Predictive Inverse Model for Advective Heat Transfer in a Planar Fracture with Heterogeneous Permeability

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Abstract

Identifying fluid flow maldistribution in planar geometries is a well-established problem in subsurface science/engineering. Of particular importance to the thermal performance of Engineered (or "Enhanced") Geothermal Systems (EGS) is identifying the existence of non-uniform (i.e., heterogeneous) permeability and subsequently predicting advective heat transfer. Here, machine learning via a Genetic Algorithm (GA) identifies the spatial distribution of an unknown permeability field in a two-dimensional Hele-Shaw geometry (i.e., parallel-plates). The inverse problem is solved by minimizing the L2-norm between simulated Residence Time Distribution (RTD) and measurements of an inert tracer breakthrough curve (BTC) (C-Dot nanoparticle). Principal Component Analysis (PCA) of spatially-correlated permeability fields enabled reduction of the parameter space by more than a factor of ten and restricted the inverse search to reservoir-scale permeability variations. Thermal experiments and tracer tests conducted at the mesoscale Altona Field Laboratory (AFL) demonstrate that the method accurately predicts the effects of extreme flow channeling on heat transfer in a single bedding-plane rock fracture. However, this is only true when the permeability distributions provide adequate matches to both tracer RTD and frictional pressure loss. Without good agreement to frictional pressure loss, it is still possible to match a simulated RTD to measurements, but subsequent predictions of heat transfer are grossly inaccurate. The results of this study suggest that it is possible to anticipate the thermal effects of flow maldistribution, but only if both simulated RTDs and frictional pressure loss between fluid inlets and outlets are in good agreement with measurements.

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2 **Heterogeneous Permeability**

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9 **Key Points:**

10 • Rapid heat transfer accurately predicted at a meso-scale field laboratory.

11 • Genetic algorithm identified channeled flow via tracer/pressure calibration.

12 • Principal component analysis reduced the number of solved variables by more than a
13 factor of ten.

14

Abstract

15

16 Identifying fluid flow maldistribution in planar geometries is a well-established problem in
17 subsurface science/engineering. Of particular importance to the thermal performance of
18 geothermal reservoirs is identifying the existence of non-uniform (i.e., heterogeneous)
19 permeability and subsequently predicting advective heat transfer. Here, machine learning via a
20 Genetic Algorithm (GA) identifies the spatial distribution of an unknown permeability field in a
21 two-dimensional Hele-Shaw geometry. The inverse problem is solved by minimizing the L^2-
22 norm between measured and simulated fluid Residence Time Distribution (RTD). Principal
23 Component Analysis (PCA) of spatially-correlated permeability fields enabled reduction of the
24 parameter space by more than a factor of ten and restricted the inverse search to large-scale
25 permeability variations. Thermal experiments and tracer tests conducted at a meso-scale field
26 laboratory in Altona, New York demonstrate that the method accurately predicts the effects of
27 extreme flow channeling on heat transfer in a bedding-plane rock fracture. However, this is only
28 true when the permeability distributions provide adequate matches to both tracer RTD and
29 frictional pressure loss. Without good agreement to frictional pressure loss, it is still possible to
30 match a simulated RTD to measurement, but subsequent predictions of heat transfer are grossly
31 inaccurate. The results of this study suggest that it is possible to anticipate the thermal effects of
32 flow maldistribution, but only if both simulated RTDs and frictional pressure loss between fluid
33 inlets and outlets are in good agreement with measurements.

1 Introduction

34

35   Accurate forecasting of heat and mass transfer in two-dimensional planes (e.g., Hele-
36 Shaw cells, groundwater aquifers, microchannels, magma flow etc.) requires adequate
37 specification of reservoir permeability. This is particularly true when non-uniform (i.e.,
38 heterogeneous) permeability disrupts the flow pattern between an inlet and outlet. If permeability
39 distributions promote preferential flow paths, such conditions can produce an inlet-outlet "short-
40 circuit" which is directly analogous to electrostatics described by Ohm's law (e.g., Darnton et al.,
41 2001; Levitov and Falkovich, 2016; Oh et al., 2012). This is a well-established problem in
42 subsurface hydrology (e.g., Alley et al., 2002; Black et al., 2017; Cathles, 1990; Doughty, 1999;
43 Guo et al., 2016; Kurowski et al., 1994; Larsson et al., 2013; Moreno and Neretnieks, 1993;
44 Mortada and Nabor, 1961; Subramanian et al., 2013; Tsang and Tsang, 1988; Zouache et al.,
45 2019; etc.). More recently, fluid flow maldistribution has been identified as a key issue that
46 "plagues" (Toma, 2017) energy functional materials and microfluidics (e.g., Al-Housseiny et al.,
47 2012; Babu and Ajayan, 2019; Lagadec et al., 2019; Stone et al., 2004; etc.) and can play a
48 significant role in blood flow maldistribution (e.g., Pries et al., 2010; Qiu et al., 2019; Zouache et
49 al., 2016; etc.)

50   In subsurface science and engineering, quantitatively characterizing the spatial
51 distribution of permeability is exceptionally challenging. This is particularly true in fractured
52 rock hydrogeology where fracture permeability, k, is proportional to the square of the spacing or
53 "aperture," b, between the upper and lower surfaces ($k = b^2/12$). Therefore, laterally extensive
54 fracture planes can transmit fluids at reservoir-scale (e.g., 100s of meters) through a single
55 fracture while permeabilities can range by orders of magnitude even at the local scale (e.g.,
56 centimeters). As a result, it is exceedingly difficult to use local-scale measurements of fracture
57 aperture (e.g., near observation wells) to anticipate reservoir-scale flow patterns.

58   This difficulty introduces considerable uncertainty into predicting heat and mass transfer
59 in rock fractures and has significant implications for the movement of both subsurface resources
60 (e.g., hot water/steam, petroleum, dissolved minerals, etc.) and contaminants (e.g., radionuclides,
61 fertilizers, chlorinated solvents, etc.). In an Engineered (or "Enhanced") Geothermal System
62 (EGS), for example, commercial success requires development of a deep subsurface heat
63 exchanger capable of heating cold injectate before circulating fluids are produced for direct use
64 or electricity generation. If inlet-outlet short-circuiting reduces the effective heat transfer area,
65 however, then production wells may cool unexpectedly prior to the plant reaching financial
66 maturity (Hawkins et al., 2018). Rapid thermal drawdown at one or more production wells is
67 often referred to as "premature thermal breakthrough" and has challenged the growth of low-
68 grade geothermal resources for decades (e.g., Brown et al., 1999; Hawkins et al., 2018; Horne,
69 1985; Murphy et al., 1981; Pruess, 1990; Robinson and Tester, 1984, 1986; Robinson et al.,
70 1988; Shook and Suzuki, 2017; Suzuki et al., 2019; Tester et al., 1982, 1989, 2006; etc.).

71   Here, an inverse model is presented for identifying the spatial distribution of fracture
72 aperture in a two-dimensional plane (i.e., a "Hele-Shaw" geometry) based on calibration to a
73 specified inlet-outlet RTD and frictional pressure loss. The inversion procedure consists of two
74 stages. In the first stage, the inverse problem is solved by minimizing an objective function that
75 compares simulated Residence Time Distribution (RTD) to the measured breakthrough curve
76 (BTC) of an inert tracer. In the second stage, the minimum and maximum fracture aperture is
77 constrained by comparing measured and simulated frictional pressure loss. The resulting fracture
78 aperture field is then used to forecast advective heat transfer using the numerical model
79 summarized in subsection 3.3 and originally presented in Fox et al. (2015).

80   A Genetic Algorithm (GA), which mimics the optimizing behavior of natural selection, is
81 employed here to solve this highly non-linear inverse problem. Principal Component Analysis
82 (PCA) of self-affine fractures made this problem computationally tractable by reducing the
83 number of optimized variables from thousands to as few as five. As a result of PCA, this global
84 optimizer if forced to investigate the effects of large-scale aperture fluctuations on transport
85 processes that dominate at the reservoir-scale.

86   The capabilities and predictive value of this novel inversion routine are demonstrated in
87 the results section based on meso-scale experiments performed at a well-characterized field site
88 known as the Altona Field Laboratory (AFL). These field tests targeted a sub-horizontal bedding
89 plane fracture located roughly 7.6 m below ground surface and this target fracture is known to
90 exhibit channeled-flow behavior. An inert tracer test performed under forced-convection
91 between an injector-producer well pair spaced 14.1 m apart provided the measured BTC data. A
92 thermal experiment originally presented in Hawkins et al. (2017a, 2018) provided the
93 temperature measurements needed to evaluate predictions that resulted from the optimized
94 aperture fields and a brief summary of the thermal experiment is provided here in subsection 2.2.

95 **2 Altona Field Laboratory (AFL)**

96   The predictive inverse model described in this study is tested against meso-scale field
97 tests conducted at the Altona Field Laboratory (AFL). This site is located in northeastern New
98 York State near the USA/Canadian border. Since the first five wells were drilled in 2004, the 10
99 x 10 m five-spot well configuration has been the target of numerous investigations into flow
100 channeling in sparsely-spaced fracture planes (Figure 1). These investigations included
101 geophysical surveys and hydraulic/tracer testing and primarily targeted a sub-horizontal bedding-

102 plane fracture located roughly 7.6 m below ground surface (Becker and Tsoflias, 2010; Guiltinan
103 and Becker, 2015; Hawkins et al., 2017b; Talley et al., 2005; Tsoflias and Becker, 2008; Tsoflias
104 et al., 2015; etc.). As discussed in Hawkins et al. (2017a, 2018) fluids are highly concentrated in
105 an east-west direction between well 204 and 304. Based on the combined results of numerical
106 models, adsorbing tracer tests, Ground Penetrating Radar (GPR) surveys, and Fiber-Optic
107 Distributed Temperature Sensing (FO-DTS), a narrow channel roughly 1-2 m wide spans the
108 14.1 m distance between well 204 to the west and well 304 to the east. The first subsection that
109 follows below describes the results of a tracer test performed in 2016 using inert carbon-cored
110 nanoparticles referred to as "C-Dots" (e.g., Hawkins et al., 2017a). Following the description of
111 this tracer test is a brief summary of the heat exchange experiment conducted in 2015 which is
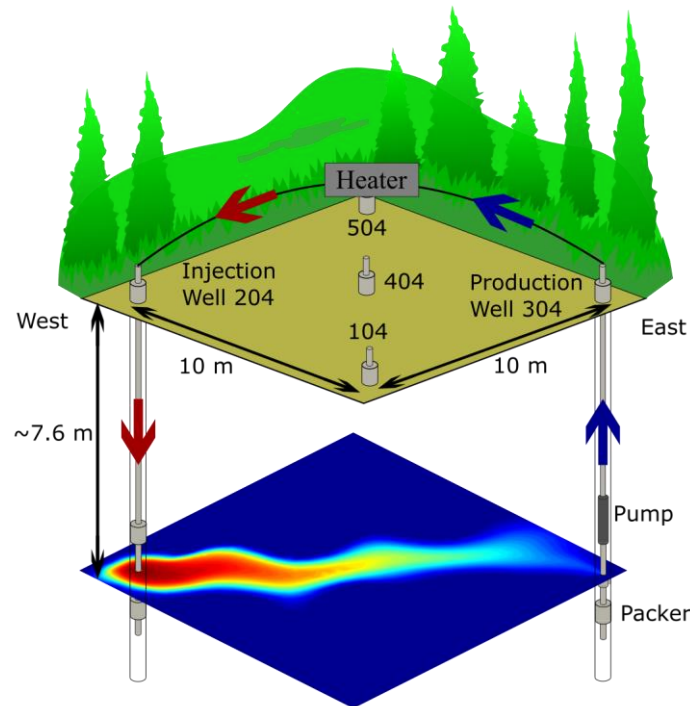112 described in detail in Hawkins et al. (2017a, 2018).

113

114 Figure 1. Three-dimensional schematic of the Altona Field Laboratory (AFL) showing the spatial
115 configuration of injection well 204 and production well 304 as well as three monitoring wells
116 (104, 404, and 504). The red and blue arrows depict the circulation pattern introduced during the
117 heat exchange experiment in which hot water continuously circulated through the target fracture
118 for 6 d. The contour map shown at ~7.6 m below ground surface is the temperature distribution
119 predicted after 24 h as a result of the aperture distribution shown in the upper-left corner of
120 Figure 6.

121   2.1 Inert Tracer Test (2016)

122   A carbon-cored nanoparticle tracer, or "C-Dots," served as an inert tracer for field testing
123 at the AFL in the summer of 2016. This fluorescent tracer consists of a carbon core decorated
124 with a highly fluorescent polymer (Krysmann et al., 2012). It is ideal for groundwater tracer
125 testing, because it is highly water soluble, inert, environmentally benign, and detectable by

126 fluorescence at concentrations as low as ~1 ppb in deionized water (Hawkins et al., 2017a; Ray
127 et al., 2009; Subramanian et al., 2013; Zhao, 2015). C-Dots were used in past field tests at the
128 Altona site in the fall of 2015 to estimate the Residence Time Distribution (RTD) between
129 injection well 204 and production well 304 (Hawkins et al., 2017a). A comparison to an organic
130 tracer (phenol) and an anion tracer (iodide) confirmed that the C-Dot nanoparticle tracer behaved
131 conservatively at the Altona site.

132 A constant volumetric flow rate of 5.82 L/min at the injection/production wells
133 established a steady-state forced-convection hydraulic field (Table 1). A C-Dot mass of 402.2 g
134 was introduced into the formation following the approach described in Hawkins et al. (2017a,
135 2018) in which a PVC bypass system allowed for introduction of the tracer with minimal
136 disturbance to the subsurface hydraulic gradient field. Polypropylene centrifuge tubes were used
137 to collect 15 mL groundwater samples for laboratory analysis. An ice-filled cooler preserved the
138 fluid samples until the time of analysis by excitation/emission spectra measured by a Cary
139 Eclipse fluorescence spectrophotometer manufactured by Agilent Technologies.

140 Produced fluids during the tracer test were immediately circulated back into the injection
141 well. A straddle-packer assembly (Lansas Products) hydraulically isolated flow to the target
142 fracture and constrained wellbore storage volume to roughly 3 to 5 L in the injection well and
143 three observation wells. At the production well, a single packer provided hydraulic isolation
144 below the fracture. In the 2015 tracer tests (Hawkins et al., 2017a), the injection well straddle
145 packer assembly reduced wellbore storage to roughly 8 L. In the present study, this volume was
146 reduced further to 5 L.

147 Table 1. Description of the 2016 inert tracer test and modeled parameters

Symbol	Description	Value	Unit
Q	volumetric flow rate	9.7e-5	m^3/s
μ_w	dynamic viscosity of water	1.0e-3	Pa-s
ρ_w	fluid density	1000	kg/m^3
M_{inj}	[a]tracer mass injected	0.402e-3	kg

148 [a]tracer injectate (2.895 L) consisted of 4.190 mL of concentrated C-Dot solution (96,000 ppm) added to 2.891 L of groundwater
149 collected from well 304

150 2.2 Heat Exchange Experiment (2015)

151 In the fall of 2015, field testing at the AFL provided a low-temperature analog to simulate
152 conditions in an operating geothermal reservoir at meso-scale. For six days, heated injectate (74
153 °C) circulated through the initially cold (11.7 °C) reservoir (Table 2). A Grundfos submersible
154 pump installed in a shroud assembly maintained a continuous fluid flow rate of 5.7 L/min out of
155 production well 304. All produced fluids were piped through a propane-fired tankless water
156 heater (AO Smith, Model #: T-H3-OS) and reintroduced into well 204. The design and results of
157 the thermal test are described in detail in Hawkins et al. (2017a, 2017b, 2018).

158 During reservoir heating, a Fiber-Optic Distributed Temperature Sensor (FO-DTS)
159 (Sensornet Ltd, Model: Oryx) recorded the spatial distribution of fracture/matrix heat exchange
160 in ten monitoring boreholes located throughout the 10 x 10 m well field. The FO-DTS measured
161 temperature along a 1400 m cable wrapped around ten externally threaded schedule-80 Polyvinyl
162 Chloride (PVC) pipes. Wrapping the fiber-optic cables resulted in roughly 70 to 140 temperature
163 measurements along the 1.5 to 3 m long pipes.

164   FO-DTS temperature measurements were recorded in the dry, bulk rock matrix above the
165 target fracture with roughly 0.5 m of undisturbed rock remaining between the deepest
166 measurement and the target fracture. Fracture fluid temperatures could then be determined based
167 on an error function extrapolation of the measured temperature data to the depth of the target
168 fracture (Hawkins et al., 2017a). In addition to the FO-DTS temperature measurements,
169 temperature sensors measured fluid temperature in three monitoring wells (Solinst Leveloggers,
170 Model #: 3001) and in the production well (RBR, Model: Solo T).

171 Table 2. Description of the 2015 thermal test and modeled parameters

Symbol	Description	Value	Unit
Q	[a]volumetric flow rate	9.5e-5	m^3/s
T_r	[b]ambient rock temperature	11.7	°C
T_w	[b]injected water temperature	74.0	°C
$C_{p,w}$	specific heat of water	4200	J/kg-K
$C_{p,r}$	[c]specific heat of sandstone	1000	J/kg-K
μ_w	dynamic viscosity of water	1.0e-3	Pa-s
λ_r	[b]thermal conductivity of Potsdam Sandstone	7.6	W/m-K
ρ_r	[b]density of Potsdam Sandstone	2500	kg/m^3
ρ_w	fluid density	1000	kg/m^3
t_{th}	total elapsed time of the thermal test	5.184e5	s

172 [a]volumetric flow rate into injection well 204 and out of production well 304 were equal and constant
173 [b]source: Hawkins et al. (2017a)
174 [c]source: Robertson (1988)

3 Inverse Problem Description and Numerical Methods

176   An inverse problem is a computational formulation where one attempts to calculate the
177 parameters that describe a system, making use of a set of data or observations. It is the reverse of
178 a forward model which calculates the response of a system for a given set of parameter values.
179 Mathematically, the inverse problem is formulated as $d_i = F(x_j)$, where d_i is the vector of data
180 or observations and F is a model often comprised of algebraic and/or differential equation or sets
181 of equations which depend on x_j. The inverse problem is solved by determining the x_j given d_i
182 based on minimizing an "objective function" which evaluates the fitness of a calculated value to
183 measurement.

184   For the inverse problem presented here, x_j represents local values of the fracture aperture
185 field, F is a numerical model describing the behavior of the injected C-Dot tracer (subsection
186 2.1), and d_i is the measured BTC of tracer introduced at a single injection well and eluted from a
187 single production well under steady-state hydraulic conditions. The description of the
188 permeability field is first discussed (subsection 3.1), followed by an explanation of how tracer
189 RTD is calculated (subsection 3.2). Finally, the numerical model of heat transfer described in
190 Fox et al. (2015) is summarized in subsection 3.3.

191   3.1 Spatially-correlated Fracture Aperture Distribution

192   Laminar fluid flow through naturally permeable materials such as rock fractures
193 commonly exhibit variable flow paths as a result of spatially fluctuating permeability. The
194 permeability field adopted here is a self-affine description where relatively large-scale variations
195 in permeability (small frequencies, long wavelengths) are more important for describing heat and

196   mass transport compared to small-scale variations (Feder, 1988; Falconer, 2003). The primary
197   purpose of the self-affine model is to systematically generate target permeability distributions
198   with different flow path geometries which, in effect, "trains" the optimization routine. Core
199   samples containing rock fractures from wide-ranging geographies have revealed fracture
200   apertures to have a self-affine spatial correlation which suggests that their subsequent
201   permeability distributions are also self-affine (Brown et al., 1986; Glover et al., 1998;
202   Plouraboué et al., 1995; Ponson et al., 2007; Schmittbuhl et al., 1993, 1995, 2008). Below is a
203   detailed description of the varying permeability field with a self-affine spatial correlation.

204   The permeability field described here is based on the equivalency to fracture aperture
205   described as $k = b^2/12$, where k and b are fracture's permeability and aperture, respectively.
206   The spatially varying aperture field can therefore be described as $b(x, y) = \langle b \rangle + h(x, y)$, where
207    $\langle b \rangle$ is the mean aperture value and h is the local aperture fluctuation. A Hurst roughness
208   coefficient, ζ, and a spatial correlation function are used to describe the spatially-correlated
209   aperture field (e.g., Di Federico and Neuman, 1997, 1998) and the variogram function, γ,
210   required for describing the spatial correlation is

211    $$\gamma(d) \equiv \langle [h(\mathbf{p}) - h(\mathbf{p} + d)]^2 \rangle = 2[\langle h^2 \rangle - \langle h(\mathbf{p})h(\mathbf{p} + d) \rangle],$$ (1)

212   where $\mathbf{p}$ is a position vector on the surface of the 2-D fracture and d is the displacement distance
213   from $\mathbf{p}$. To have a self-affine geometry, Equation 1 must have a scaling behavior that follows
214    $\gamma(d) = a|d|^{2\zeta}$, where a is an arbitrary constant.

215   Based on core samples, Hurst coefficients range from 0.5 for sandstone (e.g., Boffa et al.,
216   1999; Ponson et al., 2007) to 0.8 for granite (e.g., Schmittbuhl et al., 1993, 1995; Plouraboué et
217   al., 1995) and a surface with a lower Hurst coefficient is rougher. An algorithm modified from
218   Méheust and Schmittbuhl (2001) generated numerous hypothetical self-affine fracture aperture
219   distributions. Random seeding using Fast Fourier Transforms (FFTs) produces hypothetical
220   fracture aperture fields that are both statistically isotropic and self-affine (Figure 3).

221   3.2 Numerical Model for Calculation of Residence Time Distribution (RTD)

222   Laminar flow of an incompressible fluid between two parallel plates is described by the
223   Hele-Shaw equation, which is often referred to as the "cubic law" in fractured rock
224   hydrogeology (Witherspoon et al.,1980). As described in Fox et al. (2015), the velocity field can
225   be derived by the finite element method treatment of

226    $$-\nabla \cdot \left(\frac{b^3}{12\mu} \nabla P \right) = 0,$$ (2)

227   where P is the hydraulic fluid pressure and μ is the fluid's dynamic viscosity, assuming the
228   lubrication approximation is valid for the specified length scales. The velocity averaged across
229   the fracture aperture is then given by $\langle \mathbf{v} \rangle = -b^2/(12\mu) \nabla P$. With velocity specified, the
230   transport of tracer particles in the permeability field follows the governing equation

231    $$b \frac{\partial c}{\partial t} + b \langle \mathbf{v} \rangle \cdot \nabla c = 0,$$ (3)

232   where c is tracer concentration and t is time.

233   The modeling approach described by Equation 3 neglects longitudinal/transverse
234   dispersion resulting from the combined effects of velocity and concentration gradients across the
235   fracture aperture (i.e., Taylor-Aris dispersion). Instead, tracer RTD variance is produced solely

236 by non-uniform velocity in the plane of the fracture. These velocity variations result from the
237 combined effects of dipole-flow and heterogeneous permeability. As discussed in Hawkins et al.
238 (2017a, 2018), this approach is more appropriate for tracer transport between parallel plates
239 where Taylor-Aris dispersion, as described by Horne and Rodriguez (1983), should be negligible
240 relative to velocity variations resulting from dipole flow and heterogeneous permeability
241 (Detwiler et al., 2000; Neretnieks, 1980; Novakowski et al., 1985; Robinson et al., 1988; Tang et
242 al., 1981).

243   Flow through the fracture plane is induced by steady-state circulation between an
244 injection and production well and there is a no flux condition on the boundaries of the fracture.
245 The wells are modeled as an internal boundary with 64 FEM computational nodes rather than a
246 single node to avoid singularities in velocity. The total flux across the internal boundaries
247 representing the wells is equal to the volumetric flow rate, Q.

248   To model a pulse injection tracer test, the initial condition is $c(x, y, 0) = 0$, and the
249 boundary condition at the location of the injection well $\mathbf{p}_{inj}$ is $c(\mathbf{p}_{inj}, t) = m_p \delta(t - t_0)/Q$,
250 where m_p is the mass of injected tracer, δ is the Dirac delta function, and t_0 is the time of
251 instantaneous tracer injection. Additionally, there is a no mass flux condition along the lateral
252 boundary of the fracture. The concentration of tracer at the production well with time produces a
253 tracer breakthrough curve that is a result of the aperture field and resulting flow in the fracture.
254 By simulating an instantaneous injection of an inert tracer, the concentration curve represents the
255 RTD of the permeable reservoir.

256   3.3 Numerical Model for Advective Heat Transfer

257   After the permeability field from the inverse search has been obtained, it is used to
258 calculate advective heat transfer and subsequently forecast production well fluid temperature as a
259 function of time. The comparison of predicted temperatures resulting from the optimized
260 aperture field with measured temperatures recorded at the AFL determines how effective the
261 inversion scheme was in predicting a reservoir's thermal performance. The temperature field is
262 determined using the numerical model presented in Fox et al. (2015). It calculates the heat
263 transport within the fracture aperture field using the FEM and couples it to a boundary integral
264 description of the heat conduction in the impermeable material surrounding the fracture plane.
265 Note that this simulation neglects buoyancy effects and assumes the fluid and rock properties are
266 negligible.

267 **4 Inversion Algorithm**

268   To solve for the permeability field, an inversion algorithm was developed using reduced
269 order modeling and a global optimizer. The optimization algorithm is used to minimize the error
270 against a simulated tracer RTD curve, expressed as an objective function. The L^2-norm error is
271 used as the objective function and is defined as $L_d^2 = \left\| \hat{d}_i - d_i \right\|^2 / \|d_i\|$, where d_i is the vector of
272 measured concentration data as a function of time while the $\hat{d}_i$ terms are the simulated values
273 resulting from a candidate aperture field of the inverse search.

274   A GA is used for the parameter optimization and a reduced order model of the aperture
275 field is obtained using PCA, an algorithm that determines the dominant modes of a system. The
276 description and application of PCA for the problem is first presented, followed by the
277 explanation of how the GA is used to solve the inverse problem.

278   4.1 Principal Component Analysis

279   Principal Component Analysis (PCA) (sometimes called "proper orthogonal
280   decomposition") is a statistical tool common to engineering sciences. PCA is particularly useful
281   in the analysis of datasets with a large number of independent variables (e.g., Cui et al., 2015;
282   Johnston et al. 2009, Li and Jarfarpour, 2010; van Doren et al., 2006; etc.). In the current study,
283   PCA reduces the number of unknown parameters requiring estimation via the inverse model by
284   targeting larger, spatially-correlated regions of the aperture field. Rather than independently
285   estimating each fracture aperture value in the FEM mesh, PCA restricts the inverse search to
286   reservoir-scale heterogeneities by varying the intensity of PCA modes (eigenvectors) with the
287   largest modal coefficients (eigenvalues).

288   The first step in PCA is to calculate the covariance matrix between independent variables
289   in a dataset. Then, the eigenvectors and eigenvalues of the covariance matrix are calculated.
290   Each eigenvector represents a "mode" of the dataset (Figure 2). The eigenvalues represent the
291   relative strength of each mode. For spatially-correlated fractures, the mode with the largest
292   eigenvalue corresponds to spatial variations over relatively large areas of the fracture.
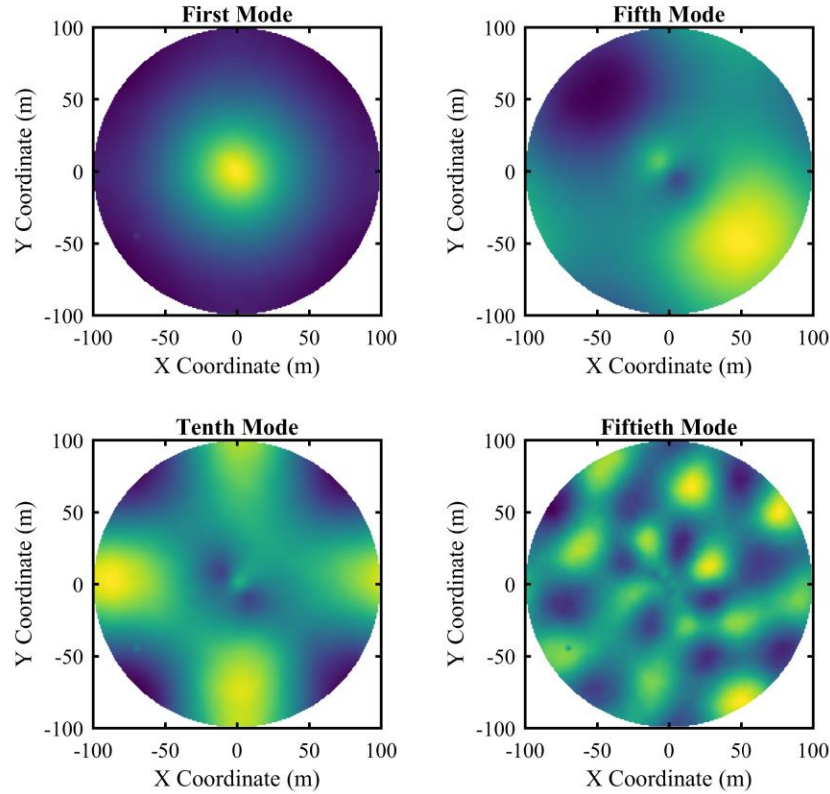
293

Figure 2. Modes resulting from PCA of a hypothetical circular fracture with a radius of 100 m.
The color scale is arbitrary as the strength of each mode fluctuates during the Genetic Algorithm
(GA) optimization. The first mode (largest-scale variations) is shown in the upper-left corner and
is compared to modes that describe spatial variations at smaller scale including the fifth mode
(upper right), the tenth mode (bottom left), and the fiftieth mode (bottom right).

For PCA of a generic data set, X_{ij}, describing realization j and variable i where n variables
each have a mean of zero across m realizations, the covariance matrix, $C_{ii'}$, between two
variables is given by

$$C_{ii'} = \sum_{j=1}^{m} X_{ij} X_{i'j}. \tag{4}$$

For a discrete heterogeneous fracture represented by nodes on a computational mesh, n is
the number of nodes and m is the number of aperture fields generated. Once the eigenvectors (or
"modes") and eigenvalues (or "coefficients") of the covariance matrix have been determined, an
aperture field in the matrix of observations can be reproduced by knowing the modes of the
entire dataset and the coefficients for the specific aperture field, such that

$$w_i = \sum_{j=1}^{p} \Phi_{ij} \beta_j + \langle w \rangle, \tag{5}$$

where w_i is the fracture aperture at node i, Φ_{ij} is the matrix of PCA modes, β_j are the coefficients
of the PCA modes, and $\langle w \rangle$ is the mean fracture aperture. The reproduced aperture field is
identical to the original if p is equal to n, the total number of nodes. If p is less than n, then
Equation 5 produces a simplified aperture field that neglects small-scale spatial variations as

313   illustrated in Figure 3. In addition, entirely new aperture distributions result from altering the
314   magnitude of the mode coefficients. Any new distribution will have a unique aperture field, but
315   the eigenvectors of its covariance matrix (i.e. its PCA modes) are identical to the original data
316   set. A description of the FEM mesh and parameters of the PCA are provided in Table 3.

317   Table 3. Description of the finite element mesh and the principal component analysis

Symbol	Description	Value
n_{nodes}	number of finite element nodes	7196
n_{elem}	number of finite elements	7100
n_{inj}	number of injection well nodes	64
n_{prod}	number of production well nodes	64
n_{PCA}	number of samples for PCA	5000
ζ	[a]Hurst roughness coefficient	0.5

318   [a]average value for sandstone based on Boffa et al. (1999) and Ponson et al. (2007)
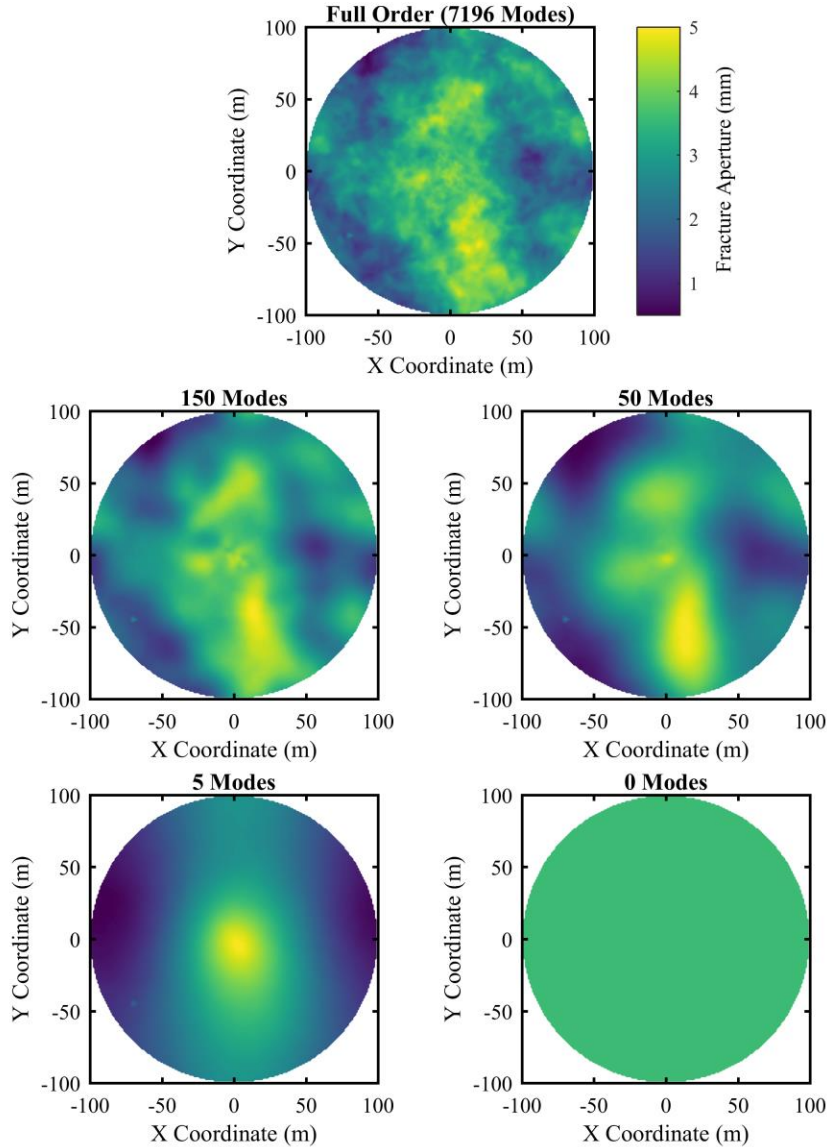
Figure 3. Hypothetical 2-D fracture aperture field (top) compared to reduced order aperture fields ranging from 150 modes (middle left) to 0 modes (i.e., homogeneous) (bottom right).

4.2 Genetic Algorithm

The challenge in solving an inverse problem for heterogeneous permeability is that the number of unknown variables is typically too large for conventional inversion routines (McLaughlin and Townley, 1996). For a circular fracture with non-uniform aperture, accurate numerical simulations of heat and mass transport could easily require thousands of specified aperture values in the two-dimensional field. Therefore, an attempt to solve an inverse problem by independently varying the values at each node in the computational mesh is inefficient considering the computational time required and the relatively small influence of local-scale aperture fluctuations on tracer and thermal transport at the reservoir scale. As discussed in the previous subsection, PCA simplified this problem by allowing the inverse model to vary the

332 aperture field in a spatially-correlated fashion. As a result, the only information that must be
333 specified, *a-priori*, is the Hurst roughness coefficient, the number of PCA modes required in
334 solving the inverse problem, and (if desired) the mean or the minimum and maximum aperture.

335   Even after PCA, there remains several challenging aspects in successfully employing an
336 inverse model for heterogeneous aperture fields, including: (1) the inverse problem is non-linear;
337 (2) there remains a relatively large number of unknown variables; and (3) there are likely many
338 local minimums of the objective function. For these reasons, a stochastic inverse search method
339 is preferable to a deterministic minimization which would require long computation times to
340 search the parameter space and suffers from inefficiencies while searching near local minima.
341 One commonly used metaheuristic optimization method is a Genetic Algorithm (GA). It is
342 designed to mimic natural selection and minimize an objective function, $f(x_j)$, where x_j is a set
343 of parameter values (e.g., Cao et al., 2012; Hornby et al., 2011; Tasan and Gen, 2012). The GA
344 is similar to natural selection in that the "fitness" of an initial population is quantified and new
345 "generations" are produced by "breeding" and/or "mutating" individuals in the population that
346 are specified using a suitable selection function. Once the desired conditions for termination are
347 met (e.g., number of generations, target objective function, computation time, etc.), the
348 optimization is complete (Figure 4).
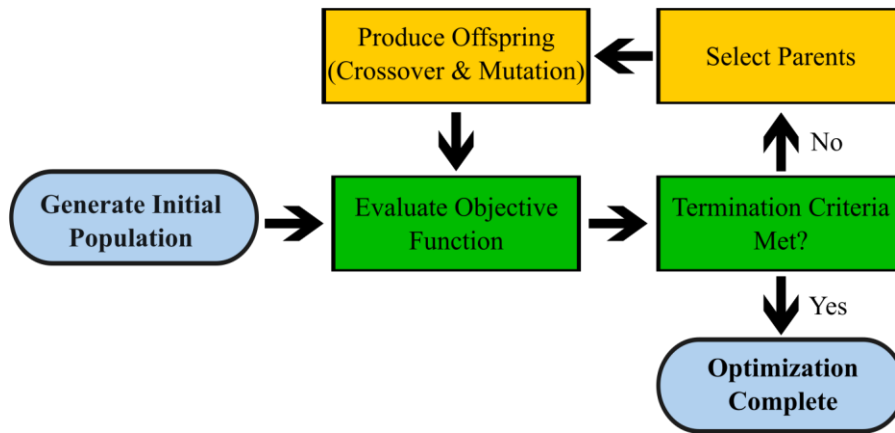

349

350 Figure 4. Flow diagram for the Genetic Algorithm (GA) demonstrating the optimization process.

351   For the first stage of the optimization, the parameter values for the GA are the PCA mode
352 coefficients, β, and the min/max fracture aperture. In stage 2, the min and max fracture apertures
353 are specified based on agreement between measured and simulated frictional pressure loss. For
354 both stages, the objective function is the L^2-norm between measured and simulated tracer BTCs.
355 Parameters of the initial population were determined using the range of β values given an
356 arbitrary standard deviation used during PCA. The parent selection function selects individuals
357 for "reproduction" and this study adopted the tournament style. A genetic crossover function
358 generated 80% of the "offspring" and a mutation function produced the remaining 20%. The
359 mutation function randomly changes the "genes" of individual parents, which in this study are
360 the PCA mode coefficients or the min/max aperture and a maximum of 60 generations is
361 specified as the termination criteria (Table 4).

362   Table 4. Input parameters of the two-stage GA optimization

Symbol	Description	Stage 1	Stage 2
n_{pop}	population size in each generation	144	120
n_{gens}	maximum number of generations	60	60
n_{elite}	elite fit carryover	2	2
n_{var}	number of solved variables	[a]1-152	150
w_{min}	minimum aperture	[b]n/a	0.02e-3 m
w_{max}	maximum aperture	[b]n/a	8e-3 m
f_{cross}	crossover fraction	0.8	0.8

363   [a]ranges from 1 (i.e., 0 modes/homogeneous) to 152 (i.e., 150 modes and minimum/maximum aperture)
364   [b]stage 1 heterogeneous scenarios consider both the minimum and maximum aperture as unknown variables solved by the GA

5 Results

366   The results of the optimization routine are broken down into two subsections. The first
367   subsection describes the results of the two-stage optimization. The second subsection compares
368   predicted temperatures with measurements at the AFL.

369   5.1 Optimized Aperture Field

370   As discussed in subsection 4.2, the main input parameter required for the first stage of
371   this optimization is the number of PCA modes selected for optimization. With the number of
372   modes selected and the initial range of mode coefficients and apertures specified, the GA
373   searches for an optimal aperture distribution by varying modal coefficients and the upper/lower
374   aperture boundaries. Four scenarios are considered in the first stage ranging from a homogeneous
375   aperture field (i.e., zero PCA modes) to a heterogeneous aperture field composed of 150 PCA
376   modes.

377   The first scenario of the first stage is presented as a base case in which the fracture
378   aperture field is homogeneous. As shown in the top row of Figure 5, a uniform aperture of 0.51
379   mm results in reasonably good agreement with tracer measurements (Table 5). Both the
380   simulated and measured modal residence times (i.e., elapsed time corresponding to peak tracer
381   concentration) occur at 20 min. The model-fit peak concentration (0.0361 min^{-1}) at the modal
382   residence time is in slight disagreement with measurement (0.04 min^{-1}) and the L^2-norm for the
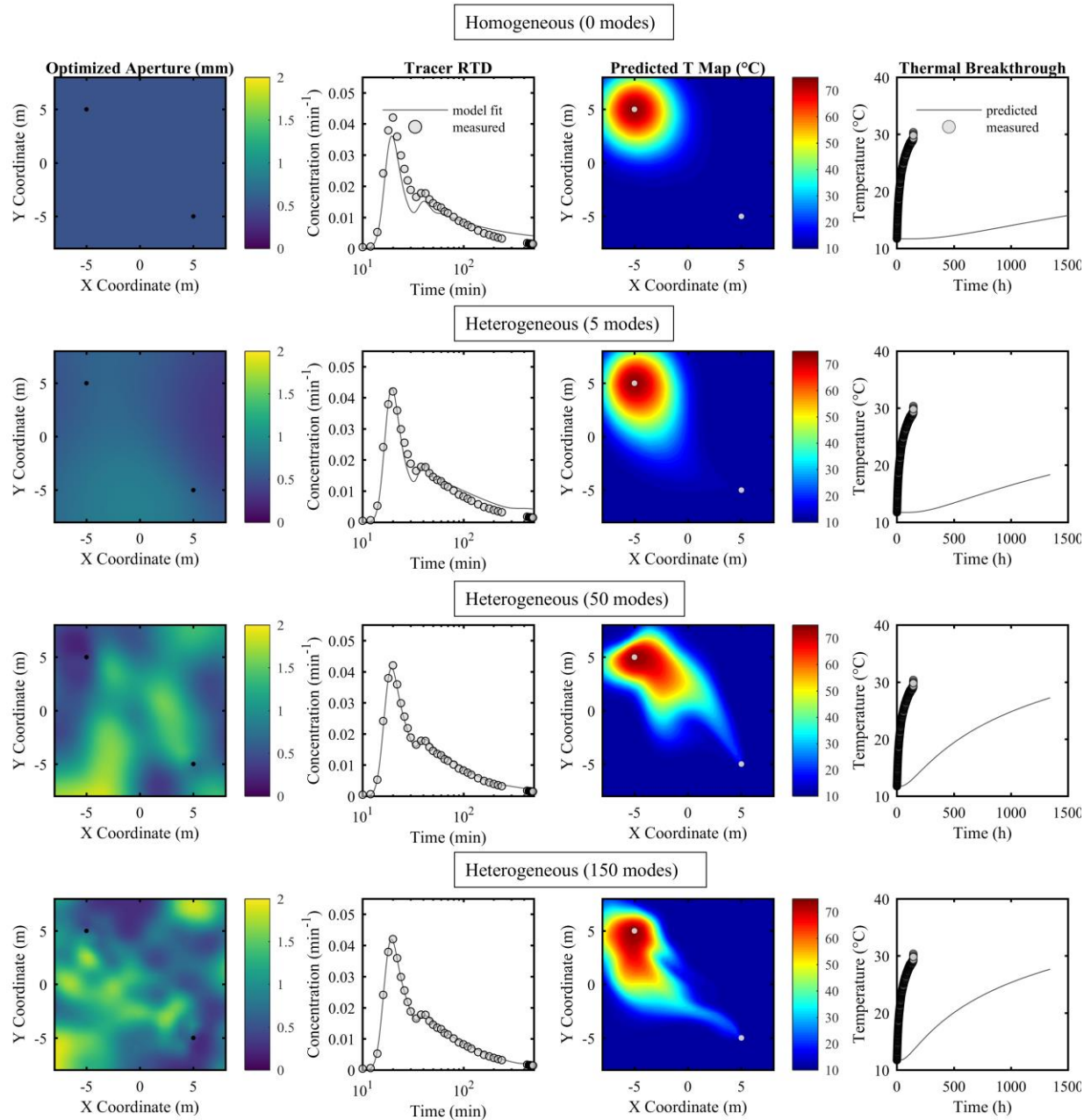383   base case is 0.22.

Figure 5. Preliminary results of the four scenarios prior to considering frictional pressure loss
(i.e. stage 1-1 to 1-4). Optimized aperture distributions are shown in the first column and their
resulting match to the measured inert tracer BTC are shown in the second column. Subsequent
predictions of advective heat transfer are shown in the third and fourth columns. The contour plot
in the third column shows the temperature distribution in the plane of the target fracture
predicted after 6 d of continuous hot water circulation and the fourth column compares predicted
and measured production well temperatures. The four scenarios include the base case
homogeneous aperture field (first row) and three heterogeneous cases in which the number of
modes solved by the GA varies from 5 (second row) to 150 (fourth row). Note that the bimodal
distributions of tracer breakthrough (second column) results from reinjection of produced fluids

395 from well 304 back into well 204. Tracer concentration is normalized by dividing the product of
396 concentration and volumetric flow rate by tracer mass injected. Measured production well
397 temperatures first published in Hawkins et al. (2017a).

398 Table 5. Final results of the two-stage optimization, preliminary results of the four scenarios
399 considered in the first stage, and predicted production well temperature, T_{prod}, after 6 d of hot
400 water circulation

Stage	Number of Modes Optimized	Mean Aperture (min, max) [mm]	L^2-norm	[a]ΔP [Pa]	[b]T_{prod} ($t = 6$ d) [°C]
2	150	2.5 [c](n/a, n/a)	0.11	-548	26.0
1-1	0 (homogeneous)	0.51 (n/a, n/a)	0.22	-11,191	11.70
1-2	5	0.61 (0.31, 0.87)	0.15	-6317	11.73
1-3	50	0.83 (0.16, 2.0)	0.06	-11,274	13.0
1-4	150	0.84 (0.24, 2.2)	0.04	-5636	13.3

401 [a]measured frictional pressure loss is -433 Pa (Table 6)
402 [b]measured temperature at the end of the 6 d heat exchange experiment is 29.3 °C (Hawkins et al., 2017a) (Figure 5)
403 [c]maximum and minimum aperture specified at fixed values of 0.02 and 8 mm, respectively

404 The second scenario considers a total of seven unknown variables which include the
405 minimum/maximum fracture aperture and the coefficients of the first five PCA modes. The GA
406 used a population of 144 and allowed up to 60 generations before terminating the optimization
407 routine. The optimization identified a minimum and maximum aperture of 0.31 and 0.87 mm,
408 respectively, which resulted in an L^2-norm of 0.15 and a noticeably improved match to the peak
409 tracer concentration (second row in Figure 5).

410 Two additional scenarios considered a larger number of PCA modes (50 and 150). With
411 50 PCA modes optimized, the GA identified a minimum and maximum aperture of 0.16 and 2.0
412 mm, respectively, and resulted in an L^2-norm of 0.06. With 150 modes optimized, the minimum
413 and maximum aperture identified are 0.24 and 2.2 mm, respectively, and resulted in an L^2-norm
414 of 0.04. For the scenarios in which the number of PCA modes were 5, 50, and 150, the resulting
415 mean of the optimized aperture field were 0.61, 0.83, and 0.84 mm, respectively.

416 Increasing the number of PCA modes identified by the GA from 0 to 5 resulted in a
417 nearly perfect match to the modal residence time and the peak tracer concentration and a minor
418 mismatch of the trough measured at ~30 min. With 50 modes, there is a nearly perfect overall
419 match between the model fit and measured breakthrough curves. With 150 modes, the roughness
420 of the permeability field is noticeably increased but the match to the inert tracer curve is
421 relatively unchanged compared to the 50 mode scenario.

422 Even though the third and fourth scenarios result in a nearly perfect match to the BTC
423 measured at the AFL, these aperture distributions do not guarantee accurate predictions of
424 advective heat transfer (Hawkins et al., 2017a). This is true because an inert tracer BTC provides
425 an estimation of interwell fluid volume whereas the extent of advective heat transfer is
426 dominated by the effective heat transfer surface area (Hawkins et al., 2018). As discussed in
427 Hawkins (2020), however, advective heat transfer can be anticipated if a model aperture field
428 provides good agreement with both frictional pressure loss and reservoir transit time.

429 The first four scenarios discussed above specified a fixed volumetric flow rate into the
430 injection well and out of the production well. The FEM solver subsequently calculated the
431 pressure field and velocity field. During the inversion routine, the objective function evaluated

432 the tracer fit, but gave no consideration to the agreement between measured and simulated
433 pressures in the injector or producer. Based on the pressure drop simulated by the permeability
434 fields shown in column 1 of Figure 5, however, there is considerable disagreement between the
435 simulated and measured frictional pressure loss between the injection and production well.
436 Simulated pressure loss between the injection well and production well for 0, 5, 50, and 150
437 modes are -11,911, -6317, -11,274, and -5636 Pa, respectively. In contrast, the pressure
438 measured in production well 304 is 433 Pa less than the pressure measured at injection well 204
439 (Table 6).

440 Table 6. Hydraulic head measurements from 2012 and parameter values used to calculate
441 frictional pressure loss

Symbol	Description	Value	Unit
Q	volumetric flow rate	7.00e-5	m^3/s
Δh_{204}	[a]injection well hydraulic head change at Q	-6.35e-3	m
Δh_{304}	[a]production well hydraulic head change at Q	-50.51e-3	m
g	gravitational acceleration	9.8	m/s^2
ρ_w	fluid density	1000	kg/m^3
ΔP	[b]calculated frictional pressure loss	-433	Pa

442 [a]hydraulic head measured via Solinst Levelloggers, Model #: 3001 and reported value is the average across 16 or more
443 measurements recorded over a time interval of at least 30 min.
444 [b]frictional pressure loss, ΔP, calculated as $\Delta P = \rho_w g(\Delta h_{304} - \Delta h_{204})$

445 Based on Equation 2, the relatively small pressure drop measured at the Altona site
446 suggests the true fracture aperture is larger than the aperture field identified in any of the four
447 scenarios considered during stage one of the optimization. Therefore, the second stage of the
448 optimization searched for a larger mean aperture by fixing the minimum and maximum apertures
449 at values of 0.02 and 8.0 mm, respectively. These boundaries were determined based on
450 manually varying the min/max apertures of the fourth scenario until the resulting aperture fields
451 provided good agreement to measured pressure loss.

452 With the inverse search forced towards larger apertures, stage 2 of the optimization
453 identified an aperture field with a mean of 2.5 mm and produced good agreement between
454 measured and simulated tracer RTD (L^2-norm is 0.11). The resulting RTD is presented in Figure
455 6 and shows relatively good agreement to measurement. The resulting pressure difference is -548
456 Pa and is in relatively good agreement with the measured value of -433 Pa compared to the
457 results of the first stage where simulated pressure losses ranged from -11,911 to -5636 Pa.
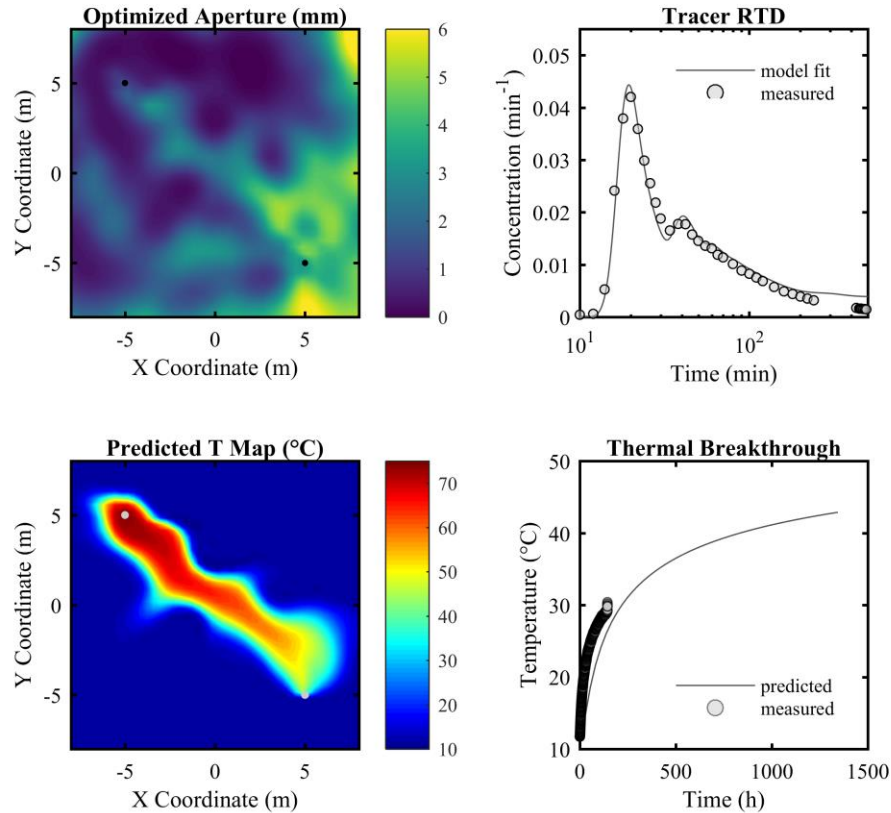
458

459 Figure 6. Optimized aperture field (top left) and the resulting simulation compared to measured
460 tracer breakthrough (top right) while considering both RTD and frictional pressure loss. The
461 bottom row shows the subsequent forward-model prediction of advective heat transport.
462 Predicted temperature simulated after 6 d of hot water injection is shown on the bottom left and a
463 comparison between measured and predicted thermal breakthrough is shown on the bottom right.

464   5.2 Temperature Predictions

465   Production well temperature predictions were made as described in subsection 3.3 for the
466 optimized aperture field (Figure 6) as well as the four preliminary aperture fields identified
467 during the first stage of the two-stage optimization (Figure 5). In this subsection, the results of
468 the final optimized aperture field will be discussed first followed by a comparison to predictions
469 resulting from the first stage of the optimization. A comparison between these simulations and
470 measured production well temperature rise at the AFL provide an evaluation of the inversion
471 routine's predictive capabilities.

472   With good agreement to both residence time distribution and frictional pressure loss, the
473 optimized aperture field predicted that production well fluid temperatures would rise 4 °C in 35
474 h. By the end of the 6 d heat exchange experiment, predicted production well temperature
475 reaches 26.0 °C. For comparison, measured production well temperature rises 4 °C in 7.3 h and
476 reaches 29.3 °C after 6 d.

477   Even though the L^2-norm of the final optimization is larger than the best results from the
478 stage 1 preliminary optimization (Table 5), the temperature predictions from stage 2 are

479 substantially more accurate than any of the four scenarios of stage 1. For all four scenarios of the
480 first stage the predicted thermal breakthrough times were substantially delayed relative to the
481 final result of the two-stage optimization. Despite the ideal match to RTDs shown in column 2 of
482 Figure 5, predicted production well temperatures for all four first-stage scenarios disagree with
483 the measured temperature recorded over the 6 d thermal test described in subsection 2.2. For the
484 homogeneous base case, predicted production well temperature rises 4 °C in ~61 d whereas
485 measured production fluids rise an equivalent temperature in just 7.3 h. The three subsequent
486 scenarios that solve for varying PCA modes (5, 50, and 150) reduces this disagreement, but
487 predictions are still off by at least a factor of 30. For the 5 and 50 mode scenarios, the predicted
488 time before production well fluids increase 4 °C is ~37 and ~11 d, respectively. For the 150
489 mode scenario, this predicted time is ~10 d.

490 **6 Conclusion**

491 We present a novel inverse model that identifies the spatial distribution of heterogeneous
492 fracture aperture/permeability based on calibration to measured inert tracer breakthrough curves
493 (BTCs). The approach used Principal Component Analysis (PCA) to force the inverse search to
494 investigate large-scale spatial fluctuations in fracture aperture. As a result, a Genetic Algorithm
495 (GA) efficiently solved the inverse problem and identifies a simplified aperture field that
496 provides good agreement to the Residence Time Distribution (RTD) and frictional pressure loss.

497 Meso-scale experiments conducted at the Altona Field Laboratory (AFL) demonstrated
498 that the method can identify an aperture field by minimizing the L^2-norm of simulated RTDs
499 compared to the breakthrough curve measured via the C-Dot nanoparticle tracer. With this
500 aperture field specified, a predictive forward model of advective heat transfer accurately
501 forecasted rapid thermal breakthrough between the same injector/producer well pair. This
502 channeled-flow behavior resulted from a narrow preferential flow path (~1-2 m) that caused an
503 inlet/outlet "short circuit." This result is in very good agreement with independent measurements
504 via Ground Penetrating Radar (GPR), a network of thermal sensors, and the results of an
505 adsorbing tracer test (Hawkins et al., 2017a, 2018).

506 Accurate predictions of heat transfer are only possible, however, if the aperture field
507 produces good agreement to both transit time and frictional pressure loss. In fact, even a uniform
508 (i.e., homogeneous) aperture field can provide reasonable agreement to the measured BTC while
509 providing a poor prediction of thermal transport. With a homogeneous aperture field, production
510 well temperature rises 4 °C in ~63 d whereas measured temperature rises 4 °C in just 7.3 h. In
511 contrast, the aperture field that provides good agreement to both frictional pressure loss and
512 residence time distribution predicts a 17.6 °C temperature rise in 8 d compared to a measured
513 temperature rise of 17.6 °C in 6 d. This result is consistent with the observation that a tracer RTD
514 alone cannot independently establish both volume and effective surface area of a one-
515 dimensional reservoir (Hawkins et al., 2018).

516 The results of this study have significant implications for tracer and hydraulic testing in
517 fractured rock reservoirs. Of particular importance is the observation that the RTD of an inert
518 tracer in a fracture is independent of the mean and range of fracture aperture magnitude. This
519 reinforces the conclusions of Hawkins et al. (2017a, 2018) which suggest that inert tracer testing
520 should not be considered as a means by which fracture aperture magnitude is estimated. Instead,
521 inert tracer testing provides an approximation of interwell fluid volume. Meanwhile, measured
522 frictional pressure loss does provide an approximation of mean fracture aperture. The predictive

523 model shown in Figure 6 provides reasonably accurate estimates of the thermal response because
524 the identified aperture field produces an interwell fluid volume in good agreement with tracer
525 tests and a mean aperture that is in good agreement with measured frictional pressure loss.

Acknowledgements

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