 1.9 ± 2.0 Sv century⁻¹, and the other⁴³ a stronger
210 southward trend of -4.2 ± 1.3 Sv century⁻¹ across 27°N (Figure 5b) where the \pm values are formal
211 estimates of the 95% confidence interval adjusted for residual autocorrelations⁴¹. Discrepancies are
212 apparent broadly over the subtropics, with one reanalysis product⁴³ suggesting spin-up of the gyre,
213 and the other⁴² spin-down. These results are unaffected if ageostrophic Ekman transports are also
214 included in the calculation (Figure 5b).

215 **Relation to the Deep Overturning**

216 The longterm weakening of the Florida Current found here is comparable to the slowing of the deep
217 overturning circulation hypothesized to have occurred over the past century^{17–19}. These hypotheses
218 are partly based on the facts that models consistently show strong correlation between overturning
219 streamfunction and sea-surface temperature in the North Atlantic subpolar gyre on decadal and
220 longer timescales^{17,18,45,46}, and that observations show a “warming hole” over the subpolar gyre,

221 where sea-surface temperatures have recently fallen by 0.3–0.9°C century⁻¹ relative to the global
222 average^{43,47,48} (Figure 6).

223 To test whether the inferred weakening of the Florida Current, observed surface cooling over
224 the subpolar gyre, and hypothesized slowdown of the deep overturning are all physically consistent
225 with one another, I consider a simple ocean heat budget for the North Atlantic poleward of 27°N
226 (see Supplementary Information). I assume that decreasing ocean heat transport across 27°N due
227 to the combined weakening of the Florida Current and deep overturning is largely balanced by
228 increasing surface turbulent (sensible and latent) heat gain across the northern North Atlantic due
229 to the cooling sea-surface temperatures⁴⁹. Ignoring local heat storage, the sea-surface-temperature
230 change per unit change in transport is a function of the background mean sea-surface temperature,
231 vertical temperature stratification, and surface wind speed over the study region, along with the area
232 across which surface cooling takes place (see Supplementary Information). Choosing reasonable
233 parameter ranges, I derive a rough, first-principles estimate of 0.3–0.6°C Sv⁻¹. This is similar to
234 values of 0.2–0.5°C Sv⁻¹ found independently by dividing the observed sea-surface-temperature
235 trends across the subpolar gyre (Figure 6) by the posterior median estimate of the trend in Florida
236 Current transport over the past century (Supplementary Figure 1b). These numbers agree with a
237 range of 0.2–0.6°C Sv⁻¹ published based on regression analyses of sea-surface temperature and
238 overturning streamfunction from climate models^{17,18,45,46}.

239 **Conclusions**

240 Lack of knowledge about decadal and longer trends in ocean currents has been a key observational
241 uncertainty related to climate change. I used Bayesian data analysis^{27,28} to assimilate data from
242 submarine cables and tide gauges and to infer the evolution of the Florida Current transport at
243 27°N during 1909–2018. I found that Florida Current transport probably declined over the last 110
244 years, such that modern submarine cable data likely represent transport in a relatively reduced state,
245 and that the weakest decadal transport since the turn of the twentieth century probably occurred in
246 the last two decades. Results are consistent with observed cooling across the subpolar sea surface
247 and suggestions of a continuous decline in the deep overturning circulation over the past century,
248 and lend support to model predictions that a reduction of the deep overturning cell under climate
249 change is mirrored by a slowdown of the surface western boundary current.

250 Future studies should identify what caused the weakening of Florida Current transport, and
251 constrain whether changes in upper mid-ocean transports also took place. While systematic issues
252 with current reanalyses preclude conclusive results, possible longterm changes in the wind-driven
253 gyre circulation cannot be ruled out. Likewise, a recent data analysis⁵ determined that the Antilles
254 Current is highly variable on interannual and shorter timescales over 2005–2015, but that current's
255 behavior across decadal and longer timescales is unclear. Future efforts should also build upon this
256 Bayesian modeling framework to incorporate altimetric observations, GPS data, and proxy records
257 to better constrain the inference and reduce uncertainty.

- 258 1. Baringer, M. O., and J. C. Larsen. Sixteen Years of Florida Current Transport at 27°N, *Geophys.*
260 *Res. Lett.*, **28**, 317–3182 (2001).
- 261 2. Meinen, C. S., M. O. Baringer, and R. F. Garcia. Florida Current transport variability: An
262 analysis of annual and longer-period signals, *Deep-Sea Res.*, **57**, 835–846 (2010).
- 263 3. Larsen, J. C. Transport and heat flux of the Florida Current at 27°N derived from cross-stream
264 voltages and profiling data: theory and observations, *Phil. Trans. R. Soc. A.*, **338**, 169–236
265 (1992).
- 266 4. Garcia, R. F., and C. S. Meinen. Accuracy of Florida Current Volume Transport Measurements
267 at 27°N Using Multiple Observational Techniques, *J. Atmos. Ocean. Tech.*, **31**, 1169–1180
268 (2014).
- 269 5. Meinen, C. S., et al. Structure and Variability of the Antilles Current at 26.5°N, *J. Geophys.*
270 *Res.-Oceans*, **124**, 3700–3723 (2019).
- 271 6. Sverdrup, H. U. Wind-Driven Currents in a Baroclinic Ocean; With Application to the Equato-
272 rial Currents of the Eastern Pacific, *Proc. Natl. Acad. Sci. U.S.A.*, **33**, 318–326 (1947).
- 273 7. Stommel, H. The Westward Intensification of Wind-Driven Ocean Currents, *EOS T. Am. Geo-*
274 *phys. Un.*, **29**, 202–206 (1948).
- 275 8. McCarthy, G. D., et al. Measuring the Atlantic Meridional Overturning Circulation at 26°N,
276 *Prog. Oceanogr.*, **130**, 91–111 (2015).

- 277 9. Lynch-Stieglitz, J. The Atlantic Meridional Overturning Circulation and Abrupt Climate
278 Change, *Annu. Rev. Mar. Sci.*, **9**, 83–104 (2017).
- 279 10. Palter, J. B. The Role of the Gulf Stream in European Climate, *Annu. Rev. Mar. Sci.*, **7**, 113–
280 137 (2015).
- 281 11. Pillsbury, J. E. The Gulf Stream—a description of the methods employed in the investigation,
282 and the results of the research, *Rept. Supt., US Coast Geod. Surv.*, **Appendix 10**, 461–620
283 (1890).
- 284 12. Hela, I. The Fluctuations of the Florida Current, *B. Mar. Sci. Gulf Carib.*, **1**, 241–248 (1952).
- 285 13. Schmitz, W. J., and W. S. Richardson. On the transport of the Florida Current, *Deep-Sea Res.*,
286 **15**, 679–693 (1968).
- 287 14. Wunsch, C., D. V. Hansen, and B. D. Zetler. Fluctuations of the Florida Current inferred from
288 sea level records, *Deep-Sea Res.*, **16**, 447–470 (1969).
- 289 15. Niiler, P. P., and W. S. Richardson. Seasonal Variability of the Florida Current, *J. Mar. Res.*,
290 **31**, 144–167 (1973).
- 291 16. Brooks, I. H., and P. P. Niiler. Energetics of the Florida Current, *J. Mar. Res.*, **35**, 163–191
292 (1977).
- 293 17. Rahmstorf, S., et al. Exceptional twentieth-century slowdown in Atlantic Ocean overturning
294 circulation, *Nature Clim. Change*, **5**, 475–480 (2015).

- 295 18. Caesar, L., S. Rahmstorf, A. Robinson, G. Feulner, and V. Saba. Observed fingerprint of a
296 weakening Atlantic Ocean overturning circulation, *Nature*, **556**, 191–196 (2018).
- 297 19. Thornalley, D. J. R., et al. Anomalously weak Labrador Sea convection and Atlantic overturn-
298 ing during the past 150 years, *Nature*, **556**, 227–230 (2018).
- 299 20. Thomas, M. D., et al. Upper ocean manifestations of a reducing meridional overturning circu-
300 lation, *Geophys. Res. Lett.*, **39**, L16609 (2012).
- 301 21. Beadling, R. L., et al. Evaluation of Subtropical North Atlantic Ocean Circulation in CMIP5
302 Models against the Observational Array at 26.5°N and Its Changes under Continued Warming,
303 *J. Clim.*, **31**, 9697–9718 (2018).
- 304 22. Maul, G. A., et al. Sea Level Variation as an Indicator of Florida Current Volume Transport:
305 Comparisons with Direct Measurements, *Science*, **227**, 304–307 (1985).
- 306 23. Park, J., and W. Sweet. Accelerated sea level rise and Florida Current transport, *Ocean Sci.*,
307 **11**, 607–615 (2015).
- 308 24. Brink, K. H. Coastal-Trapped Waves and Wind-Driven Currents over the Continental Shelf,
309 *Annu. Rev. Fluid Mech.*, **23**, 389–412.
- 310 25. Tamisiea, M. E., and J. X. Mitrovica. The moving boundaries of sea level change: Understand-
311 ing the origins of geographic variability, *Oceanography*, **24**, 24–39 (2011).
- 312 26. Holgate, S. J., et al. New Data Systems and Products at the Permanent Service for Mean Sea
313 Level, *J. Coastal Res.*, **29**, 3, 493–504 (2013).

- 314 27. Tingley, M. P., et al. Piecing together the past: statistical insights into paleoclimatic recon-
315 structions, *Quaternary Sci. Rev.*, **35**, 1–22 (2012).
- 316 28. Ashe, E. L., et al. Statistical modeling of rates and trends in Holocene relative sea level, *Qua-*
317 *ternary Sci. Rev.*, **204**, 58–77 (2019).
- 318 29. Sheinbaum, J. et al. Flow structure and transport in the Yucatan Channel, *Geophys. Res. Lett.*,
319 **29**, 1040 (2002).
- 320 30. Jones, P. D., T. Jonsson, and D. Wheeler. Extension to the North Atlantic Oscillation Using
321 Early Instrumental Pressure Observations from Gibraltar and South-West Iceland, *Int. J. Cli-*
322 *matol.*, **17**, 1433–1450 (1997).
- 323 31. DiNezio, P. N., et al. Observed Interannual Variability of the Florida Current: Wind Forcing
324 and the North Atlantic Oscillation, *J. Phys. Oceanogr.*, **39**, 721–736 (2009).
- 325 32. Enfield, D. B., et al. The Atlantic multidecadal oscillation and its relation to rainfall and river
326 flows in the continental U.S., *Geophys. Res. Lett.*, **28**, 2077–2080 (2001)
- 327 33. Bingham, R. J., and C. W. Hughes. Signature of the Atlantic meridional overturning circulation
328 in sea level along the east coast of North America, *Geophys. Res. Lett.*, **36**, L02603 (2009).
- 329 34. Little, C. M., et al. The Relationship between United States East Coast Sea Level and the
330 Atlantic Meridional Overturning Circulation: a Review, *J. Geophys. Res.-Oceans*, **124**, 6435–
331 6458 (2019).

- 332 35. Leaman, K. D., R. L. Molinari, and P. S. Vertes. Structure and Variability of the Florida Current
333 at 27°N: April 1982–July 1984, *J. Phys. Oceanogr.*, **17**, 565–583 (1987).
- 334 36. Giese, B. S., and S. Ray. El Niño variability in simple ocean data assimilation (SODA), 1871–
335 2008, *J. Geophys. Res.*, **116**, C02024 (2011).
- 336 37. Santamaría-Gómez, A., et al. Uncertainty of the 20th century sea-level rise due to vertical land
337 motion errors, *Earth Planet. Sci. Lett.*, **473**, 24–32 (2017).
- 338 38. Khan, N. S., et al. Drivers of Holocene sea-level change in the Caribbean, *Quaternary Sci.*
339 *Rev.*, **155**, 13–36 (2017).
- 340 39. Love, R., et al. The contribution of glacial isostatic adjustment to projections of sea-level
341 change along the Atlantic and Gulf coasts of North America, *Earth's Future*, **4**, 440–464 (2016).
- 342 40. Birol, F., et al. Coastal applications from nadir altimetry: Example of the X-TRACK regional
343 products, *Adv. Space Res.*, **59**, 936–953 (2017).
- 344 41. Theiler, J., et al. Testing for nonlinearity in time series: the method of surrogate data, *Physica*
345 *D*, **58**, 77–94 (1992).
- 346 42. Compo, G. P., et al. The Twentieth Century Reanalysis Project, *Q. J. R. Meteorol. Soc.*, **37**,
347 1–28 (2011).
- 348 43. Poli, P., et al. ERA-20C: An Atmospheric Reanalysis of the Twentieth Century, *J. Clim.*, **29**,
349 4083–4097 (2016).

- 350 44. Gray, A. R., and S. C. Riser. A Global Analysis of Sverdrup Balance Using Absolute
351 Geostrophic Velocities from Argo, *J. Phys. Oceanogr.*, **44**, 1213–1229 (2014).
- 352 45. Drijfhout, S., et al. Is a Decline of AMOC Causing the Warming Hole above the North Atlantic
353 in Observed and Modeled Warming Patterns?, *J. Clim.*, **25**, 8373–8379 (2012).
- 354 46. Roberts, C. D., F. K. Garry, and L. C. Jackson. A Multimodel Study of Sea Surface Temper-
355 ature and Subsurface Density Fingerprints of the Atlantic Meridional Overturning Circulation,
356 *J. Clim.*, **26**, 9155–9174 (2013).
- 357 47. Rayner, N. A., et al. Global analyses of sea surface temperature, sea ice, and night marine air
358 temperature since the late nineteenth century, *J. Geophys. Res.*, **108**, 4407 (2003).
- 359 48. Kaplan, A., et al. Analyses of global sea surface temperature 1856–1991, *J. Geophys. Res.*,
360 **103**, 18567–18589 (1998).
- 361 49. Gulev, S. K., M. Latif, N. Keenlyside, W. Park, and K. P. Koltermann. North Atlantic Ocean
362 control on surface heat flux on multidecadal timescales, *Nature*, **499**, 464–467 (2013).
- 363 50. Lumpkin, R., and G. C. Johnson. Global ocean surface velocities from drifters: Mean, vari-
364 ance, El Niño-Southern Oscillation response, and seasonal cycle, *J. Geophys. Res.-Oceans*, **118**,
365 2992–3006 (2013).

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370 **Author Contributions** C.G.P. conceived the study, formulated the model framework, performed the anal-
371 yses, and wrote the manuscript.

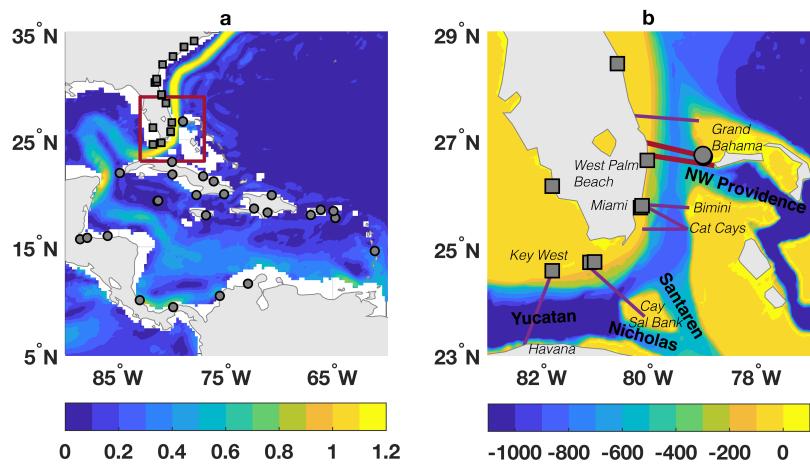
372 **Competing Interests** The author declares that they have no competing financial interests.

373 **Correspondence** Correspondence and requests should be addressed to C.G.P. (cpiecuch@whoi.edu).

374 **Data Availability** The tide gauge data and submarine cable data that support the main findings of this
375 study are available from the Permanent Service for Mean Sea Level (PSMSL; <http://www.psmsl.org/>) and
376 the National Oceanic and Atmospheric Administration Atlantic Oceanographic and Meteorological Labora-
377 tory (NOAA AOML; <https://www.aoml.noaa.gov/>), respectively. Ancillary data sets, used for interpretation
378 and not incorporated into the Bayesian model, and their availabilities are as follows: Surface drifter data
379 of surface current speeds shown in Figure 1a are available from NOAA AOML; Global Digital Elevation
380 Model bathymetry shown in Figures 1b and 4 are available from NOAA National Geophysical Data Center
381 (NGDC; <https://www.ngdc.noaa.gov/>); Cruise data from the R/V Walton Smith shown in Figure 4 are avail-
382 able from NOAA AOML; Reanalysis wind-stress fields shown in Figure 5 are available from the Woods
383 Hole Oceanographic Institution Community Storage Server (WHOI CCS; <https://cmip5.whoi.edu>); Gridded
384 data sets of sea-surface temperature shown in Figure 6 are available from the WHOI CCS, UK Met Of-
385 fice Hadley Centre (<https://www.metoffice.gov.uk/hadobs/>), and NOAA Earth System Research Laboratory
386 Physical Sciences Division (ESRL PSD; <https://www.esrl.noaa.gov/psd/>); Time series of climate indices
387 shown in Supplementary Figure 2 are available from NOAA ESRL PSD; Global Positioning System data
388 of vertical land motion rates shown in Supplementary Table 1 are available from Système d’Observation du

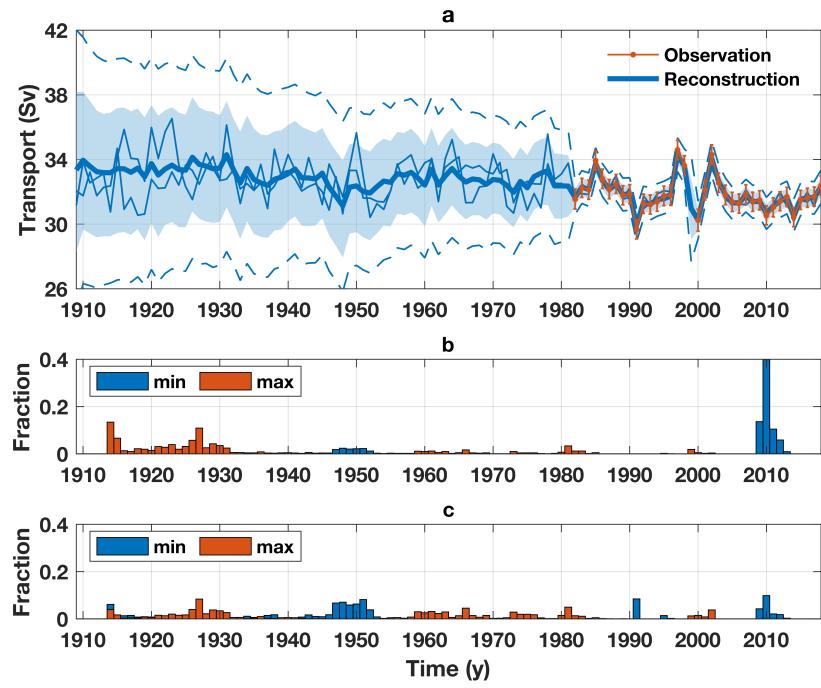
389 Niveau des Eaux Littorales (SONEL; <http://www.sonel.org/>); Proxy relative sea-level index points shown
390 in Supplementary Table 2 are taken from Khan et al.³⁸ and Love et al.³⁹; Satellite-altimetric time series of
391 sea-surface height shown in Supplementary Figure 6 are available from Centre of Topography of the Oceans
392 and the Hydrosphere (CTOH; <http://ctoh.legos.obs-mip.fr/>); Model estimates of glacial isostatic ad-
393 justment rates used in the synthetic data experiments are available from PSMSL; Global-mean thermosteric
394 sea level from the Community Climate System Model Version 4 used in the synthetic data experiments
395 was downloaded from the WHOI CCS; Model solutions from the Simple Ocean Data Experiment (SODA)
396 shown in Supplementary Figures 4, 15 and used in the synthetic data experiments are available from the
397 University of Hawaii Asia-Pacific Data-Research Center (<http://apdrc.soest.hawaii.edu/>). Maps in display
398 items were produced using the Mapping Toolbox in MATLAB.

399 **Code Availability Statement** The computer code used to run the Bayesian model and produce the results
400 in this study, written in the MATLAB software environment, is available at the corresponding author's
401 GitHub website (<https://github.com/christopherpiecuch>).



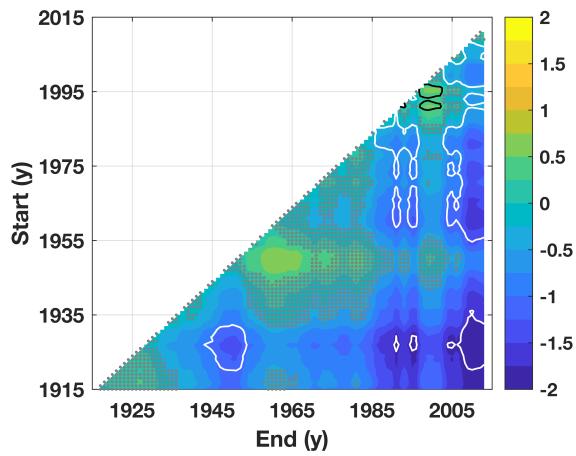
402

[Fig. 1]



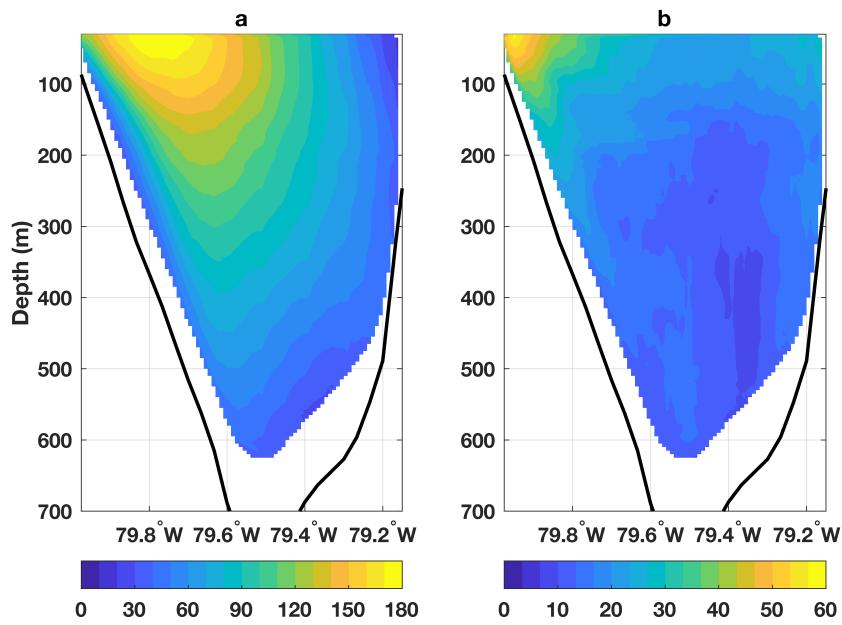
403

[Fig. 2]



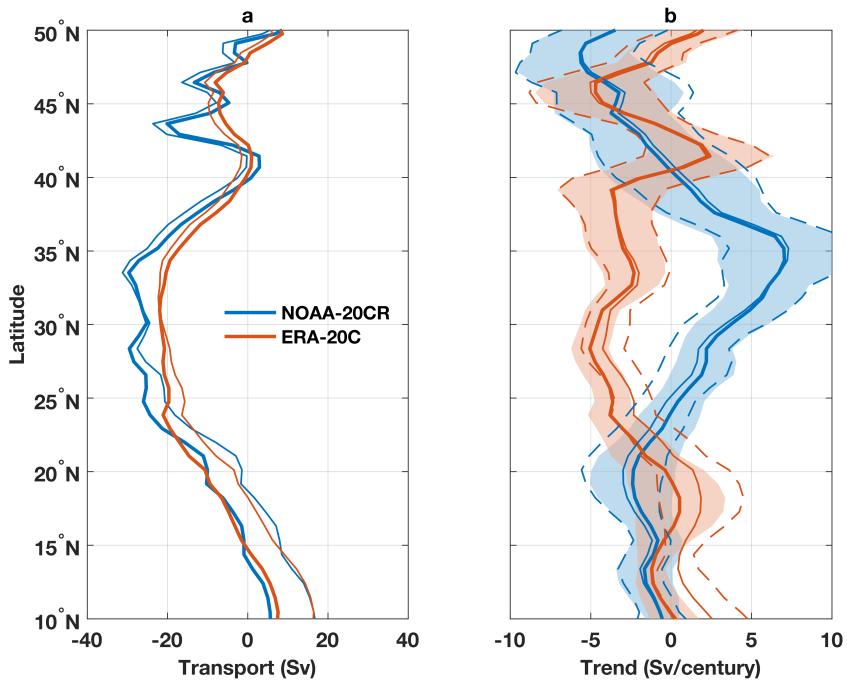
404

[Fig. 3]



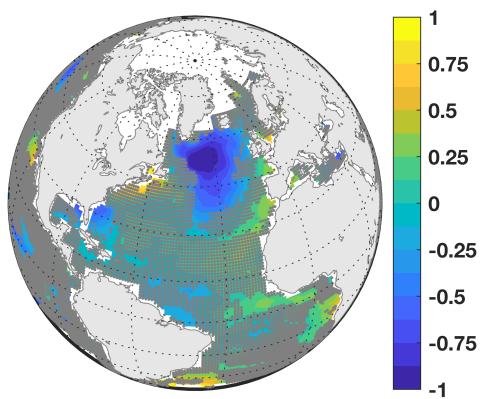
405

[Fig. 4]



406

[Fig. 5]



407

[Fig. 6]

408 **Methods**

409 **Observational data used in the Bayesian model** I use annual relative sea level from 46 tide
410 gauges in the southeastern USA (19 records), Caribbean Islands (20 records), southeastern Central
411 America (5 records), and northern South America (2 records) during 1909–2018 (Figure 1a; Sup-
412 plementary Figure 7; Supplementary Table 3). Data were downloaded from the Permanent Service
413 for Mean Sea Level (PSMSL) Revised Local Reference (RLR) database²⁶ on 4 February 2019.
414 The study period is the longest interval such that, for each year, data is available from at least one
415 southeastern USA tide gauge and at least one gauge in the Caribbean Islands, southeastern Central
416 America, or northern South America. Over the study period, each tide gauge returns on average
417 ~ 30 y of data, but some have as few as ~ 10 y of data, whereas others have as many as ~ 100 y.
418 The time series together constitute 1,390 y of data over the study period ($\sim 27\%$ completeness).

419 I also use Florida Current transport from submarine telephone cables at 27°N between West
420 Palm Beach and Grand Bahama (Figure 1b)^{1–4}. Using electromagnetic theory, changes in the flow
421 can be estimated from voltages induced across the cable due to the transport of charged particles
422 by the variable current³. The original cable spanned from Jupiter Inlet to Settlement Point, giving
423 measurements from 18 March 1982 to 22 October 1998; observations resumed on 19 June 2000
424 based on a cable running from West Palm Beach to Eight Mile Rock (Figure 1b). Transports are
425 provided by the National Oceanic and Atmospheric Administration (NOAA) at 1-day intervals, but
426 the data have an effective sampling rate of 3 days, due to low-pass filtering applied to the original
427 observations. I use annual averages of the daily data (Figure 2a). Given a standard error of 1.7 Sv

428 on the daily values⁴, I estimate standard errors on the annual averages of 0.30–0.35 Sv, depending
429 on data availability in any given year, consistent with values computed by Garcia and Meinen⁴.

430 **Bayesian framework**

431 I apply a hierarchical dynamical spatiotemporal model^{27,28,51,52} to the submarine-cable data and
432 tide-gauge records to infer annual changes in Florida Current transport and coastal sea level. The
433 model comprises three levels: a process level describing how the quantities of interest relate to
434 one another, and vary in space and time; a data level specifying how the imperfect available data
435 correspond to the quantities of interest; and a parameter level placing prior constraints on the un-
436 certain parameters in the process and data levels. My model builds on the Bayesian algorithm of
437 Piecuch et al.⁵³, who studied the origin of spatial variation in sea-level trends on the east coast of
438 the USA during 1900–2017. Here I develop new equations to consider an expanded geographic re-
439 gion, incorporate the submarine-cable data, and represent the relationship between Florida Current
440 transport and the difference in coastal sea level across the Florida Straits. See the Supplementary
441 Information for residual analyses and synthetic data experiments that establish the appropriate-
442 ness of the model given the data, and exemplify its ability to accurately estimate the quantities of
443 interest given the available incomplete, noisy, biased data.

444 **Process level**

445 **Coastal sea level** Coastal relative sea level is a process with spatiotemporal covariance^{54,55}. As
 446 in Piecuch et al.⁵³, I model sea level, $\boldsymbol{\eta}_k = [\eta_{1,k}, \dots, \eta_{N,k}]^\top$, at steps $k \in \{1, \dots, K\}$ and sites
 447 $n \in \{1, \dots, N\}$, as the sum of a spatially correlated autoregressive process of order 1 and a large-
 448 scale spatial field of linear temporal trends,

$$\boldsymbol{\eta}_k - \mathbf{b}t_k = r(\boldsymbol{\eta}_{k-1} - \mathbf{b}t_{k-1}) + \mathbf{e}_k. \quad (1)$$

449 In Eq. (1), t_k is the time at step k , r is the lag-1 autocorrelation coefficient, \mathbf{b} is the spatial
 450 vector of temporal trends, and \mathbf{e}_k is an innovation sequence driving the autoregressive process.

451 Supplementary Table 4 describes all of the model parameters. I set $\sum_{k=1}^K t_k = 0$ to represent $\boldsymbol{\eta}_k$ as
 452 anomalies from a time mean. The trend vector \mathbf{b} is modeled as a random normal field with spatial
 453 structure, $\mathbf{b} \sim \mathcal{N}(\mu \mathbf{1}_N, \Pi)$, such that μ is the spatial mean, $\mathbf{1}_X$ is a $X \times 1$ column vector of ones,
 454 and,

$$\Pi_{ij} = \pi^2 \exp(-\lambda |\mathbf{s}_i - \mathbf{s}_j|). \quad (2)$$

455 Here π^2 is the partial sill, λ is the inverse range, and $|\mathbf{s}_i - \mathbf{s}_j|$ is distance between target sites \mathbf{s}_i and
 456 \mathbf{s}_j . The symbol \sim means “is distributed as” and $\mathcal{N}(\mathbf{p}, \mathbf{q})$ is the multivariate normal distribution
 457 with mean vector \mathbf{p} and covariance matrix \mathbf{q} .

458 I cast \mathbf{e}_k as a temporally independent, identically distributed (iid), but spatially correlated
 459 vector with zero mean, $\mathbf{e}_k \sim \mathcal{N}(\mathbf{0}_N, \Sigma)$, where $\mathbf{0}_X$ is a $X \times 1$ column vector of zeros, and,

$$\Sigma_{ij} = (c_{ij}) \sigma^2 \exp(-\phi |\mathbf{s}_i - \mathbf{s}_j|). \quad (3)$$

460 Here σ^2 is the partial sill and ϕ is the inverse range. Matrix element $c_{ij} = 1$ if locations \mathbf{s}_i and
 461 \mathbf{s}_j are either both on the southeastern USA or both along the Caribbean, Central America, or

462 South America. Otherwise, $c_{ij} = 0$. That is, sea level covaries within, but not between, these re-
 463 gions. This spatial covariance structure is motivated by previous analyses of tide-gauge records and
 464 satellite-altimetry data. Thompson and Mitchum⁵⁶ applied clustering methods to low-pass-filtered
 465 tide-gauge records during 1952–2001, finding that the Caribbean Sea (which in their analysis com-
 466 prised Cuba, Puerto Rico, and Colombia) formed one cluster of coherent sea-level variation, and
 467 the southeastern USA (from Florida to North Carolina) formed another cluster. Zhao and Johns⁵⁷
 468 determined that Florida Current transports over 1993–2011 were positively correlated with sea-
 469 surface height over the Caribbean Sea (including the Bahamas) and along southeastern Central
 470 America, but negatively correlated with sea-surface height on the southeastern USA coast on in-
 471 terannual timescales.

472 **Florida Current transport** For periods longer than a day, the momentum balance across Florida
 473 Straights will be nearly geostrophic. Assuming that subsurface pressure signals are vertically
 474 coherent^{33,34}, variations in Florida Current transport should therefore be correlated with changes
 475 in the sea-level difference across Florida Straits. Based on this reasoning, I assume that the re-
 476 lationship between annual Florida Current transport, $\mathbf{T} = [T_1, \dots, T_K]^\top$, and coastal sea level,
 477 $\eta = [\eta_1, \dots, \eta_K]$, at times $\mathbf{t} = [t_1, \dots, t_K]^\top$ can be written as,

$$\mathbf{T} = \bar{T}\mathbf{1}_K + \rho\eta^\top\Delta + \alpha\mathbf{t} + \mathbf{w}. \quad (4)$$

478 Here \bar{T} is the time-mean transport and ρ is a scalar coefficient representing the change in transport
 479 per unit change in sea-level difference across the Florida Straits. I assume that ρ is a constant,
 480 and does not vary with time period or frequency band. While it might appear simplistic, this

481 assumption is justified based on admittance and coherence analysis applied to output from an ocean
482 general circulation model (see Supplementary Information). The $N \times 1$ vector Δ is a differencing
483 operator, such that $\Delta_i = 1$ if site i is Settlement Point (the tide gauge nearest to the eastern end of
484 the submarine cable in the Bahamas), $\Delta_i = -1$ if site i is West Palm Beach (the closest tide gauge
485 to the western end of the cable in southeastern Florida), and zero otherwise. Hence, $\rho\eta^T\Delta$ is the
486 sea-level difference across Florida Straits converted into units of a transport.

487 The remaining terms in Eq. (4) account for other effects unrelated to large-scale geostrophic
488 ocean dynamics. The scalar α represents an apparent trend in T , included to correct for longterm
489 static sea-level changes unrelated to ocean dynamics, for example, due to glacial static adjustment²⁵.
490 That is, $b^T\Delta$ is the difference in sea-level trends across Florida Straits, resulting from both dynamic
491 processes and static effects. Hence, in Eq. (4), $b^T\Delta + \alpha/\rho$ represents the dynamic component of
492 the difference in sea-level trends across Florida Current, and $-\alpha/\rho$ constitutes the static component
493 of the trend in sea-level differences across the Florida Straits (Supplementary Figure S3). Satellite
494 altimetry, GPS data, and proxy sea-level index points support this interpretation of Eq. (4) (cf. the
495 Main Text). I also include $w = [w_1, \dots, w_K]^T$, which is modeled as iid uncorrelated white noise,
496 $w_k \sim \mathcal{N}(0, \omega^2)$, with variance ω^2 , to parameterize the response to local atmospheric or terrestrial
497 forcing, such as river runoff, air pressure, or wind stress across Florida Straits.

498 **Data level**

499 **Tide-gauge records** Following Piecuch et al.⁵³, I represent annual data from tide gauges, $\mathbf{z}_k =$
 500 $[z_{1,k}, \dots, z_{M_k,k}]^\top$, at $M_k \leq N$ locations at time step k , as corrupted (incomplete, noisy, biased)
 501 versions of the underlying $\boldsymbol{\eta}_k$ process,

$$\mathbf{z}_k = \mathsf{H}_k \boldsymbol{\eta}_k + \mathbf{d}_k + \mathsf{F}_k (\mathbf{a} t_k + \boldsymbol{\ell}). \quad (5)$$

502 Here \mathbf{d}_k is a random error sequence, which is modeled as a spatially and temporally uncorrelated
 503 normal field, $\mathbf{d}_k \sim \mathcal{N}(\mathbf{0}_{M_k}, \delta^2 \mathbf{I}_{M_k})$, with variance δ^2 . A vector of location-specific offsets $\boldsymbol{\ell}$ are
 504 imposed and represented as a spatially uncorrelated Gaussian field, $\boldsymbol{\ell} \sim \mathcal{N}(\nu \mathbf{1}_M, \tau^2 \mathbf{I}_M)$, with
 505 mean ν , variance τ^2 , and where M is the number of tide gauges, such that $N \geq M \geq M_k \forall k$.
 506 Purely local error trends in the data \mathbf{a} are also modeled as a random normal field without spatial
 507 correlation, $\mathbf{a} \sim \mathcal{N}(\mathbf{0}_M, \gamma^2 \mathbf{I}_M)$, with variance γ^2 . Finally, H_k and F_k are selection matrices, filled
 508 with zeros and ones, which pick out $\boldsymbol{\eta}_k$, \mathbf{a} , or $\boldsymbol{\ell}$ values at the observation sites for time t_k .

509 **Submarine-cable measurements** I assume that L annual data values from the submarine cable,
 510 $\mathbf{x} = [x_1, \dots, x_L]^\top$, are available and represent imperfect (incomplete and noisy) versions of the
 511 underlying \mathbf{T} process,

$$\mathbf{x} = \mathbf{G}\mathbf{T} + \mathbf{u}. \quad (6)$$

512 Here \mathbf{G} is a $L \times K$ selection matrix, picking out years when cable data are available, and $\mathbf{u} =$
 513 $[u_1, \dots, u_L]^\top$ is a zero-mean random data error sequence, $u_l \sim \mathcal{N}(0, \xi_l^2)$, where the ξ_l^2 are set

514 equal to the corresponding submarine-cable data standard error variances mentioned above and
515 computed based on the availability of data in any given year.

516 **Parameter level** To close the model, priors are placed on the parameters in the process- and
517 data-level equations. Similar to Piecuch et al.⁵³, I use proper, mostly conjugate prior forms. Prior
518 forms and hyperparameter values are given in Supplementary Table 5. The selection of the hyper-
519 parameter values follows the basic logic in Piecuch et al.⁵³. My philosophy is to employ diffuse
520 and uninformative priors. To quantify the importance of priors relative to the data, after I compute
521 the posterior solutions (see immediately below), I compare widths of the 95% credible intervals
522 from the posterior and prior probability distribution functions for each parameter (Supplementary
523 Table 6). If prior and posterior credible intervals have similar widths, then the posterior solutions
524 are largely determined by the prior assumptions. If posterior credible intervals are much narrower
525 than the prior credible intervals, then the posterior solutions are mostly constrained by the obser-
526 vations. Almost universally, the 95% posterior credible intervals are much narrower than the 95%
527 prior credible intervals (Supplementary Table 6), implying that posterior inference is drawn pre-
528 dominantly from the information content of the observations, and not overly influenced by prior
529 beliefs encoded into the model.

530 **Drawing samples from the posterior distribution** Given the model equations, I use Bayes'

531 rule, and assume that the posterior probability distribution function takes the form,

$$\begin{aligned} p(\eta, \mathbf{T}, \Theta | \mathbf{Z}, \mathbf{x}) &\propto p(\mathbf{Z}, \mathbf{x} | \eta, \mathbf{T}, \Theta) \times p(\eta, \mathbf{T} | \Theta) \times p(\Theta) \\ &= p(\boldsymbol{\eta}_0) \times p(\bar{T}) \times p(r) \times p(\sigma^2) \times p(\phi) \times p(\mu) \times p(\pi^2) \\ &\quad \times p(\lambda) \times p(\delta^2) \times p(\nu) \times p(\tau^2) \times p(\gamma^2) \times p(\rho) \times p(\alpha) \\ &\quad \times p(\omega^2) \times p(\mathbf{b} | \mu, \pi^2, \lambda) \times p(\boldsymbol{\ell} | \nu, \tau^2) \times p(\mathbf{a} | \gamma^2) \times p(\mathbf{x} | \mathbf{T}) \\ &\quad \times p(\mathbf{T} | \eta, \rho, \alpha, \omega^2, \bar{T}) \times \prod_{k=1}^K [p(\mathbf{z}_k | \boldsymbol{\eta}_k, \mathbf{a}, \boldsymbol{\ell}, \delta^2) \times p(\boldsymbol{\eta}_k | \boldsymbol{\eta}_{k-1}, \mathbf{b}, r, \sigma^2, \phi)] \end{aligned} \tag{7}$$

532 In Eq. (7), \mathbf{Z} is the structure of all tide-gauge data points, p is used to represent probability distribu-
533 tion function, $|$ is conditionality, \propto is proportionality, and $\Theta \doteq \{r, \sigma^2, \phi, \dots\}$ is used to represent
534 the set of all model parameters. I assume that the observations are conditionally independent,
535 provided the process and parameters.

536 Draws from the posterior probability distribution function are made as in Piecuch et al.⁵³.

537 I use Markov chain Monte Carlo (MCMC) methods, evaluating the full conditional distributions
538 for process and parameter values using a Gibbs sampler, but using Metropolis steps for the in-
539 verse range parameters. I run 200,000 MCMC iterations, setting initial process values to zero,
540 and drawing initial parameter values randomly from the respective prior distribution. To remove
541 initialization transients, I discard the first 100,000 iterations as burn in, and then I keep only 1
542 out of every 100 of the remaining 100,000 iterations to reduce serial correlation effects between
543 draws. Results shown here are based on a 3,000-element chain produced by performing the above
544 procedure 3 times and stitching together the resulting 1,000-member chains. Solutions for scalar

parameters are summarized in Supplementary Table 6. To evaluate convergence of the solution for each parameter, I compute the convergence monitor \hat{R} of Gelman and Rubin⁵⁸, which compares the variance within and between the 3 different 1,000-member solutions. In each case, $\hat{R} \sim 1.00$ (Supplementary Table 6), indicating that the solutions have converged.

549 Local and global uncertainty measures

The probabilistic nature of the model solutions allows for the calculation of both pointwise and pathwise uncertainty measures⁵⁹. Pointwise statistics measure probabilities locally. The light blue shading in Figure 2a represents the 95% pointwise posterior credible intervals computed from the transport solutions at each year of the reconstruction. The interpretation is that, for each year, there is a 95% chance that the true transport value falls within this blue shading.

Pathwise statistics measure probabilities more globally. The dashed blue lines in Figure 2a represent the 95% pathwise posterior credible intervals calculated from the transport estimates across all years of the reconstruction. These values are computed by widening the 95% pointwise posterior credible intervals until 95% of modeled transport time series are captured in their entirety. That is, there is a 95% chance that the full time series of transport does not stray outside the bounds of these pathwise credible intervals.

Other examples of pathwise statistics include values quoted in the text for the minimum and maximum decadal-average transports and the corresponding histograms of their timing shown Figure 2b, 2c. For each of the 3,000 ensemble members comprising the posterior solution, I smooth

564 the transport time series using an 11-point boxcar window, and then identify the minimum and
565 maximum transport values along with the times at which they occurred. These values vary from
566 one ensemble member to the next, and so performing this procedure for each ensemble member
567 allows me to populate histograms for the transport extrema and their occurrence times.

568 **Hypothesis testing**

569 In addition to generating posterior solutions for transport and sea level, the Bayesian model pro-
570 vides data-constrained estimates of the various model parameters (e.g., Supplementary Table 6).
571 This allows for rigorous hypothesis testing through simulation experiments. For example, in Fig-
572 ure 3, I show the change in decadal-average Florida Current transport between all possible pairs of
573 decades, and indicate the probability that such changes would have occurred given a stationary red-
574 noise process with the same autocorrelation and variance characteristics. As a specific instance, I
575 state that decadally averaged transport declined by 1.2 ± 1.2 Sv from 1970–1980 to the present,
576 and that this decline is extremely likely (probability $P = 0.96$) more than would be expected from
577 stationary red noise. This conclusion was determined as follows. First, I use the posterior trans-
578 port solutions to compute a histogram of transport averaged over 2008–2018 minus the transport
579 averaged over 1970–1980. Next, I use the posterior solutions for the scalar model parameters as
580 the basis for the simulation of a parallel set of 3,000 synthetic transport time series following Eqs.
581 (1) and (4) but with the trends (b and α) set to zero. Then, I populate histograms of the difference
582 between decadally averaged synthetic transport between 1970–1980 and 2008–2018. Finally, I
583 compute what fraction of the original posterior transport solutions shows more of a decline than is

- 584 shown by the stationary synthetic transport process, which, in this example case, is 0.96.
- 585 51. Cressie, N., & C. K. Wikle. Statistics for Spatio-Temporal Data, John Wiley & Sons, 588 pp
586 (2011).
- 587 52. Banerjee, S., B. P. Carlin, & A. E. Gelfand. Hierarchical Modeling and Analysis for Spatial
588 Data, Chapman and Hall, Boca Raton, 448 pp (2004).
- 589 53. Piecuch, C. G., et al. Origin of spatial variation in US East Coast sea-level trends during 1900–
590 2017, *Nature*, **564**, 400–404 (2018).
- 591 54. Bos, M. S., S. D. P. Williams, I. B. Araújo, & L. Bastos. The effect of temporal correlated
592 noise on the sea level rate and acceleration uncertainty, *Geophys. J. Int.*, **196**, 1423–1430.
- 593 55. Hughes, C. W., and M. P. Meredith. Coherent sea-level fluctuations along the global continental
594 slope, *Phil. Trans. R. Soc. A.*, **364**, 885–901 (2006).
- 595 56. Thompson, P. R., and G. T. Mitchum. Coherent sea level variability on the North Atlantic western
596 boundary, *J. Geophys. Res. Oceans*, **119**, 5676–5689, doi:10.1002/2014JC009999 (2014).
- 597 57. Zhao, J., and W. Johns. Wind-forced interannual variability of the Atlantic Meridional Over-
598 turning Circulation at 26.5°N, *J. Geophys. Res.-Oceans*, **119**, 2403–2419.
- 600 58. Gelman, A., and D. B. Rubin. Inference from iterative simulation using multiple sequences,
601 *Stat. Sci.*, **7**, 457–472 (1992).

- 602 59. Tingley, M. P., & P. Huybers. Recent temperature extremes at high northern latitudes unprece-
- 603 dented in the past 600 years, *Nature*, **496**, 201–205 (2013).

604 **Figure 1** Florida Current and study region. **a**, Gray squares (circles) are locations of tide
605 gauges in the southeastern USA (Caribbean). Shading is mean ocean surface current
606 speed (m s^{-1}) from surface-drifter data⁵⁰. Red box is area shown in (b). **b**, Details of
607 Florida Straits. Shading is ocean depth (m). Bold (light oblique) font indicates ocean
608 channels (land locations) mentioned in the text. Thick red lines are locations of submarine
609 cable measurements. Thin purple lines are locations of *in situ* measurements from past
610 studies^{13–16}.

611 **Figure 2** Florida Current transport. **a**, Blue shows posterior median (thick line), 95%
612 pointwise (light shading) and pathwise (dash dot) credible intervals, and two arbitrary
613 ensemble members (thin lines) from the probabilistic Florida Current transport solution.
614 Orange shows annual transport from raw submarine cable data plus and minus twice
615 the standard error². **b**, Histograms of modeled probabilities that the minimum (blue) and
616 maximum (orange) decadal average transport occurred centered on a given year. **c**, As
617 in (b) but histograms were calculated after having removed the corresponding longterm
618 trend. See Methods for discussion of statistics and uncertainty measures.

619 **Figure 3** Weakening of Florida Current transport over different periods. Shading shows
620 posterior median estimates of the change in decadal-average Florida Current transport
621 between all pairs of decades (Sv). Negative values indicate that transport fell from the
622 start to the end decade. Stippling indicates that it is as likely as not ($0.33 < P < 0.67$)
623 that transport rose or fell. White (black) contours encircle periods when it is very likely

624 ($P > 0.90$) that transport weakened (strengthened) from the start to the end decade
625 more than expected from a stationary red noise process (see Methods for discussion of
626 calculations of significance).

627 **Figure 4 Structure of the Florida Current within Florida Straits.** **a**, Mean northward
628 velocities (m s^{-1}) through Florida Straits from shipboard acoustic doppler current profiler
629 data from 70 research cruises of the R/V Walton Smith between 2001–2018. Values are
630 computed by interpolating all data between 26.9°N and 27.1°N from a given cruise onto a
631 common grid, and then averaging over all cruises. For a value to be shown at a longitude
632 and depth, data must have been available from at least 14 cruises. **b**, As in (a) but showing
633 the standard deviation in meridional velocities (m s^{-1}) across cruises.

634 **Figure 5 Changes in wind-stress curl and gyre circulation.** **a**, Thick lines are time-mean
635 geostrophic Sverdrup streamfunction⁶ computed from wind-stress curl from NOAA 20CR⁴²
636 and ERA 20C⁴³ reanalyses over 1900–2010 as a function of latitude in the North Atlantic.
637 Thin lines are the same, but also incorporate the ageostrophic Ekman transport integrated
638 across the basin. **b**, Median estimates (thick lines) and formal 95% confidence intervals
639 (colored shading) of the trend in Sverdrup streamfunction versus latitude during 1900–
640 2010 from the two reanalyses. Thin and dashed lines represent median estimates and
641 confidence intervals, respectively, with Ekman transports included.

642 **Figure 6 Changes in sea-surface temperature.** Shaded values are sea-surface tempera-
643 ture trends ($^\circ\text{C century}^{-1}$) since 1909 averaged over three products: ERA-20C⁴³, HadISST⁴⁷,

⁶⁴⁴ and Kaplan⁴⁸. Stippling indicates that the magnitude of the average trend is less than 2.35
⁶⁴⁵ (the $P = 0.95$ -value of the inverse t -distribution for 3 degrees of freedom) times the sam-
⁶⁴⁶ ple standard deviation computed across the three different products at a given grid cell,
⁶⁴⁷ and is meant as a rough indicator of where values are not significant.

¹ **Supplementary Information for “Weakening of the Gulf
2 Stream at Florida Straits over the past century inferred
3 from coastal sea-level data”**

⁴ Christopher G. Piecuch¹

⁵ ¹Woods Hole Oceanographic Institution, Woods Hole, Massachusetts, USA

⁶ **S1 Heat budget**

⁷ Previous studies interpret a “warming hole” over the subpolar North Atlantic Ocean, where surface
⁸ temperatures have cooled relative to the global mean over the past century (Figure 6), in terms of a
⁹ slowing deep Atlantic meridional overturning circulation^{1–3}. These interpretations are partly based
¹⁰ on regression analyses of climate-model output, which suggest that, for every 1-Sv decline in the
¹¹ deep overturning, subpolar sea-surface temperatures cool by 0.2–0.6°C^{1,2,4,5}.

¹² To assess whether observed cooling of the subpolar-Atlantic sea surface and hypothesized
¹³ slowdown of the deep overturning circulation are consistent with my independent determination of
¹⁴ a weakening Florida Current transport, I formulate simple ocean heat budget. The control volume
¹⁵ is taken to be the full-depth Atlantic Ocean north of 27°N. I ignore transports through Bering
¹⁶ Strait, and changes in evaporation minus precipitation over the basin. I also assume that, on these
¹⁷ space- and timescales, local heat storage is negligible to leading order (see below). Thus, the heat
¹⁸ budget is a balance between heat transport divergence due to weakening of the Florida Current and

- 19 deep overturning circulation at 27°N, and turbulent ocean heat gain at the surface due to cooling
 20 sea-surface temperatures in the subpolar region.

21 Following Marshall et al.⁶, the volume-integrated ocean heat transport divergence \mathcal{Q}_{moc} can
 22 be written as,

$$\mathcal{Q}_{moc} \doteq -\rho_o c_o \overline{\Delta T}^z \Psi, \quad (\text{S1})$$

23 where $\rho_o = 1029 \text{ kg m}^{-3}$ is ocean water density, $c_o = 3994 \text{ J kg}^{-1} \text{ °C}^{-1}$ is seawater's specific heat
 24 capacity, Ψ is the overturning streamfunction, and $\overline{\Delta T}^z$ is the temperature difference between the
 25 warm waters in Florida Straits and cool waters below $\sim 1000 \text{ m}$ over the open Atlantic Ocean.

26 The area-integrated turbulent heat gain at the surface \mathcal{Q}_{surf} is expressed,

$$\mathcal{Q}_{surf} \doteq A (Q_E + Q_H), \quad (\text{S2})$$

27 where A is the ocean surface area over which the heat gain occurs, and Q_E and Q_H are latent and
 28 sensible heat fluxes, respectively. After Large and Yeager⁷, the turbulent heat fluxes are written as,

$$Q_E \doteq \Lambda_v \rho_a C_E [q(z_q) - q_{sat}(SST)] |\Delta \vec{U}|, \quad (\text{S3})$$

29 and,

$$Q_H \doteq \rho_a c_a C_H [\theta(z_\theta) - SST] |\Delta \vec{U}|, \quad (\text{S4})$$

30 where $\Lambda_v = 2.5 \times 10^6 \text{ J kg}^{-1}$ is the latent heat of vaporization, $\rho_o = 1.22 \text{ kg m}^{-3}$ is near-surface
 31 density of air, $c_a = 1000.5 \text{ J kg}^{-1} \text{ °C}^{-1}$ is the specific heat capacity of air, $C_H \doteq 0.018\sqrt{C_D}$ (where
 32 C_D is the drag coefficient), $C_E \doteq 0.0346\sqrt{C_D}$, $q(z_q)$ and $\theta(z_\theta)$ are potential air temperature and

³³ specific humidity, respectively, $|\Delta\vec{U}|$ is surface wind speed, SST is sea-surface temperature, and,

$$q_{sat} \doteq \frac{q_1}{\rho_a} \exp(q_2/SST), \quad (S5)$$

³⁴ where $q_1 = 0.98 \times 640380 \text{ kg m}^{-3}$ and $q_2 = -5107.4 \text{ K}$.

³⁵ Using Eq. (S5) to linearize Eq. (S3) about a background sea-surface temperature \overline{SST} gives,

$$Q_E = \Lambda_v \rho_a C_E \left[q(z_q) + \frac{q_2}{\overline{SST}^2} (SST - \overline{SST}) q_{sat}(\overline{SST}) \right] |\Delta\vec{U}|. \quad (S6)$$

³⁶ Equating Q_{moc} and Q_{surf} , making use of Eqs. (S4) and (S6), and solving for SST yields,

$$SST = \frac{-\frac{\rho_o c_o \overline{\Delta T}^z \Psi}{A |\Delta\vec{U}| \rho_a} - c_a C_H \theta(z_\theta) - \Lambda_v C_E q(z_q) + \Lambda_v C_E \frac{q_2}{\overline{SST}^2} q_{sat}(\overline{SST})}{\Lambda_v C_E \frac{q_2}{\overline{SST}^2} q_{sat}(\overline{SST}) - c_a C_H}. \quad (S7)$$

³⁷ Finally, differentiating with respect to Ψ gives,

$$\frac{\partial SST}{\partial \Psi} = -\frac{\rho_o c_o \overline{\Delta T}^z}{A |\Delta\vec{U}| \rho_a \left[\Lambda_v C_E \frac{q_2}{\overline{SST}^2} q_{sat}(\overline{SST}) - c_a C_H \right]}. \quad (S8)$$

³⁸ This expression represents the SST change expected for a unit Ψ change under this simple model.

³⁹ To compute an estimate of $\partial SST / \partial \Psi$, I must choose appropriate values for the remaining

⁴⁰ parameters. I choose $A = 6.7 \times 10^{12} \text{ m}^2$, which represents the area of the shaded (unhatched)

⁴¹ subpolar region of cooling shown in Figure 6. Based on a contemporary ocean state estimate⁸, I

⁴² use $C_D = 0.0011$ so $C_E = 0.0012$ and $C_H = 0.00061$, $6.5\text{--}11.5^\circ\text{C}$ for \overline{SST} , and $7.5\text{--}10.5 \text{ m s}^{-1}$

⁴³ for $|\Delta\vec{U}|$ as reasonable values for the subpolar Atlantic. I also use $10\text{--}15^\circ\text{C}$ as a range for $\overline{\Delta T}^z$

⁴⁴ judging from that same ocean state estimate. These parameter choices lead to an estimated range

⁴⁵ for $\partial SST / \partial \Psi$ of $0.3\text{--}0.6^\circ\text{C Sv}^{-1}$. This range is consistent with values of $0.2\text{--}0.5^\circ\text{C Sv}^{-1}$ found

46 by dividing the observed sea-surface-temperature trends across the subpolar gyre (Figure 6) by the
 47 model's posterior median estimate of the trend in Florida Current transport during the past century
 48 (Supplementary Figure 1b). These values also agree with the values of 0.2–0.6°C Sv⁻¹ published
 49 based on regression analyses of sea-surface temperature and Atlantic overturning streamfunction
 50 from climate models^{1,2,4,5}.

51 **Ignoring local heat storage** In the heat budget, I assumed that local heat storage is negligible.
 52 This assumption is based on a simple scaling argument. Suppose that, in contrast, changes in local
 53 heat storage are in fact important, and have similar magnitude to the change in ocean heat transport
 54 divergence. In this case I can consider the quasi-balance between local storage and advection,

$$\left| V \frac{\delta}{\delta t} \left(\frac{\delta \Theta}{\delta t} \right) \right| \approx \left| \overline{\Delta T}^z \frac{\delta \Psi}{\delta t} \right|, \quad (\text{S9})$$

55 where $\delta \Theta / \delta t$ is the rate of change in ocean temperature averaged over the control region, V is the
 56 volume of the control region, and δt is a time increment. That is, $\delta (\delta \Theta / \delta t)$ is the change in local
 57 ocean heat storage rate required to balance the change in heat transport convergence or divergence
 58 due to a trend in Ψ over the study period. Rearranging to solve for $\delta (\delta \Theta / \delta t)$ gives,

$$\left| \delta \left(\frac{\delta \Theta}{\delta t} \right) \right| \approx \left| \frac{1}{V} \overline{\Delta T}^z \frac{\delta \Psi}{\delta t} \delta t \right|, \quad (\text{S10})$$

59 I take $\overline{\Delta T}^z = 10\text{--}15^\circ\text{C}$ as before, $\delta \Psi / \delta t = 1.7 \text{ Sv century}^{-1}$ (the magnitude of the posterior
 60 median model estimate of the centennial trend in Florida Current transport), and $\delta t = 100 \text{ y}$. Now,
 61 if $V = 6.6\text{--}8.7 \times 10^{16} \text{ m}^3$ (the volume of the full-depth North Atlantic north of 27°N, depending
 62 on whether marginal seas are included), then $\delta (\delta \Theta / \delta t) \sim 0.6\text{--}1.2^\circ\text{C century}^{-1}$. In other words,
 63 for a change in the local heat storage rate to be comparable to the change in ocean heat transport

64 divergence, there would need to be a change in centennial temperature trends averaged over the
65 full-depth control region of this magnitude. If, instead, I take $V = 1.2\text{--}1.9 \times 10^{16} \text{ m}^3$ (the volume
66 of the top 700 m in the northern North Atlantic), then the required change in centennial temperature
67 trends becomes $\delta(\delta\Theta/\delta t) \sim 2.8\text{--}6.6^\circ\text{C century}^{-1}$. Such magnitudes are substantially larger than
68 estimated changes in large-scale temperature trends in the Atlantic over the 20th century compared
69 to previous centuries⁹. So, I conclude that local heat storage, while possibly making higher-order
70 contributions to the budget, can be neglected in this lowest-order-of-magnitude exercise.

71 S2 Residual analysis

72 Various residual terms appear in the Bayesian model equations (see Methods). When building
73 the algorithm, I made certain assumptions regarding the spatial and temporal structures of these
74 residuals. To test whether these assumptions are appropriate given the data, I undertake a residual
75 analysis, using the model equations to solve for the sea-level innovations e_k , tide-gauge errors d_k ,
76 transport innovations w_k , cable-data errors u , tide-gauge error trends a and tide-gauge data bias ℓ .

77 I made the assumption that e_k , d_k , w_k , and u behave as iid temporal white noise. If this
78 assumption is reasonable, then the posterior solutions should look random in time. However, if
79 systematic structure is present, it would mean that this assumption is inappropriate, and that the
80 model is misspecified given the data. Time series of posterior e_k and d_k solutions are shown
81 in Supplementary Figure 8a, 8b for an arbitrary target location, while model solutions for w_k
82 and u are shown in Supplementary Figure 8c, 8d. These time series look random in time, and

83 there are no obvious signs of autocorrelation. The amplitudes of e_k , d_k , and w_k variations are
84 consistent with posterior solutions for the respective variance or partial sill parameters σ^2 , δ^2 , and
85 ω^2 (Supplementary Table 6), and the magnitude of fluctuations in u is in keeping with the prior
86 error variances placed on the submarine-cable data.

87 To be more thorough, I compute sample autocorrelation coefficients directly from the poste-
88 rior solutions for e_k , d_k , w_k , and u across all space and time points. I compare those values to the
89 autocorrelation coefficients expected theoretically for temporal white noise, given the same num-
90 ber of time steps. Supplementary Figure 9 compares the empirical and theoretical autocorrelation
91 coefficients for time lags between 1 and 20 y. Values calculated empirically from the posterior
92 solutions are consistent with the theoretically expected values. More quantitatively, 96%, 95%,
93 93%, and 95% of empirical autocorrelation coefficients computed respectively from e_k , d_k , w_k ,
94 and u are captured by the theoretical 95% confidence intervals.

95 In addition to being random in time, e_k and d_k are supposed to have spatially invariant
96 amplitudes. In Supplementary Figure 10, I map median estimates of standard deviations computed
97 empirically from the posterior model solutions of e_k and d_k at each tide-gauge location. While
98 there is some higher-order spatial variation, these values are to lowest order fairly uniform and
99 constant in space, and very similar to the posterior estimates of the partial sill σ^2 and variance
100 parameter δ^2 (Supplementary Table 6).

101 Motivated by past studies^{10,11}, I assume that e_k is spatially structured, such that there is co-
102 variance between sites along the Caribbean, Central America, and South America, and between

103 sites on the southeastern USA, but no covariance between these two broad regions. These as-
104 sumptions are reflected in the block structure of the theoretical covariance matrix Σ shown in
105 Supplementary Figure 11b computed from the posterior median solution for the partial sill σ^2
106 (Supplementary Table 6). This theoretical covariance matrix is very similar to the covariance ma-
107 trix determined empirically by comparing all pairs of posterior solutions for e_k (Supplementary
108 Figure 11a). Indeed, the Pearson correlation coefficient between the two matrices in Supplemen-
109 tary Figure 11 is 0.91, and the theoretical covariance matrix explains 82% of the variance in the
110 empirical covariance matrix.

111 Finally, I consider residual spatial fields of the tide-gauge data biases $\ell - \nu\mathbf{1}$ and error trends
112 a . According to the data-level Eq. (5) for the tide gauges, these two vectors should have zero mean,
113 no spatial correlation, and spatial variances of τ^2 and γ^2 , respectively. Supplementary Figure 12
114 facilitates an assessment of these assumptions, showing both posterior solutions for $\ell - \nu\mathbf{1}$ and a
115 as well as the solutions expected for a zero-mean random process given the posterior solutions for
116 τ^2 and γ^2 (Supplementary Table 6). Consistent with model assumptions, these vector fields look
117 fairly random, scattered about zero. The spatial spread in $\ell - \nu\mathbf{1}$ and a appears consistent with the
118 posterior τ^2 and γ^2 solutions. Indeed, 95% of the posterior $\ell - \nu\mathbf{1}$ solutions are captured by the 95%
119 credible intervals predicted for a zero-mean, spatially uncorrelated Gaussian process with variance
120 τ^2 , and similarly 95% of posterior solutions for a fall within the 95% credible interval produced
121 by simulating a zero-mean random normal field with variance γ^2 (Supplementary Figure 12).

122 In conclusion, the design of my Bayesian algorithm is supported by residual analysis, which

123 demonstrates that the model structure is appropriate and warranted given the available data.

124 **S3 Sensitivity of model solutions to input data**

125 Posterior solutions for Florida Current transports presented in the main text are based on the as-
126 simulation of submarine cable data over 1982–2018 with specified standard errors of 0.30–0.35 Sv
127 (see Methods). To quantify how robust or sensitive the solutions are to the duration of the data
128 and the selected standard errors, I perform two additional data assimilation experiments. In the
129 first sensitivity experiment, I double the standard errors on the cable data given to the Bayesian
130 algorithm during 1982–2018. I refer to this experiment as the “double-error” experiment. For
131 clarity, in this section, I call the Bayesian model solution presented in the main text the “baseline”
132 experiment. In the second sensitivity experiment, I maintain the original standard errors, but I give
133 the Bayesian algorithm cable measurements for the period 2000–2018, withholding data values
134 during 1982–1998. (Due to an outage in the observing system, no data are available for 1999.) I
135 call this experiment the “half-data” experiment.

136 Salient features of the two sensitivity experiments are summarized alongside the baseline ex-
137 periment in Supplementary Figure S13. Baseline and double-error solutions are, in many respects,
138 very similar. For example, time series of Florida Current transport, transport trend over 1909–2018,
139 and regression coefficient between transport and sea-level difference across the Florida Straits from
140 these two experiments are nearly the same (cf. blue and orange in Supplementary Figure S13). One
141 difference is that the widths of the posterior 95% credible intervals on the transport during 1982–

¹⁴² 2018 (i.e., the period when transport observations are available) are about twice as large in the
¹⁴³ double-error experiment compared to the baseline experiment (Supplementary Figure S13a). This
¹⁴⁴ is consistent with the larger standard errors placed on the data in the former experiment. In sum,
¹⁴⁵ I conclude that model solutions are generally quantitatively insensitive to reasonable alternative
¹⁴⁶ specifications of the standard error on the cable transport measurements.

¹⁴⁷ Solutions from the half-data experiment (yellow in Supplementary Figure S13) show simi-
¹⁴⁸ larities to the other two solutions, but can show larger uncertainty. This is unsurprising, since the
¹⁴⁹ half-data experiment has fewer data constraints. For example, whereas the posterior 95% credible
¹⁵⁰ intervals on the 110-y transport trend are -1.7 ± 3.7 and -1.6 ± 3.9 Sv century⁻¹ in the baseline
¹⁵¹ and double-error experiments, in the half-data experiment it is -2.3 ± 6.9 Sv century⁻¹. The fact
¹⁵² that uncertainties from the double-error experiment are smaller than from the half-data experiment
¹⁵³ suggests that having more observations with larger errors is more informative for constraining the
¹⁵⁴ transport history than having fewer observations that have smaller errors. Importantly, although the
¹⁵⁵ trend from the half-data experiment is more uncertain in an absolute sense, the sign of the trend is
¹⁵⁶ similarly determined in all three experiments. I find that 82%, 80%, and 77% of trend solutions in
¹⁵⁷ the baseline, double-error, and half-data experiments are negative (Supplementary Figure S13b).
¹⁵⁸ That is, all three experiments suggest that Florida Current transport probably declined over the past
¹⁵⁹ century. Thus, I reason that the main findings in this study are qualitatively robust to reasonable
¹⁶⁰ alternative choices for the duration of the transport data assimilated into the Bayesian algorithm.

161 **S4 Synthetic data experiments**

162 In the half-data experiment, $\sim 90\%$ of the observed but withheld Florida Current transport values
163 during 1982–1999 fall within the pointwise posterior 95% credible intervals on the transport. This
164 suggests that the uncertainties estimated by the Bayesian algorithm are reasonable. To more thor-
165oughly evaluate the meaningfulness of the posterior solutions generated by the Bayesian algorithm,
166 I perform a number of synthetic data (or pseudo-proxy) experiments. In these experiments, I take a
167 set of known processes and corrupt them to look like the observations, and I then apply the model
168 to these corrupted process values. By comparing the posterior solutions to the known (withheld)
169 values, I can quantify the accuracy and precision of the error bars furnished by the model (e.g., are
170 $\sim 95\%$ of the true values actually captured by the posterior 95% credible intervals?).

171 **First experiment—perfect model** I run a perfect model experiment. I choose, from the ensemble
172 of posterior model solutions presented in the main text, the array of scalar parameter solutions
173 $(\bar{T}, r, \sigma^2, \dots)$ from the ensemble member that minimizes the Mahalanobis distance to the mean
174 parameter array. Using these scalar parameter values, I simulate synthetic versions of the sea-
175 level and transport processes based on the process-level equations. Using the data-level equations,
176 I generate synthetic tide-gauge and submarine-cable data by adding noise, bias, and gaps to the
177 simulated processes, as in the real world, and I apply the Bayesian model to these synthetic data.

178 Results are summarized in Supplementary Table 7 and Supplementary Figure 14. For 13
179 out of the 14 scalar parameters, or $\sim 93\%$, the true value is captured by the corresponding 95%
180 posterior credible interval from the model (Supplementary Table 7). Considering vector fields, I

181 find that 100%, 98%, and 100% of the true values for regional sea-level trends b , tide-gauge biases
182 ℓ , and tide-gauge error trends a respectively fall within the corresponding pointwise posterior
183 95% credible intervals (not shown). In terms of the processes, 98% of the true sea-level values and
184 99% of true transport values fall within the estimated pointwise 95% credible intervals, and the
185 true transport time series is entirely encompassed by the pathwise 95% posterior credible intervals
186 (e.g., Supplementary Figure 14). Together, these results show that the model performs well, and
187 that the posterior credible intervals are meaningful, if slightly conservative, roughly capturing the
188 correct fraction of true process and parameter values.

189 **Second experiment—more realistic case** The first synthetic data experiment is informative, show-
190 ing that the processes and parameters are identifiable given incomplete, noisy, biased data. It is
191 also potentially idealistic, since the model is perfectly specified. The equations governing the
192 spatiotemporal evolution of the processes, and the relationship between the observations and the
193 processes were known perfectly, and the task was to infer the uncertain values of the processes and
194 parameters appearing in those equations. While residual analysis suggests that they are appropri-
195 ate given the data, the model equations probably represent a simplification of the complex, myriad
196 oceanographic and geophysical processes contributing to changes in sea level and transport, and
197 their correspondence to observations in the real world. While some degree of model misspecifica-
198 tion is inevitable, the salient question is whether the model is robust to misspecification and still
199 provides meaningful posterior estimates.

200 So, I perform a second synthetic data experiment. Rather than use the process equations to

201 simulate sea level and transport, I bring together output from more complex physical models. I
202 begin with ocean dynamics. I take 110 y of monthly Florida Current transport near 27°N, and
203 sea level from each of the model grid cells nearest to the 46 tide gauges from version 2.2.4 of the
204 Simple Ocean Data Assimilation (SODA) product¹³. This version of SODA represents a solution
205 to an ocean general circulation model forced at the surface by an atmospheric reanalysis over the
206 period 1871–2010 (I use the past 110 y of output covering from 1901 to 2010). The model has
207 moderate spatial resolution, with 40 vertical levels and a native $0.25^\circ \times 0.40^\circ$ horizontal grid in
208 longitude and latitude. A version of the solution, which was interpolated onto a regular $0.5^\circ \times 0.5^\circ$
209 horizontal grid, was downloaded from the Asia-Pacific Data-Research Center (APDRC) of the
210 University of Hawai'i School of Ocean and Earth Science and Technology. After downloading I
211 removed the monthly time series of global-mean sea level and computed annual means from the
212 resulting monthly sea-level values.

213 The SODA solution represents a tradeoff between spatial resolution and temporal cover-
214 age. Coupled climate models are available that cover a comparable or longer time period¹⁴, but
215 most publicly available solutions have coarser horizontal resolution (nominally $\sim 1^\circ$ in longitude
216 and latitude), and may not faithfully represent the Florida Current and coastal sea level. While
217 much higher-resolution ocean models are available¹⁵ that more accurately portray the complexity
218 of Florida Current transport and coastal sea level, these model runs are typically short, and do not
219 span the centennial timescales of primary interest here. Thus, while it has its deficiencies (see
220 below), SODA is perhaps one of the best-suited ocean models for my purposes. For example,
221 Chepurin et al.¹⁶ show that version 2.2.4 of SODA simulates interannual-to-decadal variations in

222 coastal sea level along the eastern USA and parts of the Caribbean reasonably well over 1950–
223 2011.

224 I superimpose static sea-level effects on the dynamic sea-level fields from SODA. I add a
225 yearly time series of global-mean sea level due to ocean warming and thermal expansion over
226 1901–2010 from the Version 4 of the Community Climate System Model¹⁷ (downloaded from
227 the Woods Hole Oceanographic Institution’s Community Storage Server). I also include, at each
228 tide-gauge location, an estimate of the trend in relative sea level due to the combined effects of
229 ongoing glacial isostatic adjustment from Peltier et al.¹² (downloaded from the PSMSL) along with
230 twentieth-century melting of mountain glaciers and ice sheets due to Hamlington et al.¹⁸ (courtesy
231 of S. Adhikari, Jet Propulsion Laboratory). Finally, I add time series of a random-in-time but
232 correlated-in-space process with zero mean and temporal variance of $\sim (1 \text{ cm})^2$ to simulate sea-
233 level changes due to the inverted barometer effect linked with the North Atlantic Oscillation¹⁹.

234 I apply the data-level equations to these transport and sea-level values, incorporating noise
235 and bias, and imparting data gaps so that the synthetic tide-gauge and submarine-cable data are
236 only available when and where the true observations are available. These synthetic datasets are
237 subsequently fed into the Bayesian model algorithm. The results of this second synthetic data
238 experiment are summarized in Supplementary Table 8 and Supplementary Figure 15. In this case,
239 only four scalar parameters (those appearing in the data-level equations) are known perfectly. For
240 three out of these four parameters, or 75%, the true value is captured by the 95% posterior credible
241 intervals from the model (Supplementary Table 8). For one parameter, δ^2 , the tide-gauge data error

242 variance, the Bayesian model slightly underestimates the true value. Considering the process time
243 series, I find that 81% of the true transport values and 95% of the true sea-level values are captured
244 by the pointwise 95% posterior credible intervals produced by the Bayesian model, and that, as in
245 the previous experiment, the full time series of the true transport is totally captured by the pathwise
246 95% posterior credible interval (Supplementary Figure 15).

247 It is worth noting that the posterior solution for α , the apparent trend in the transport process
248 Eq. (4), suggests that sea level at Settlement Point on Grand Bahama must have risen 0.2 ± 1.6
249 mm y^{-1} faster than at West Palm Beach near West Palm Beach due to processes unrelated to ocean
250 dynamics. This is consistent with the trend difference of $\sim 0.1 \text{ mm y}^{-1}$ I imposed between these
251 two sites based on model estimates of GIA and contemporary ice melt^{12,18}, demonstrating that the
252 model succeeds in separating static and dynamic sea-level trends.

253 Recall that my Bayesian model assumes that the transfer coefficient ρ between sea level and
254 transport is a fixed constant. To test this assumption, I consider in more detail time series of Florida
255 Current transport and sea-level difference across Florida Straits from SODA. Transport and sea-
256 level difference are highly correlated with one another (Pearson correlation coefficient of ~ 0.9),
257 and a linear regression suggests that transport increases by $\sim 0.9 \text{ Sv}$ for every 1-cm increase in
258 sea level difference, consistent with a visual inspection of the two time series (Supplementary Fig-
259 ure 4a). To study the correspondence as a function of frequency band, I apply admittance and
260 coherence analysis to the model output. Transport and sea-level difference are significantly coher-
261 ent at all accessible periods from 2- to 32-y (Supplementary Figure 4b), in agreement with basic

expectations from geostrophy. Moreover, the transfer function (using sea-level difference as the input and transport as the output) is qualitatively insensitive to frequency band, with similar values found at interannual and multidecadal timescales (Supplementary Figure 4c). Importantly, the Bayesian model posterior estimate for the transfer coefficient ρ is consistent with SODA and overlaps the values obtained from the admittance analysis (Supplementary Figure 4c). This suggests that it is reasonable to assume that there is a constant transfer coefficient between sea-level difference and transport on the timescales of this study, and also that the Bayesian model successfully infers the correct transfer-coefficient value.

Note that the Florida Current transport from SODA is suspicious (Supplementary Figure 15c). Mean transport is ~ 51 Sv, growing from ~ 42 Sv at the beginning of the period to ~ 56 Sv at the end. This value is $\sim 60\%$ larger than the average value observed by submarine cable since 1982, and ~ 10 Sv larger than the largest annual transport value inferred at any time in the original Bayesian model solution discussed in the main text. The striking increase of ~ 14 Sv over the 110-y run is extreme in light of the more subtle trend estimates produced by the original Bayesian model solution (cf. Figure 2a; Supplementary Figure 15c). Although it is imperfect, in that it does not realistically represent the true evolution of the Florida Current over the past century, SODA is nevertheless informative in the present context. For establishing the ability of the Bayesian algorithm to infer the parameters and processes from imperfect data, I do not require that the SODA reproduces observed reality, but rather that it portrays a physically plausible scenario, and that the basic “statistics” (e.g., spatiotemporal covariance structure, relationship between state variables, etc.) are believable.

283 In sum, I conclude that, even in a more complex setting, my Bayesian model performs rea-
284 sonably well, giving uncertainty estimates that roughly capture the correct fraction of true values.

285 **References**

- 286 1. Rahmstorf, S., et al. Exceptional twentieth-century slowdown in Atlantic Ocean overturning
287 circulation, *Nature Clim. Change*, **5**, 475–480 (2015).
- 288 2. Caesar, L., S. Rahmstorf, A. Robinson, G. Feulner, and V. Saba. Observed fingerprint of a
289 weakening Atlantic Ocean overturning circulation, *Nature*, **556**, 191–196 (2018).
- 290 3. Thornalley, D. J. R., et al. Anomalously weak Labrador Sea convection and Atlantic overturning
291 during the past 150 years, *Nature*, **556**, 227–230 (2018).
- 292 4. Drijfhout, S., et al. Is a Decline of AMOC Causing the Warming Hole above the North Atlantic
293 in Observed and Modeled Warming Patterns?, *J. Clim.*, **25**, 8373–8379 (2012).
- 294 5. Roberts, C. D., F. K. Garry, and L. C. Jackson. A Multimodel Study of Sea Surface Temperature
295 and Subsurface Density Fingerprints of the Atlantic Meridional Overturning Circulation, *J.*
296 *Clim.*, **26**, 9155–9174 (2013).
- 297 6. Marshall, J., H. Johnson, and J. Goodman. A Study of the Interaction of the North Atlantic
298 Oscillation with Ocean Circulation, *J. Clim.*, **14**, 1399–1421 (2001).

- 299 7. Large, W. G., and S. G. Yeager. Diurnal to Decadal Global Forcing For Ocean and Sea-Ice
300 Models: The Data Sets and Flux Climatologies, *NCAR Technical Note*, NCAR/TN-460+STR,
301 112 pp.
- 302 8. Forget, G., et al. ECCO version 4: an integrated framework for non-linear inverse modeling and
303 global ocean state estimation, *Geosci. Model Dev.*, **8**, 3071–3104 (2015).
- 304 9. Gebbie, G. Atlantic Warming Since the Little Ice Age, *Oceanography*, **32**, 220–230 (2019).
- 305 10. Thompson, P. R., and G. T. Mitchum. Coherent sea level variability on the North Atlantic west-
306 ern boundary, *J. Geophys. Res. Oceans*, **119**, 5676–5689, doi:10.1002/2014JC009999 (2014).
- 307 11. Zhao, J., and W. Johns. Wind-forced interannual variability of the Atlantic Meridional Over-
308 turning Circulation at 26.5°N, *J. Geophys. Res.-Oceans*, **119**, 2403–2419.
- 309 12. Peltier, W. R., et al. Space geodesy constrains ice age terminal deglaciation: The global ICE-
310 6G_C (VM5a) model, *J. Geophys. Res.-Solid Earth*, **120**, 450–487 (2015).
- 311 13. Giese, B. S., and S. Ray. El Niño variability in simple ocean data assimilation (SODA), 1871–
312 2008, *J. Geophys. Res.*, **116**, C02024 (2011).
- 313 14. Taylor, K. E., et al. An Overview of CMIP5 and the Experimental Design, *B. Am. Meteorol.
314 Soc.*, **93**, 485–498 (2012).
- 315 15. Gula, J., M. J. Molemaker, and J. C. McWilliams. Topographic generation of submesoscale
316 centrifugal instability and energy dissipation, *Nat. Commun.*, **7**, 12811 (2016).

- 317 16. Chepurin, G. A., J. A. Carton, and E. Leuliette. Sea level in ocean reanalyses and tide gauges,
- 318 *J. Geophys. Res.-Oceans*, **119**, 147–155 (2014).
- 319 17. Gent, P. R., et al. The Community Climate System Model Version 4 , *J. Clim.*, **24**, 4973–4991
- 320 (2011).
- 321 18. Hamlington, B. D., et al. Observation-Driven Estimation of the Spatial Variability of 20th
- 322 Century Sea Level Rise, *J. Geophys. Res.-Oceans*, **123**, 2129–2140 (2018).
- 323 19. Piecuch, C. G., and R. M. Ponte. Inverted barometer contributions to recent sea level changes
- 324 along the northeast coast of North America, *Geophys. Res. Lett.*, **42**, 5918–5925 (2015).
- 325 20. Santamaría-Gómez, A., et al. Uncertainty of the 20th century sea-level rise due to vertical land
- 326 motion errors, *Earth Planet. Sci. Lett.*, **473**, 24–32 (2017).
- 327 21. Love, R., et al. The contribution of glacial isostatic adjustment to projections of sea-level
- 328 change along the Atlantic and Gulf coasts of North America, *Earth's Future*, **4**, 440–464 (2016).
- 329 22. Khan, N. S., et al. Drivers of Holocene sea-level change in the Caribbean, *Quaternary Sci.*
- 330 *Rev.*, **155**, 13–36 (2017).
- 331 23. Gelman, A, and D. B. Rubin. Inference from iterative simulation using multiple sequences,
- 332 *Stat. Sci.*, **7**, 457–472 (1992).
- 333 24. Birol, F., et al. Coastal applications from nadir altimetry: Example of the X-TRACK regional
- 334 products, *Adv. Space Res.*, **59**, 936–953 (2017).

Region	Site	Lon ($^{\circ}$ W)	Lat ($^{\circ}$ N)	Duration (years)	Rate (mm y^{-1})	Error (mm y^{-1})
Florida	AOML	80.1622	25.7347	6.37	0.42	0.74
Florida	CCV6	80.5455	28.4600	6.93	-2.80	0.74
Florida	MIA3	80.1602	25.7328	11.00	-0.17	0.80
Bahamas	EXU0	75.8734	23.5640	6.50	-1.70	0.74
Bahamas	NAS0	77.4623	25.0525	6.51	-2.03	2.42

Table S1: Summary of GPS data from Version 6a of the dataset from Université de la Rochelle²⁰ used to estimate the difference in static sea-level rate across Florida Straits due to differential land motion quoted in the main text. Duration is the length of the data record. Error is twice the formal standard error provided with the dataset. Assuming errors are independent, the average rate across the two Bahamas sites is $-1.87 \pm 1.27 \text{ mm y}^{-1}$ and the average rate across the three southeastern Florida sites is $-0.85 \pm 0.44 \text{ mm y}^{-1}$. The difference between the former and latter average values is $-1.02 \pm 1.34 \text{ mm y}^{-1}$, which represents the rate of differential vertical land motion across Florida Straits quoted in the main text. Multiplying by -1 to convert from the land-motion frame to the sea-level frame gives the value of $1.0 \pm 1.3 \text{ mm y}^{-1}$ quoted in the main text.

Region	Site	Reference	Lon (°W)	Lat (°N)	age (y BP)	sea level (m)
Florida	Florida Bay	Love et al. ²¹	80.6	25	1260 ± 275	-1.34 ± 1.27
					890 ± 290	-0.83 ± 1.39
					400 ± 335	-1.00 ± 1.26
Florida	Bear Point	Love et al. ²¹	80.3	27.4	1930 ± 350	-0.93 ± 1.45
					1380 ± 225	-1.13 ± 1.45
					1120 ± 215	-0.83 ± 1.45
Bahamas	Acklins Island	Khan et al. ²²	73.9	22.5	1048 ± 490	-1.64 ± 1.14
					698 ± 392	-1.23 ± 1.26
					398 ± 500	-1.08 ± 1.22
					242 ± 484	-0.97 ± 1.18

Table S2: Proxy sea-level index points from southeastern Florida and the Bahamas used to estimate the difference in the rate of late-Holocene sea-level change across Florida Straits quoted in the main text. Latitudes and longitudes have been rounded to the nearest tenth of a degree. The “y BP” abbreviation stands for years before present, where present is 1950. The \pm values are twice the standard errors on the age and sea-level values provided in the given references. Using ordinary least squares to fit a trend line to the index points at each site, and ignoring age and sea-level uncertainty, I compute trends of 0.36 ± 0.97 , 0.05 ± 0.73 and 0.81 ± 0.22 mm y^{-1} at Florida Bay, Bear Point, and Acklins Island, respectively, where \pm is twice the formal standard error furnished by ordinary least squares assuming independent data. The average of the two trends from southeastern Florida is thus 0.20 ± 0.61 mm y^{-1} and so the difference between the Bahamas and southeastern Florida is 0.6 ± 0.6 , which is the value quoted in the main text.

No.	Location	Lon (°E)	Lat (°N)	Timespan (Completeness)	Coast
1	Cristóbal	-79.9167	9.35	1909–1979 (100%)	904
2	Puerto Limón	-83.0333	10	1949–1968 (90%)	906
3	Cartagena	-75.55	10.4	1949–1992 (68%)	902
4	Riohacha	-72.9167	11.55	1953–1969 (82%)	902
5	Fort-de-France II	-61.0632	14.6015	2006–2017 (100%)	912
6	Santo Tomás de Castilla	-88.6167	15.7	1965–1980 (75%)	916
7	Puerto Cortés	-87.95	15.8333	1948–1968 (100%)	908
8	Puerto Castilla	-86.0333	16.0167	1956–1968 (100%)	908
9	Lime Tree Bay	-64.7533	17.6933	1986–2015 (80%)	939
10	Port Royal	-76.85	17.9333	1955–1969 (100%)	932
11	Magueyes Island	-67.045	17.97	1955–2016 (90%)	938
12	Barahona	-71.0833	18.2	1955–1969 (67%)	936
13	Charlotte Amalie	-64.92	18.335	1976–2016 (61%)	939
14	San Juan	-66.115	18.4583	1963–2016 (81%)	938
15	Port-au-Prince	-72.35	18.5667	1950–1961 (100%)	934
16	South Sound	-81.3833	19.2667	1976–1993 (89%)	931
17	North Sound	-81.3167	19.3	1976–1996 (86%)	931
18	Puerto Plata	-70.7	19.8167	1950–1969 (70%)	936
19	Cabo Cruz	-77.7333	19.8333	1993–2017 (76%)	930
20	Guantanamo Bay	-75.1467	19.9067	1938–1971 (85%)	930
21	Gibara	-76.125	21.1083	1976–2016 (100%)	930
22	Nuevitas Punta Practico	-77.1095	21.5913	1992–2017 (35%)	930
23	Casilda II	-79.9917	21.7533	1984–2014 (48%)	930
24	Cabo de San Antonio	-84.9	21.9	1973–2017 (60%)	930
25	Isabela de Sagua	-80.0167	22.9333	2000–2016 (71%)	930
26	Key West	-81.8067	24.555	1913–2018 (97%)	940
27	Vaca Key	-81.105	24.7117	1990–2017 (79%)	940
28	Key Colony Beach	-81.0167	24.7183	1978–1994 (71%)	940
29	Virginia Key	-80.1617	25.73	1995–2017 (87%)	960
30	Miami Beach	-80.1317	25.7683	1932–1980 (92%)	960
31	Naples	-81.8067	26.1317	1966–2017 (83%)	940
32	West Palm Beach	-80.0333	26.6117	1974–2017 (36%)	960
33	Settlement Point	-78.9833	26.6833	2005–2015 (82%)	941
34	Settlement Point	-78.9967	26.71	1986–2000 (67%)	941
35	Trident Pier	-80.5917	28.415	1995–2017 (91%)	960
36	Daytona Beach Shores	-80.9633	29.1467	1967–1983 (71%)	960
37	Daytona Beach	-81	29.2333	1925–1969 (51%)	960
38	Jacksonville	-81.6167	30.35	1954–1967 (100%)	960
39	Mayport	-81.4317	30.3933	1929–1999 (99%)	960
40	Mayport	-81.4283	30.3983	2001–2017 (94%)	960
41	Fernandina Beach	-81.465	30.6717	1909–2018 (78%)	960
42	Fort Pulaski	-80.9017	32.0333	1935–2018 (95%)	960
43	Charleston	-79.925	32.7817	1922–2018 (100%)	960
44	Springmaid Pier	-78.9183	33.655	1978–2017 (60%)	960
45	Myrtle Beach	-78.885	33.6833	1958–1977 (55%)	960
46	Wilmington	-77.9533	34.2267	1936–2018 (95%)	960

Table S3: Descriptions of tide-gauge sea-level records used in this study. “Completeness” is the percentage of timespan during which data are available. “Coast” number is the code used by the PSMSL to indicate the country and coastline of measurement.

Parameter	Description
η_0	Sea-level initial condition
η_k	Sea-level values at time t_k
\bar{T}	Transport time-mean value
T_k	Transport value at time t_k
b	Spatial vector of regional trends in sea level
a	Spatial vector of local trends in sea level
ℓ	Spatial vector of tide-gauge biases
r	AR(1) coefficient of sea level
μ	Mean value of regional trends in sea level
ν	Mean value of tide-gauge biases
ρ	Transport change per unit sea-level difference
α	Transport trend correction
π^2	Partial sill of regional trends in sea level
σ^2	Partial sill of sea-level innovations
δ^2	Tide-gauge error variance
τ^2	Spatial variance in observational biases
γ^2	Variance of local trends in sea level
ω^2	Variance of transport noise correction
ϕ	Inverse range of sea-level innovations
λ	Inverse range of regional trends in sea level

Table S4: Descriptions of model processes and parameters.

Parameter	Prior Distribution	Hyperparameter Values
η_0	$\mathcal{N}(\tilde{\eta}_{\eta_0}, \tilde{\zeta}_{\eta_0}^2)$	$\tilde{\eta}_{\eta_0} = -0.2 \text{ m}$, $\tilde{\zeta}_{\eta_0}^2 = (7.6 \times 10^{-2} \text{ m})^2$
\bar{T}	$\mathcal{N}(\tilde{\eta}_{\bar{T}}, \tilde{\zeta}_{\bar{T}}^2)$	$\tilde{\eta}_{\bar{T}} = 32 \text{ Sv}$, $\tilde{\zeta}_{\bar{T}}^2 = (5.2 \text{ Sv})^2$
r	$\mathcal{U}(\tilde{u}_r, \tilde{v}_r^2)$	$\tilde{u}_r = 0.0$, $\tilde{v}_r^2 = 1.0$
μ	$\mathcal{N}(\tilde{\eta}_\mu, \tilde{\zeta}_\mu^2)$	$\tilde{\eta}_\mu = 3.4 \times 10^{-3} \text{ m y}^{-1}$, $\tilde{\zeta}_\mu^2 = (2.7 \times 10^{-2} \text{ m y}^{-1})^2$
ν	$\mathcal{N}(\tilde{\eta}_\nu, \tilde{\zeta}_\nu^2)$	$\tilde{\eta}_\nu = 7.0 \text{ m}$, $\tilde{\zeta}_\nu^2 = (0.6 \text{ m})^2$
ρ	$\mathcal{N}(\tilde{\eta}_\rho, \tilde{\zeta}_\rho^2)$	$\tilde{\eta}_\rho = 0.0 \text{ Sv m}^{-1}$, $\tilde{\zeta}_\rho^2 = (190 \text{ Sv m}^{-1})^2$
α	$\mathcal{N}(\tilde{\eta}_\alpha, \tilde{\zeta}_\alpha^2)$	$\tilde{\eta}_\alpha = 0.0 \text{ Sv y}^{-1}$, $\tilde{\zeta}_\alpha^2 = (0.3 \text{ Sv y}^{-1})^2$
π^2	$\mathcal{IG}(\tilde{\xi}_{\pi^2}, \tilde{\chi}_{\pi^2}^2)$	$\tilde{\xi}_{\pi^2} = 0.5$, $\tilde{\chi}_{\pi^2}^2 = (1.9 \times 10^{-3} \text{ m y}^{-1})^2$
σ^2	$\mathcal{IG}(\tilde{\xi}_{\sigma^2}, \tilde{\chi}_{\sigma^2}^2)$	$\tilde{\xi}_{\sigma^2} = 0.5$, $\tilde{\chi}_{\sigma^2}^2 = (1.8 \times 10^{-2} \text{ m})^2$
δ^2	$\mathcal{IG}(\tilde{\xi}_{\delta^2}, \tilde{\chi}_{\delta^2}^2)$	$\tilde{\xi}_{\delta^2} = 0.5$, $\tilde{\chi}_{\delta^2}^2 = (7.1 \times 10^{-3} \text{ m})^2$
τ^2	$\mathcal{IG}(\tilde{\xi}_{\tau^2}, \tilde{\chi}_{\tau^2}^2)$	$\tilde{\xi}_{\tau^2} = 0.5$, $\tilde{\chi}_{\tau^2}^2 = (8.5 \times 10^{-2} \text{ m})^2$
γ^2	$\mathcal{IG}(\tilde{\xi}_{\gamma^2}, \tilde{\chi}_{\gamma^2}^2)$	$\tilde{\xi}_{\gamma^2} = 0.5$, $\tilde{\chi}_{\gamma^2}^2 = (7.1 \times 10^{-4} \text{ m y}^{-1})^2$
ω^2	$\mathcal{IG}(\tilde{\xi}_{\omega^2}, \tilde{\chi}_{\omega^2}^2)$	$\tilde{\xi}_{\omega^2} = 0.5$, $\tilde{\chi}_{\omega^2}^2 = (0.7 \text{ Sv})^2$
ϕ	$\mathcal{LN}(\tilde{\eta}_\phi, \tilde{\zeta}_\phi^2)$	$\tilde{\eta}_\phi = -7.0 \log \text{ km}^{-1}$, $\tilde{\zeta}_\phi^2 = (2.2 \log \text{ km}^{-1})^2$
λ	$\mathcal{LN}(\tilde{\eta}_\lambda, \tilde{\zeta}_\lambda^2)$	$\tilde{\eta}_\lambda = -6.9 \log \text{ km}^{-1}$, $\tilde{\zeta}_\lambda^2 = (0.4 \log \text{ km}^{-1})^2$

Table S5: Prior distributions and hyperparameters. Hyperparameters are denoted with tildes to distinguish them from the other (uncertain) model parameters. The scripts are: \mathcal{N} normal (or multivariate normal) distribution with mean $\tilde{\eta}$ and variance $\tilde{\zeta}^2$; \mathcal{U} uniform distribution with lower bound \tilde{u} and upper bound \tilde{v} ; \mathcal{IG} inverse-gamma distribution with shape ξ and scale χ ; \mathcal{LN} log-normal distribution with “mean” $\tilde{\eta}$ and “variance” $\tilde{\zeta}^2$.

Parameter	Units	\hat{R}	Median Value	95% CI	Width Ratio
T	Sv	1.001	32.6317	[31.2047, 34.0538]	0.13837
α	Sv y ⁻¹	1.0007	-0.013584	[-0.054013, 0.0293]	0.085205
r	—	1.0066	0.55246	[0.47413, 0.63057]	0.16441
$\mu (\times 10^3)$	m y ⁻¹	1.0007	2.6671	[1.1105, 4.2612]	0.028929
ν	m	0.99976	6.9845	[6.9619, 7.0065]	0.018982
ρ	Sv m ⁻¹	0.9996	21.3501	[10.4544, 32.4271]	0.029465
$\pi^2 (\times 10^6)$	(m y ⁻¹) ²	1.0001	$(1.1673)^2$	$[(0.75971)^2, (1.9104)^2]$	0.00056614
$\sigma^2 (\times 10^6)$	m ²	1.0019	$(26.2588)^2$	$[(24.4292)^2, (28.3339)^2]$	0.00024641
$\delta^2 (\times 10^6)$	m ²	0.99995	$(8.3539)^2$	$[(7.3177)^2, (9.4754)^2]$	0.00037666
$\tau^2 (\times 10^6)$	m ²	0.99973	$(66.9832)^2$	$[(54.0808)^2, (85.3079)^2]$	0.00040194
$\gamma^2 (\times 10^6)$	(m y ⁻¹) ²	0.99995	$(0.6992)^2$	$[(0.40244)^2, (1.1171)^2]$	0.00090338
ω^2	Sv ²	0.9997	$(0.708)^2$	$[(0.4832)^2, (1.0033)^2]$	0.00058865
$\phi (\times 10^3)$	km ⁻¹	1.0025	0.68742	[0.52277, 0.87158]	0.0040641
$\lambda (\times 10^3)$	km ⁻¹	1.0005	0.8429	[0.43847, 1.6407]	0.80349

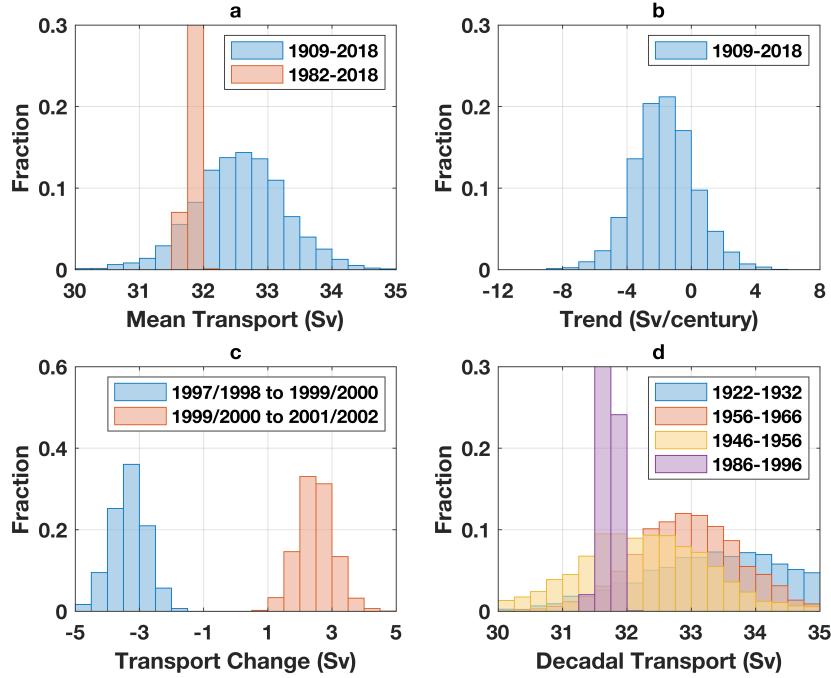
Table S6: Summary of posterior solutions for scalar parameters. The symbol \hat{R} is a convergence monitor of Gelman and Rubin²³, such that values near 1 indicate convergence. Median Value and 95% credible interval (CI) are computed from the ensemble of posterior model solutions. The Width Ratio is defined as ratio of the width of the posterior 95% CI to the prior 95% CI width.

Parameter	Units	Truth	Median Value	95% CI
\bar{T}	Sv	32.8942	32.0523	[30.9524, 33.0873]
α	Sv y ⁻¹	-0.018899	-0.023436	[-0.059135, 0.0090315]
r	—	0.54595	0.53247	[0.46355, 0.60654]
$\mu (\times 10^3)$	m y ⁻¹	2.977	3.1574	[1.2438, 5.1438]
ν	m	6.9876	6.9947	[6.9739, 7.0165]
ρ	Sv m ⁻¹	23.5497	20.974	[14.9067, 27.6991]
$\pi^2 (\times 10^6)$	(m y ⁻¹) ²	(1.078) ²	(1.4473) ²	[(0.94505) ² , (2.2444) ²]
$\sigma^2 (\times 10^6)$	m ²	(26.443) ²	(25.5557) ²	[(23.6732) ² , (27.7207) ²]
$\delta^2 (\times 10^6)$	m ²	(8.7092) ²	(9.2437) ²	[(8.3297) ² , (10.1856) ²]
$\tau^2 (\times 10^6)$	m ²	(67.1828) ²	(66.178) ²	[(54.0051) ² , (83.3185) ²]
$\gamma^2 (\times 10^6)$	(m y ⁻¹) ²	(0.64645) ²	(0.80521) ²	[(0.54918) ² , (1.1481) ²]
ω^2	Sv ²	(0.77083) ²	(0.34671) ²	[(0.23695) ² , (0.51894) ²]
$\phi (\times 10^3)$	km ⁻¹	0.63572	0.60636	[0.46714, 0.78344]
$\lambda (\times 10^3)$	km ⁻¹	0.79168	0.83584	[0.44534, 1.6007]

Table S7: Summary of first synthetic data experiment. Comparison between the true (withheld) parameter values and the posterior model estimates.

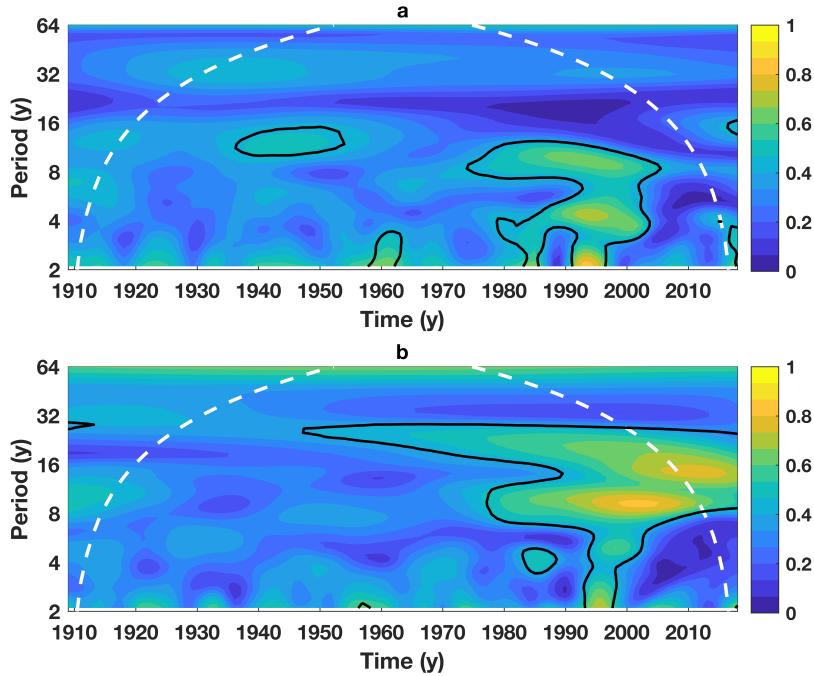
Parameter	Units	True Value	Median Value	95% CI
ν	m	6.9876	6.9707	[6.9506, 6.9918]
$\delta^2 (\times 10^6)$	m^2	$(8.7092)^2$	$(7.2674)^2$	$[(6.4296)^2, (8.1361)^2]$
$\tau^2 (\times 10^6)$	m^2	$(67.1828)^2$	$(62.0712)^2$	$[(50.8668)^2, (78.9978)^2]$
$\gamma^2 (\times 10^6)$	$(m \text{ } y^{-1})^2$	$(0.64645)^2$	$(0.80316)^2$	$[(0.55894)^2, (1.1291)^2]$

Table S8: **Summary of second synthetic data experiment.** Comparison between the true (withheld) parameter values and the posterior model estimates.



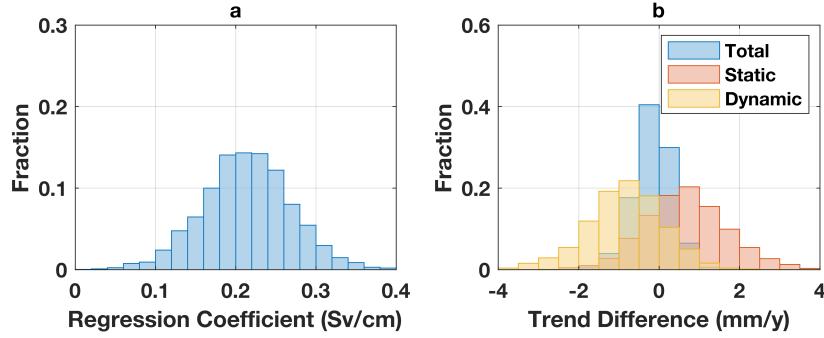
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336 **Figure S1 Some aspects of the posterior solution.** **a**, Blue (orange) is the histogram of
 337 the mean of the transport T in units of Sv over the 1909–2018 study period (1982–2018).
 338 **b**, Histogram of the transport trend $\rho b^T \Delta + \alpha$ over 1909–2018 (Sv century^{-1}). **c**, Blue
 339 (orange) is the histogram of the change in transport T in units of Sv between 1997/1998
 340 to 1999/2000 (1999/2000 to 2001/2002). **d**, Histograms of decadally averaged transport
 341 T in units of Sv: blue 1922–1932; orange 1956–1966; yellow 1946–1956; and purple
 342 1986–1996. See Supplementary Table 4 for descriptions of symbols.



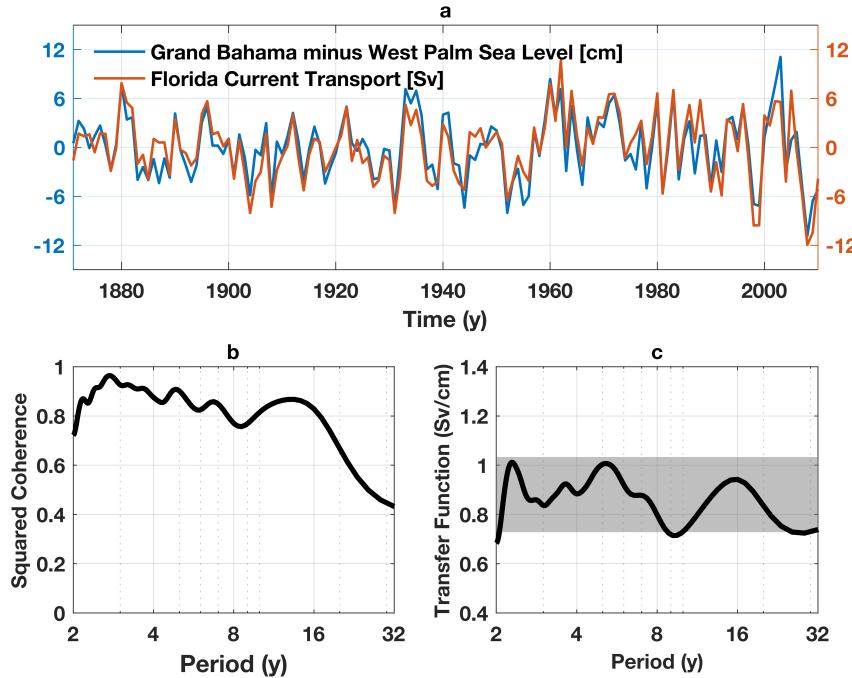
343

344 **Figure S2 Wavelet coherences.** Magnitude squared wavelet coherence between Florida
 345 Current transport T and **a**, North Atlantic Oscillation and **b**, Atlantic Multidecadal Variabil-
 346 ity. Values are computed as follows. For each ensemble member, the wavelet coherence
 347 is computed between the transport solution and the climate index. For the same ensemble
 348 member, two random time series are generated, which have identical Fourier amplitudes
 349 to the transport solution and climate index, but randomized phases, and the wavelet co-
 350 herence between the random time series is computed. Shaded colors represent medians
 351 of the set of wavelet-coherence values computed between all transport solutions and
 352 the given climate index. Black contouring indicates where 68% of wavelet coherences
 353 computed between transport solutions and the climate index exceed the value calculated
 354 between the pairs of random time series.



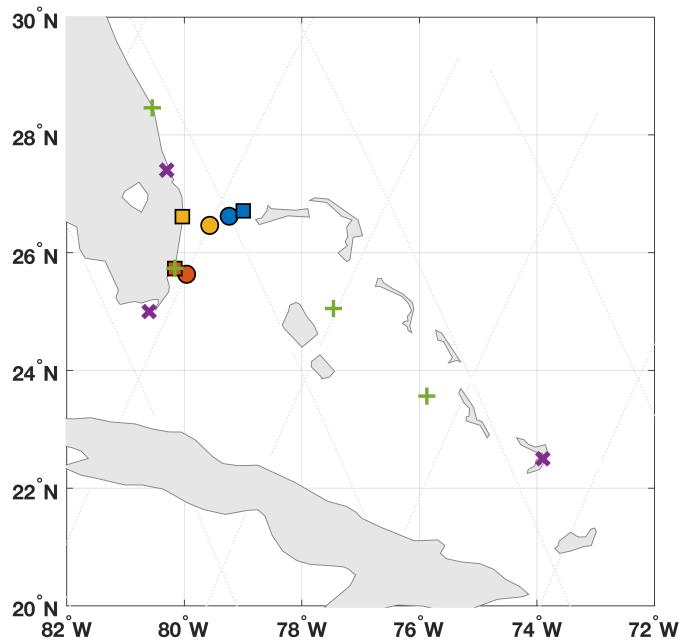
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356 **Figure S3 More aspects of the posterior solution.** a, Histogram of posterior solutions for
 357 the regression coefficient ρ (Sv cm^{-1}) between sea-level difference across Florida Straits
 358 and Florida Current transport. b, Histogram of posterior solutions for the total (blue), static
 359 (orange), and dynamic (yellow) trends in sea-level difference across Florida Straits, which
 360 are computed respectively as $b^T \Delta$, $-\alpha/\rho$, and $b^T \Delta + \alpha/\rho$ (mm y^{-1}) (cf. Methods). See
 361 Supplementary Table 4 for descriptions of symbols.



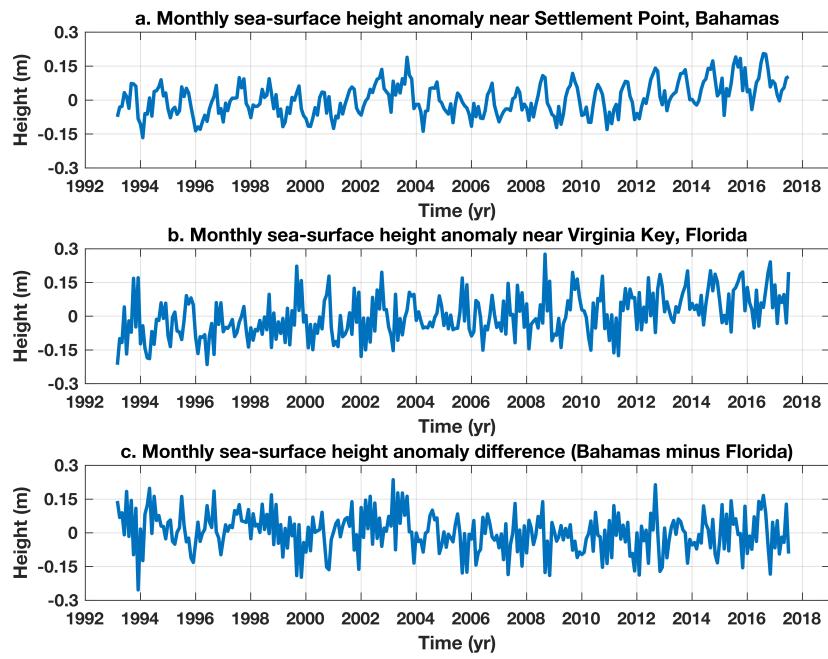
362

363 **Figure S4 Sea level and transport from SODA.** a, Blue (orange) is the annual time se-
 364 ries of sea-level difference across Florida Straits between Grand Bahama and West Palm
 365 Beach (Florida Current transport) during 1871–2010 from the SODA ocean model solu-
 366 tion. Note that both time series have been detrended. b, Black line is magnitude-squared
 367 coherence between sea-level difference and transport the first 128 y of the SODA model
 368 solution (1871–1998). All values are statistically significant at the 95% confidence level
 369 based on comparison against calculations with synthetic time series. c, Black line is am-
 370 plitude of the transfer function from an admittance calculation using sea-level difference
 371 as the input and transport as the output. Gray shading is the 95% posterior credible in-
 372 terval on the transfer coefficient ρ from a synthetic data experiment based on the SODA
 373 model solution (see Supplementary Information). Admittance and coherence calculations
 374 are based on Welch's method using a window length of 32 and 50% overlap.



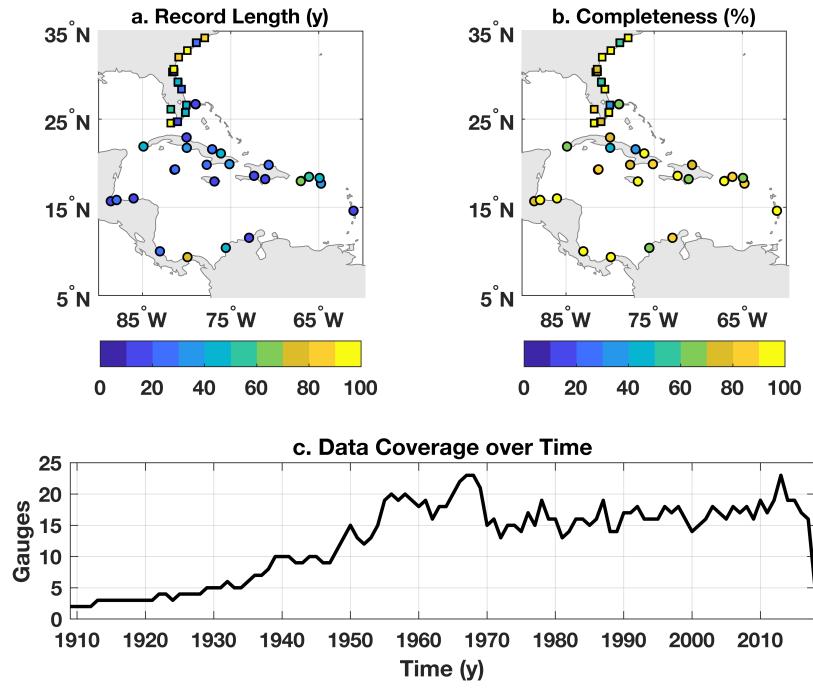
375

376 **Figure S5 Locations of ancillary observational assets.** Shaded squares are tide-gauge
 377 locations (blue is Settlement Point; orange is Virginia Key; yellow is West Palm Beach).
 378 Shaded circles are the along-track satellite-altimeter data points that are nearest the cor-
 379 responding tide gauge. Light gray criss-crossing marks ascending and descending al-
 380 timeter tracks. Green + symbols denote locations of GPS stations (cf. Supplementary
 381 Table 1). Purple × symbols are the locations of proxy sea-level indicators (cf. Supple-
 382 mentary Table 2).



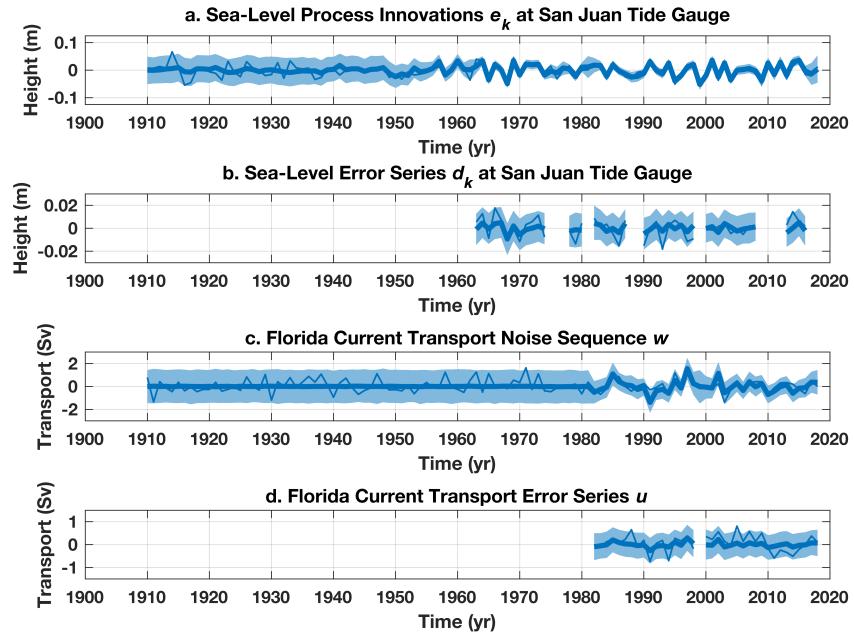
383

384 **Figure S6 Altimetric sea-surface height.** Monthly time series of anomalous sea-surface
 385 height from satellite altimetry near **a**, Settlement Point, Bahamas, **b**, Virginia Key, Florida,
 386 and **c**, the difference between the two time series. Values shown here are calculated by
 387 bin averaging the raw 1-Hz data provided by Birol et al.²⁴ by year and month. A mean
 388 seasonal cycle (annual and semi-annual harmonics) has been removed in each case.
 389 See Supplementary Figure S5 for the locations of the time series.



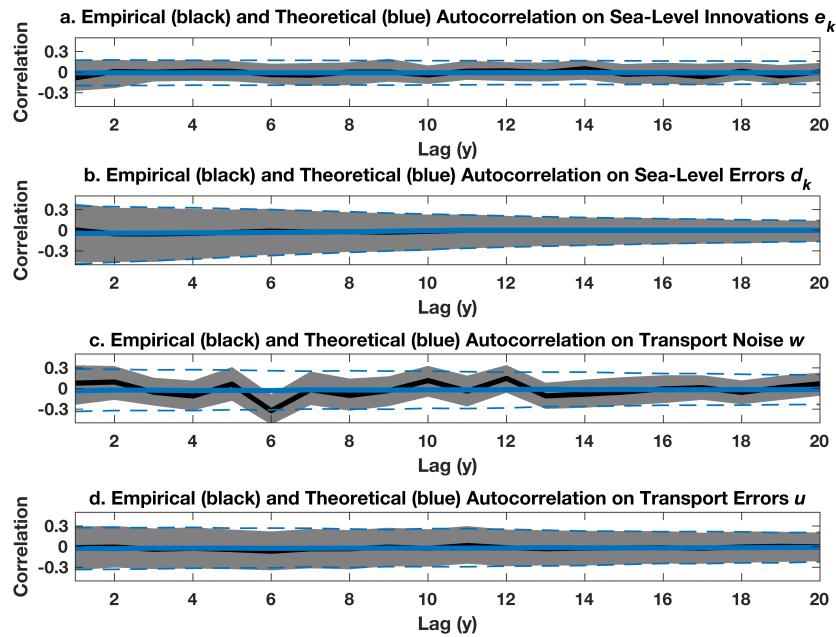
390

391 **Figure S7 Characteristics of tide-gauge data.** **a**, Record length of tide-gauge records
 392 (number of y between the first and last measurements made during the study period).
 393 Yellower (bluer) colors indicate longer (shorter) records. **b**, Record completeness (per-
 394 centage of y during the record length for which annual data are available). Yellower (bluer)
 395 colors indicate more (less) complete records. **c**, Number of tide gauges returning annual
 396 sea-level data during the course of the study period.



397

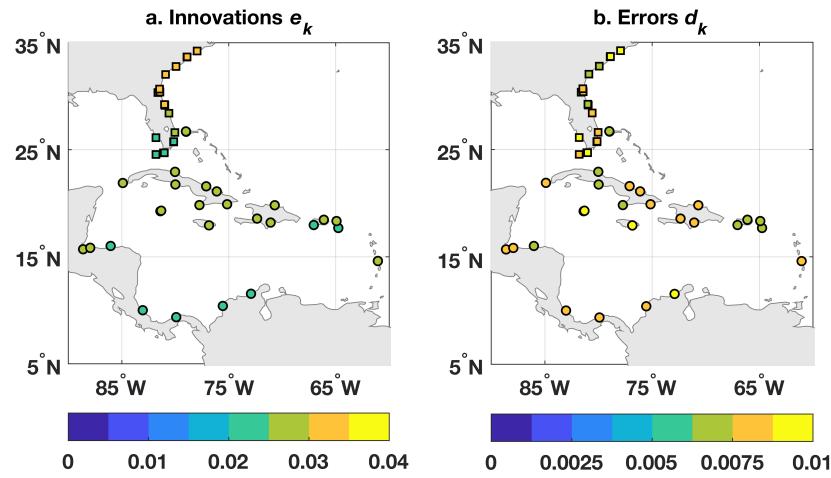
398 **Figure S8 Examples of residual time series.** Posterior median (solid lines) and pointwise
 399 95% credible intervals (light shading) of the sea-level **a**, process innovations e_k and **b**,
 400 data errors d_k at the San Juan (Puerto Rico) tide gauge. Posterior median (solid lines)
 401 and pointwise 95% credible intervals (light shading) of the transport **c**, noise sequence w_k
 402 and **d**, data errors u_k .



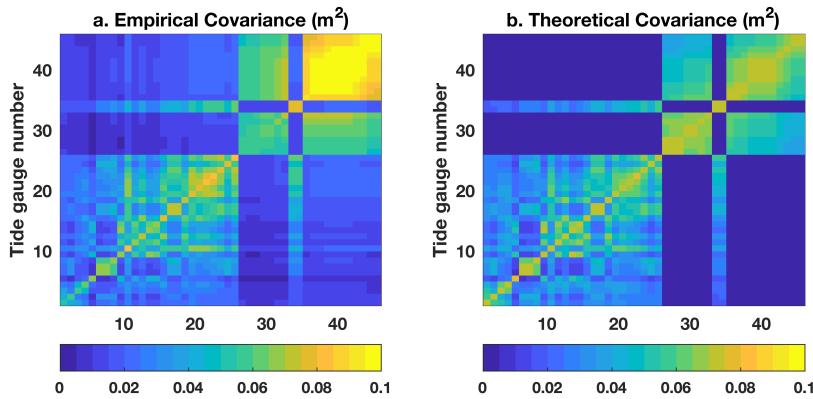
403

404 **Figure S9 Autocorrelation of the residuals.** Posterior medians (solid black) and pointwise
 405 95% credible intervals (gray shading) of the sample autocorrelation coefficient computed
 406 empirically from posterior solutions for the **a**, sea-level process innovations e_k , **b**, sea-
 407 level data errors d_k , **c**, transport noise sequence w_k , and **d**, transport data errors u_k . Solid
 408 and dashed blue lines are the mean \pm twice the standard error on the autocorrelation
 409 coefficients expected theoretically from white noise with the same temporal degrees of
 410 freedom.

411

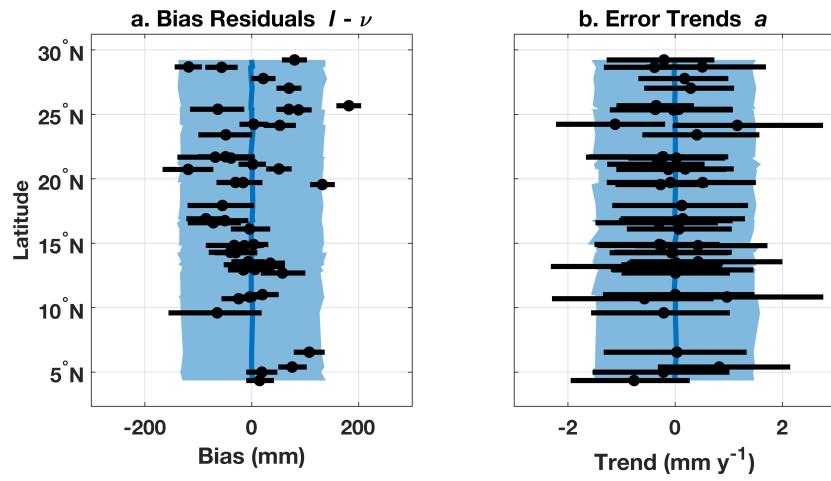


412 **Figure S10** Amplitude of sea-level residual time series. Median values of the standard
413 deviation (m) computed from posterior solutions for the sea-level **a**, process innovations
414 e_k and **b**, data errors d_k at all tide-gauge locations.



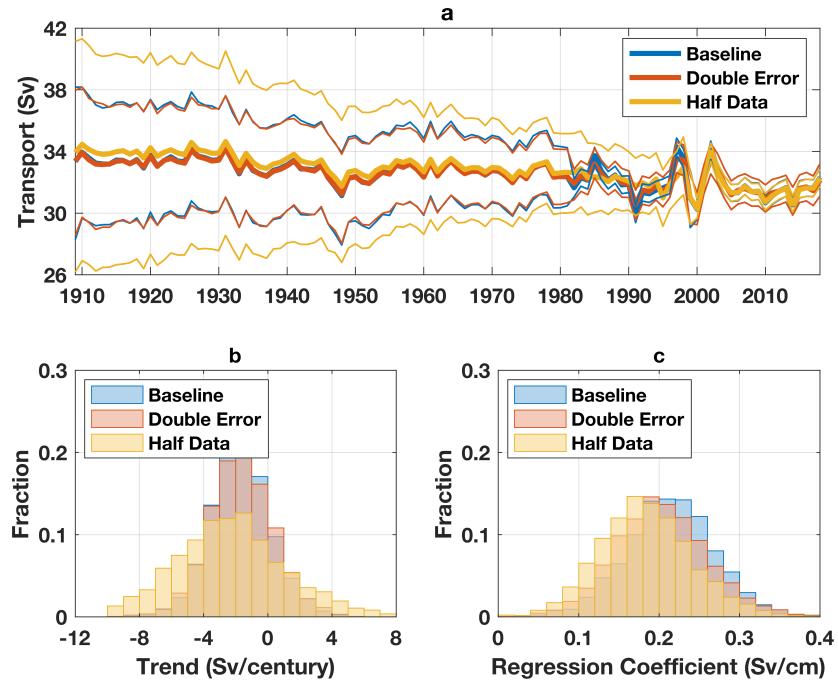
415

416 **Figure S11 Spatial covariance of sea-level process innovations.** Covariance (m^2) between
 417 all pairs of sea-level process innovations e_k computed **a**, empirically based on posterior
 418 solutions for e_k and **a**, theoretically using posterior solutions for σ^2 (Supplementary Ta-
 419 ble 6) and the assumed covariance structure Eq. (2). The “tide-gauge number” along x -
 420 and y -axes refer to the values given in the leftmost column in Supplementary Table 3.



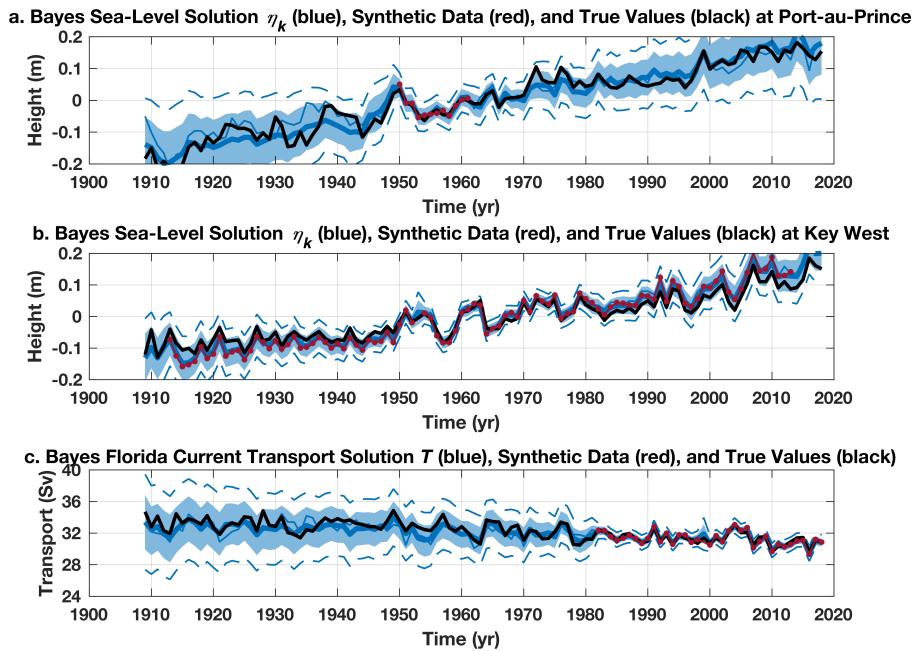
421

422 **Figure S12 Spatial structure of tide-gauge residual vectors.** Posterior medians (black
 423 dots) and pointwise 95% credible intervals (black lines) for the tide-gauge **a**, data-bias
 424 anomalies $\ell - \nu_1$ (m) and **b**, error trends a (mm y^{-1}). Also shown are the means (solid
 425 blue) and 95% credible intervals on these fields estimated from their assumed functional
 426 forms and posterior solutions for the respective variance parameters τ^2 and γ^2 (Supple-
 427 mentary Table 6).



428

429 **Figure S13** Sensitivity of Bayesian model solution to input transport data. Summary of
 430 results from sensitivity experiments using different forms of the Florida Cable transport
 431 data. **a**, Time series of transport (thick lines are posterior medians; thin lines bound
 432 the posterior 95% pointwise credible intervals). **b**, Histograms of the 110-y trend (1909–
 433 2018) in Florida Current transport. **c**, Regression coefficient between sea-level difference
 434 across Florida Straits and Florida Current transport. Blue values are from the “baseline”
 435 model experiment discussed in the main text. Orange values are based on an “double
 436 error” experiment wherein the standard errors on the transport data during 1982–2018
 437 are doubled. Yellow values are based on a “half data” experiment where the algorithm
 438 is only given the cable data during the period 2000–2018 and the 1982–1998 values are
 439 withheld. (There is no transport data value for 1999 due to a 20-month outage in the cable
 440 observations.)



441

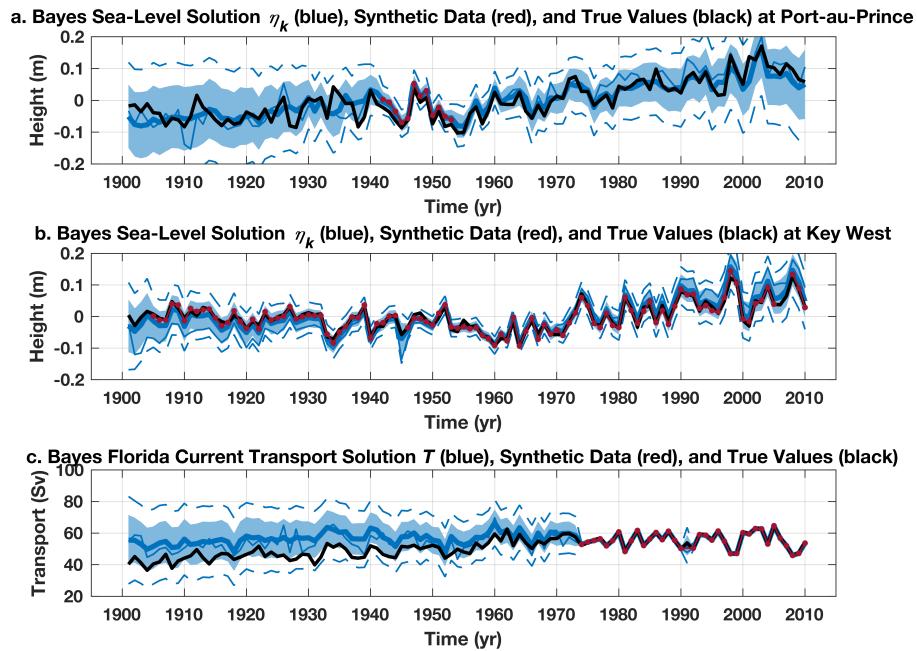
442 **Figure S14 Examples of results from first synthetic data experiment.** Synthetic observa-

443 tions (red), true values (black), and posterior medians (thick blue), pointwise (blue shad-

444 ing) and pathwise (dashed blue) 95% credible intervals, and an arbitrary ensemble mem-

445 ber (thin blue) of a, sea level at the Port-au-Prince (Haiti) tide gauge, a, sea level at the

446 Key West (USA) tide gauge, and c, Florida Current transport.



447

448 **Figure S15 Examples of results from second synthetic data experiment.** Synthetic ob-
 449 servations (red), true values (black), and posterior medians (thick blue), pointwise (blue
 450 shading) and pathwise (dashed blue) 95% credible intervals, and an arbitrary ensemble
 451 member (thin blue) of a, sea level at the Port-au-Prince (Haiti) tide gauge, a, sea level at
 452 the Key West (USA) tide gauge, and c, Florida Current transport.