

1 **Temporal variation and frequency dependence of seismic ambient noise**
2 **on Mars from polarization analysis**

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19 **Key points:**

- 20
21 • A polarization analysis of InSight seismic data enables estimates of temporal
22 variation and frequency dependence of ambient noise on Mars.
23 • Back-azimuths of low-frequency (<1 Hz) P-waves and Rayleigh waves show diurnal
24 variations due to distant and nearby winds, respectively.
25 • The back-azimuth at high frequency points in the direction of the lander, indicating
26 that wind-induced lander noise is dominant.
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28

29 **Abstract**

30 We applied a polarization analysis of InSight seismic data to estimate the temporal
31 variation and frequency dependence of the Martian ambient noise field. Low-frequency
32 (<1 Hz) P-waves show a diurnal variation in their dominant back-azimuths that are
33 apparently related to wind and the direction of sunlight in a distant area. Low-frequency
34 Rayleigh waves (0.25–1 Hz) show diurnal variations and a dominant back-azimuth
35 related to the wind direction in a nearby area. Low-frequency signals that are derived
36 mainly from wind may be sensitive to subsurface structure deeper than the lithological
37 boundary derived from an autocorrelation analysis. On the other hand, dominant back-
38 azimuths of high-frequency (>1 Hz) waves point toward the InSight lander, especially in
39 daytime, indicating that wind-induced lander noise is dominant at high frequencies. These
40 results point to the presence of several ambient noise sources as well as geologic structure
41 at the landing site.

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43 **Plain Language Summary**

44 Seismic ambient noise (microtremors) is continuously generated not only on Earth but
45 also on Mars. We used data from the seismometer on the InSight lander to make estimates
46 of microtremor characteristics and identified possible underground structures that
47 influence the propagation of microtremors. Low-frequency P-waves derived from
48 microtremors show daily variations that appear to be induced by wind and changes of
49 sunlight during the Martian day in distant areas, whereas low-frequency Rayleigh waves
50 show daily variations that may be generated by wind in nearby areas. High-frequency
51 signals appear to originate from vibrations of the lander associated with wind.
52 Microtremors in other frequency ranges have different characteristics. These results
53 suggest that depending on their frequency, microtremors can be induced by wind and
54 other sources, and may then be influenced by geological structures. This study
55 demonstrates that ambient noise data will be helpful for imaging and monitoring Mars'
56 interior structure and natural resources, such as ice deposits, without the need for data
57 from marsquakes and artificial seismic sources.

58

59 **Keywords:** InSight, ambient noise, polarization analysis, autocorrelation function, wind

60

61 **1. Introduction**

62 When NASA's Interior Exploration using Seismic Investigations, Geodesy and
63 Heat Transport (InSight) lander touched down in Elysium Planitia on 26 November 2018,
64 it went on to deploy a geophysical observatory on Mars. One of its primary scientific
65 investigations is the Seismic Experiment for Interior Structure (SEIS; InSight Mars SEIS
66 Data Service, 2019; Lognonné et al., 2019). The lander also includes a set of
67 environmental sensors, including temperature and wind sensors (Banfield et al., 2019;
68 Spiga et al., 2018). The InSight seismometer has detected several hundred marsquakes,
69 most of them much smaller than earthquakes typically felt on Earth, but some were nearly
70 as large as magnitude 4 (Witze, 2019). The instrument is especially useful to identify
71 small earthquakes at night, when the strong ambient noise generated during the day by
72 wind is subdued (Witze, 2019). Martian ambient noise detected by the seismometer on
73 the Viking 2 lander has been correlated with wind speed (Anderson et al. 1977; Nakamura
74 and Anderson, 1979). However, that seismometer did not obtain seismic signals directly
75 because it was deployed on the lander and not on the surface (Knapmeyer-Endrun et al.,
76 2017).

77 Analysis of seismic ambient noise is a technique widely used on Earth to image
78 and monitor the subsurface (e.g., Nimiya et al., 2017; Nishida et al., 2008), and several
79 studies have made similar use of ambient noise on the Moon (e.g., Larose et al., 2005;
80 Tanimoto et al., 2008). If ambient noise can be used to image and monitor the interior
81 structures of Mars, this technique will be a powerful tool because it does not require any
82 natural marsquakes or expensive artificial seismic sources.

83 In this paper, we characterize the ambient noise on Mars relying on the recent
84 data from the InSight seismometer. We applied a polarization analysis to the InSight
85 seismic records (InSight Mars SEIS Data Service, 2019) to extract the dominant back-
86 azimuth and directional intensity of ambient noise (Takagi et al., 2018). Furthermore, by
87 comparing the characteristics of Rayleigh waves with autocorrelation functions (i.e.,
88 reflectivity), we achieved insight into the relationship between lithology and the sensitive
89 frequency of Rayleigh waves included in ambient noise. By demonstrating the feasibility
90 of ambient noise methods on Mars, this study shows that future seismic network projects

91 on Mars will contribute to not only modeling and monitoring of Mars' interior structure,
92 but also exploration for Martian resources (e.g., ice deposits).

93

94 **2. Data and Method**

95 **2.1. Data Preparation**

96 The SEIS instrument includes a long-period, very broad band seismometer
97 (SEIS-VBB) with a sampling rate of 20 Hz and a natural frequency of 0.5 Hz (Lognonné
98 et al., 2019; InSight Mars SEIS Data Service, 2019). This seismometer was placed in
99 Elysium Planitia in particular to satisfy the constraints on landing safety and the
100 instrument deployment requirements (Golombek et al., 2017). In this study, we used
101 continuous seismic records from SEIS-VBB between February and June 2019. We
102 corrected the data for the instrumental response using ObsPy (Beyreuther et al., 2010).
103 The SEIS-VBB is a triaxial seismometer in which the three mutually perpendicular
104 pendulums are mounted obliquely. Therefore, our first step was to numerically rotate the
105 axes of the seismometer and construct seismic records with vertical and horizontal
106 components ([see Text S1 in the supporting information](#)).

107 We then converted the seismic data from Earth time (UTC; Coordinated
108 Universal Time) to the Mars time domain (LMST: Local Mean Solar Time) by using the
109 procedures of Allison (1997) and Allison and McEwen (2000). The power spectral density
110 on Mars calculated from ambient noise shows that the noise on Mars is lower at most
111 frequencies than that of the Earth noise model ([see Fig. S1 in the supporting information](#)).
112 The power spectral densities of the horizontal and vertical components from Sols 194 to
113 197 (**Fig. 1**) are an example of the typical daily cycle, in which signal amplitudes are
114 greater during the day than during the night. On Mars, high variability of wind in daytime
115 is caused by convective mixing in the planetary boundary layer that results from near-
116 surface gradients of atmospheric temperature (e.g., Smith et al., 2006; Spiga et al., 2018).
117 At frequencies higher than ~ 1 Hz, we observed large noise amplitudes in narrow
118 frequency ranges. These local noise peaks correspond to the elastic resonances of the
119 lander excited by the wind (Murdoch et al., 2017; Lognonné et al., 2020). These results
120 demonstrate that the amplitude of ambient noise is strongly associated with the wind
121 strength.

122 We divided continuous seismic data into 1-min segments because short time

123 windows are suitable to remove glitches and other high-amplitude signals (Takagi et al.,
 124 2018). We excluded time segments whose root-mean-squared (RMS) amplitudes
 125 exceeded 10 times the median RMS amplitude, treating daytime hours (from 6:00 to
 126 18:00 LMST) and nighttime hours (from 18:00 to 6:00 LMST) separately because the
 127 surface wind velocity was high during the daytime at the InSight landing site (**Fig. 1**) as
 128 anticipated by Spiga et al. (2018).

129

130 **2.2. Polarization Analysis**

131 We conducted a polarization analysis of the ambient seismic wave field recorded
 132 by the InSight station using the method developed by Takagi et al. (2018). This analysis
 133 uses a simple relationship between the vertical-horizontal cross spectra and the azimuthal
 134 energy distributions of incident waves in ambient noise. The real part of the cross spectra
 135 is related to linearly polarized waves and the imaginary part is related to elliptically
 136 polarized waves. We computed vertical-horizontal cross spectra from 1-min segments
 137 data using the equations

138

$$\Phi_{ZN} = \frac{u_Z^* u_N}{u_Z^* u_Z}, \quad (1)$$

$$\Phi_{ZE} = \frac{u_Z^* u_E}{u_Z^* u_Z}, \quad (2)$$

139

140 where Φ is the vertical-horizontal cross spectrum, u is the seismic record in the
 141 frequency domain of each component, and the subscripts Z , N and E indicate vertical,
 142 north-south and east-west component, respectively. The asterisk indicates the complex
 143 conjugate. The cross spectra are normalized by the power spectra of the vertical
 144 component so as to equally weight each data segment. In this study, the cross spectra were
 145 calculated at each frequency and the results were averaged within each of six single-
 146 octave frequency bands: 0.125–0.25, 0.25–0.5, 0.5–1, 1–2, 2–4 and 4–8 Hz.

147 Following Takagi et al. (2018), the dominant direction and directional intensity
 148 of a Rayleigh wave (elliptically polarized wave) are given by

149

$$\varphi_{R1} = \arctan\left(\frac{\text{Im}\langle\Phi_{ZE}\rangle}{\text{Im}\langle\Phi_{ZN}\rangle}\right) + \pi, \quad (3)$$

$$A_{R1} = \sqrt{(\text{Im}\langle\Phi_{ZN}\rangle)^2 + (\text{Im}\langle\Phi_{ZE}\rangle)^2}, \quad (4)$$

150

151 and for a P-wave (linearly polarized wave) by

152

$$\varphi_{P1} = \arctan\left(\frac{\text{Re}\langle\Phi_{ZE}\rangle}{\text{Re}\langle\Phi_{ZN}\rangle}\right) + \pi, \quad (5)$$

$$A_{P1} = \sqrt{(\text{Re}\langle\Phi_{ZN}\rangle)^2 + (\text{Re}\langle\Phi_{ZE}\rangle)^2}, \quad (6)$$

153

154 where $\langle \rangle$ denotes the ensemble average and φ_{R1} and φ_{P1} represent the phase angles of
 155 first-order terms of the azimuthal power spectra added to π , which provide the dominant
 156 back-azimuths of Rayleigh waves and P-waves, respectively. A_{R1} and A_{P1} indicate the
 157 amplitudes of the first-order terms representing the intensity of the directionality of the
 158 Rayleigh wave and P-wave, respectively.

159

In the determination of Rayleigh wave azimuth, there is a 180-degree ambiguity
 160 depending on the direction of motions (prograde or retrograde). To evaluate the motion
 161 of Rayleigh waves on Mars, we computed analytical solutions of Rayleigh waves for the
 162 layered model of Knapmeyer-Endrun et al. (2017) of the InSight landing site (see [Text S2](#)
 163 [and Fig. S2 in the supporting information](#)). The results indicate that the fundamental mode
 164 of Rayleigh waves with retrograde motions is mostly dominant in our analyzed frequency
 165 range, whereas the first higher mode with prograde motions is dominant at some
 166 frequencies higher than 4 Hz. We therefore defined the azimuth of Rayleigh waves
 167 assuming retrograde motions, although the first higher mode with prograde motions might
 168 influence our results at frequencies higher than 4 Hz. Note that the azimuth of prograde
 169 or retrograde Rayleigh waves depends on the sign of the exponent in the Fourier transform.
 170 We used equation (3) to estimate the back-azimuth of retrograde Rayleigh waves because
 171 our analysis used the Fourier transform with a negative exponent. Shear waves with
 172 displacement in the vertical-horizontal plane (SV-waves) also contribute to vertical-
 173 horizontal cross-spectra (Takagi et al., 2018). Vertically incident SV-waves contribute to
 174 the real part of the vertical-horizontal cross spectra, whereas horizontally incident SV-
 175 waves with post-critical incident angles contribute to the imaginary part. For simplicity,

176 we assumed that the contribution of P-waves is dominant in the real part of the cross
177 spectra and the contribution of Rayleigh waves is dominant in the imaginary part. Under
178 the assumption that Rayleigh and Love waves are random uncorrelated waves, Love
179 waves make no contribution to vertical and horizontal cross spectra (Takagi et al., 2018).

180

181 **2.3. Autocorrelation Analysis**

182 To estimate the geological structure beneath the InSight landing site, we applied
183 autocorrelation analysis to the vertical and horizontal motions of the seismometer record.
184 Autocorrelation of ambient noise records yields the zero-offset shot gather (e.g., Minato
185 et al., 2012; Wapenaar & Fokkema, 2006). The method assumes that the noise source is
186 randomly distributed and mutually uncorrelated for different source positions (e.g., Roux
187 et al., 2005; Wapenaar and Fokkema, 2006; Weaver & Lobkis, 2004). In this analysis, we
188 applied a bandpass filter of 5–7 Hz to each component record of 1-min segments, because
189 we found clear reflectors of autocorrelation function in that frequency band. Furthermore,
190 we sought to find and integrate information independent of the polarization analysis for
191 the investigation of the lander site. We applied one-bit normalization (e.g., Bensen et al.,
192 2007) to ensure the exclusion of energetic signals. We calculated autocorrelation
193 functions of the vertical component and the horizontal components in each sol to extract
194 P- and S-wave reflections, respectively. Even if the lander near the seismometer generates
195 vibration and becomes a noise source, the autocorrelation analysis with one-bit
196 normalization reduces the influence of the source but enhances the contribution of
197 reflected waves from the source. Thus, we expect that autocorrelation analysis is suitable
198 for subsurface imaging.

199

200 **3. Results**

201 **Figures 2a** and **2b** show the temporal variations of dominant back-azimuths and
202 directional intensity of P-waves and Rayleigh waves from Sols 75 to 211 in the six
203 frequency bands. The cross spectra are averaged for each hour in the Mars time domain
204 (LMST). The dominant back-azimuths were different for each frequency band. The
205 directional intensity of Rayleigh waves was less than that of P-waves in all frequency
206 bands.

207 To illustrate the daily temporal variation, we present results from Sols 194 to 197
208 (**Fig. 3**). For low-frequency P-waves (<1 Hz), the back-azimuths shifted between east and
209 north during the course of the day (**Fig. 3a**). The back-azimuths at the lowest frequencies
210 pointed southeast in daytime, roughly consistent with the wind direction. At night, the
211 back-azimuths of 0.25–1 Hz P-waves usually pointed east, except just after sunset; more
212 precisely, they differed notably from the wind direction at night, pointing east several
213 hours before sunrise and pointing west to north after sunset. For 0.25–1 Hz Rayleigh
214 waves (**Fig. 3b**), the back-azimuth pointed southwest before sunrise, south or southeast
215 during the day, and southwest at night, similar to the wind direction.

216 At high frequencies (>1 Hz), the dominant back-azimuths of P-waves (**Fig. 3a**)
217 pointed northeast in daytime, as did the back-azimuths of high frequency Rayleigh waves
218 (>2 Hz). As we discuss later, the Insight lander is located northeast of the seismometer.

219 **Figure 4** shows the temporal variation of the autocorrelation function during the
220 observation period. The autocorrelation function of the vertical component (**Fig. 4a**)
221 indicates the presence of reflectors at 0.6 s and 1.1 s. Because these reflectors persisted
222 throughout the observation period, they appear to be reliable and may represent a
223 lithological boundary that imposes a contrast in acoustic impedance. The autocorrelation
224 functions of the two horizontal components (**Figs. 4b and 4c**) display dominant reflectors
225 at 1.1 s. They show evidence of polarization anisotropy of S-waves, in that the reflector
226 at ~ 1.1 s is more prominent in the EW component (**Fig. 4c**) than in the NS component
227 (**Fig. 4b**).

228

229 **4. Discussion**

230 The temporal variation of the dominant back-azimuth of <1 Hz P-waves could
231 be related to the direction of sunlight (or related thermal effects) in addition to the wind
232 direction when noise derived from the lander is absent. During the several hours before
233 sunrise, the area east of the lander site is in daylight and the wind speed is high, thus the
234 dominant P-wave back-azimuth could point east before sunrise (**Fig. 3a**). This
235 interpretation would also explain the westward P-wave back-azimuth after sunset,
236 although the back-azimuths are scattered from west to north. These results demonstrate
237 that low-frequency P-waves observed at the InSight site may be derived from wind and

238 insolation effects (e.g., thermal cracking) in distant areas. Indeed, P-waves on Earth are
239 strongly influenced by distant events (Takagi et al., 2018). Seismic sources induced by
240 temperature variation are capable of generating low-frequency ambient noise.

241 Because the variation of the directionality of 0.25–1 Hz Rayleigh waves was
242 closely related to the wind direction (**Fig. 3b**), low-frequency Rayleigh waves were likely
243 derived from winds relatively close to the seismometer. The back-azimuth of Rayleigh
244 waves could be influenced by the radiation pattern of Rayleigh waves. Assuming that
245 horizontal single forces exerted in the wind directions on rough surface topography
246 excited seismic waves including Rayleigh waves, a symmetric radiation pattern (i.e., with
247 180-degree ambiguity) could be expected in the back-azimuths (**Fig. 3b**). Although
248 stacking the cross spectra for each hour improves the stability of estimated dominant
249 back-azimuths (**Fig. 3b**), using shorter time windows could make it possible to extract
250 secondary dominant back azimuths. To investigate this possibility, we computed the back-
251 azimuth from every 1-min segment (**Fig. S3 in supporting information**). The results of
252 this exercise show that directionalities of 0.125–1 Hz Rayleigh waves have two trends
253 180 degrees apart during certain periods; thus, the radiation pattern of Rayleigh waves
254 could influence the observed back-azimuth. Rayleigh waves in the 0.25–1 Hz range
255 would be sensitive to the depth range of 0.8–3.2 km, if we assume a Rayleigh wave
256 velocity of 2400 m/s (Knapmeyer-Endrun et al., 2017). Therefore, the wind may be
257 responsible for 0.25–1 Hz Rayleigh waves that are sensitive to the crustal structure
258 beneath the shallow regolith layer.

259 At high frequencies, the back-azimuths of P-waves >1 Hz and Rayleigh waves
260 >2 Hz are northeast in daytime. The direction is consistent with the location of the InSight
261 lander (**Fig. 3**), which generates mechanical noise as wind acts on the lander (Murdoch et
262 al., 2017; Lognonné et al., 2020). If wind-induced lander noise is dominant at high
263 frequencies, it would be difficult to observe high-frequency Rayleigh waves with the
264 seismometer because the distance between the lander and the seismometer is too short
265 (several meters) for surface waves to emerge. Therefore, instead of referring to “P- and
266 Rayleigh waves” in high-frequency (>1 Hz) results, it is preferable to refer to “linearly
267 and elliptically polarized components of observed waves” as we have in **Figs. 2 and 3**.

268 These frequency-dependent variations of ambient noise characteristics could be

269 mainly related to ambient noise sources and lithology beneath the seismometer. Ambient
270 noise on Earth is caused by wind (Lepore et al., 2016) as well as ocean gravity waves,
271 volcanic activity, and anthropogenic sources (e.g., Longuet-Higgins, 1950; Takagi et al.,
272 2018; Nimiya et al., 2017; Nakata et al., 2019). Before the InSight project, a main source
273 of ambient noise on Mars was expected to be the direct interaction between the
274 atmosphere and the solid surface of the planet (Knapmeyer-Endrun et al., 2017). On the
275 Moon, high-frequency Rayleigh waves are induced by ambient noise resulting from
276 thermal events (Larose et al., 2005; Tanimoto et al., 2008). On Mars, there are numerous
277 small craters near the InSight landing site (Warner et al., 2016) that could be locations of
278 thermally triggered soil slumping (Knapmeyer-Endrun et al., 2017) that could generate
279 high-frequency surface waves. Thus wind-induced noises, thermal effects, surface
280 pressure, or other sources may induce the ambient noise around the InSight landing area.

281 To further consider the relationship between the frequency dependence of
282 Rayleigh waves (**Fig. 3b**) and the lithology of the site, we investigated the autocorrelation
283 results (**Fig. 4**), in which several reflectors beneath the InSight landing site are evident.
284 The P-wave reflectors at 0.6 and 1.1 s in the vertical component (**Fig. 4a**) are stable,
285 suggesting the existence of a significant lithological boundary. Furthermore, an S-wave
286 reflector appeared at 1.1 s in the horizontal component results (**Fig. 4b and 4c**). If the 1.1
287 s S-wave reflector is the same as the 0.6 s P-wave reflector, we can estimate the ratio of
288 ~ 1.83 between the P-wave and S-wave velocities. Because we cannot estimate the seismic
289 velocity of the subsurface formation, we cannot accurately estimate the depth of the
290 reflectors from the autocorrelation functions. However, we can estimate the frequency of
291 Rayleigh waves that are sensitive to the depth of a reflector from the autocorrelation
292 function. Under the assumption that the autocorrelation function of the horizontal
293 component represents S-wave reflectivity, the depth of a reflector at two-way travel time
294 t can be estimated as $Z = t V_S/2$, where V_S is S-wave velocity. The sensitive depth of
295 Rayleigh waves is $Z = 1/3 \lambda$ (or $Z = V_S/3f$) (e.g., Foti et al., 2014; Hayashi, 2008), where
296 λ is wavelength and f is frequency. Therefore, the sensitive frequency of a Rayleigh
297 wave for a reflector at two-way travel time t can be estimated as $f = 2/(3t)$. From this
298 relationship, the frequency of a Rayleigh wave that is sensitive to a 1.1 s reflector shown
299 in **Figs. 4b and 4c** can be estimated as ~ 0.6 Hz. Below 0.6 Hz, the azimuths of Rayleigh

300 waves are associated with wind direction. Therefore, Rayleigh waves that are sensitive to
301 depths beneath the lithological boundary identified by reflectivity could be extracted from
302 wind-induced ambient noise. However, it would be difficult to extract Rayleigh waves
303 propagating above the lithological boundary close to the landing site, because they are
304 contaminated by lander-induced noise.

305

306 **5. Conclusions**

307 We have conducted a polarization analysis of InSight seismic data to estimate
308 temporal variations of the ambient noise field on Mars. Our findings are these:

309

- 310 • Low-frequency (<1 Hz) P-waves show a diurnal variation, and the dominant back-
311 azimuth is related to the wind and the direction of sunlight in distant regions.
- 312 • Low-frequency (0.25–1 Hz) Rayleigh waves show a diurnal variation, and the
313 dominant back-azimuth points toward the wind direction in nearby regions.
- 314 • The dominant back-azimuth at high-frequency (>1 Hz for linearly polarized
315 components and >2 Hz for elliptically polarized components) points in the direction
316 of the lander, indicating that the wind-induced lander noise is dominant.

317

318 These results suggest that the dominant sources of ambient noise on Mars differ
319 with frequency and wave type, and there may be several different ambient noise sources
320 despite the absence of oceans on Mars. The high repeatability of P-waves and Rayleigh
321 waves derived from ambient noise suggests the feasibility of utilizing ambient noise for
322 subsurface imaging and monitoring on Mars. Further studies are necessary to clarify the
323 contribution of SV-waves in ambient noise on Mars, which influences the results of our
324 polarization analysis.

325

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330 seismometer data from IRIS-DMC. The wind and temperature data were downloaded

331 from the following URL;

332 [https://atmos.nmsu.edu/data_and_services/atmospheres_data/INSIGHT/insight.html#Se-](https://atmos.nmsu.edu/data_and_services/atmospheres_data/INSIGHT/insight.html#Selecting_Data)
 333 [lecting_Data.](https://atmos.nmsu.edu/data_and_services/atmospheres_data/INSIGHT/insight.html#Selecting_Data)

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340 **References**

341 Allison, M. (1997). Accurate analytic representations of solar time and seasons on Mars
 342 with applications to the Pathfinder/Surveyor missions. *Geophysical Research*
 343 *Letters*. <https://doi.org/10.1029/97GL01950>

344 Allison, M., & McEwen, M. (2000). A post-Pathfinder evaluation of areocentric solar
 345 coordinates with improved timing recipes for Mars seasonal/diurnal climate
 346 studies. *Planetary and Space Science*, *48*(2–3), 215–235.
 347 [https://doi.org/10.1016/s0032-0633\(99\)00092-6](https://doi.org/10.1016/s0032-0633(99)00092-6)

348 Anderson, D. L., Miller, W. F., Latham, G. V., Nakamura, Y., Toksöz, M. N., Dainty,
 349 A. M., et al. (1977). Seismology on Mars. *Journal of Geophysical Research*,
 350 *82*(28), 4524–4546. <https://doi.org/10.1029/JS082i028p04524>

351 Banfield, D., Rodriguez-Manfredi, J. A., Russell, C. T., Rowe, K. M., Leneman, D.,
 352 Lai, H. R., et al. (2019). InSight Auxiliary Payload Sensor Suite (APSS). *Space*
 353 *Science Reviews* (Vol. 215). *Springer Nature B.V.* [https://doi.org/10.1007/s11214-](https://doi.org/10.1007/s11214-018-0570-x)
 354 [018-0570-x](https://doi.org/10.1007/s11214-018-0570-x)

355 Bensen, G. D., Ritzwoller, M. H., Barmin, M. P., Levshin, A. L., Lin, F., Moschetti, M.
 356 P., et al. (2007). Processing seismic ambient noise data to obtain reliable broad-
 357 band surface wave dispersion measurements. *Geophysical Journal International*,
 358 *169*(3), 1239–1260. <https://doi.org/10.1111/j.1365-246X.2007.03374.x>

359 Beyreuther, M., R. Barsch, L. Kischer, T. Megies, Y. Behr, & J. Wassermann (2010).
 360 ObsPy: A Python toolbox for seismology, *Seismol. Res. Lett.* *81*, no. 3, 530–533.
 361 <https://doi.org/10.1785/gssrl.81.3.530>

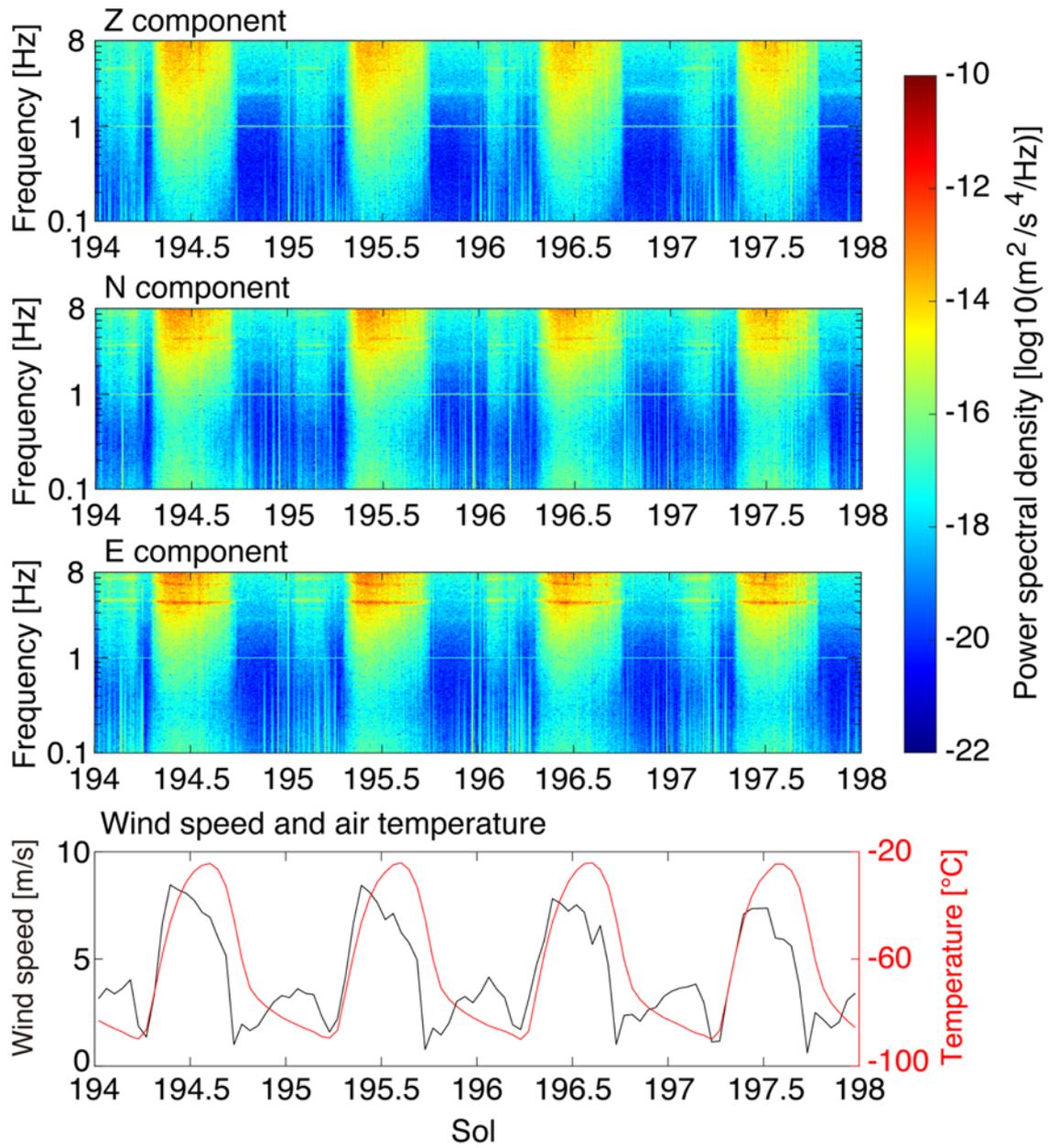
362 Foti, S., Lai, G. C., Rix, G. J., & Strobbia, C. (2014). Surface wave methods for near-
 363 surface site characterization. London: CRC Press. <https://doi.org/10.1201/b17268>

364 Golombek, M., Kipp, D., Warner, N., Daubar, I. J., Ferguson, R., Kirk, R. L., et al.
 365 (2017). Selection of the InSight Landing Site. *Space Science Reviews*, *211*(1–4), 5–
 366 [95. https://doi.org/10.1007/s11214-016-0321-9](https://doi.org/10.1007/s11214-016-0321-9)

- 367 Harkrider, D. G. (1970). Surface waves in multilayered elastic media. Part II. Higher
368 mode spectra and spectral ratios from point sources in plane layered earth models.
369 *Bull. Seismol. Soc. Am.* 60, 1937–1987.
- 370 Hayashi, K. (2008). Development of surface-wave methods and its application to site
371 investigations: Ph.D. dissertation, Kyoto University.
- 372 InSight Mars SEIS Data Service. (2019). SEIS raw data, Insight Mission. IPGP, JPL,
373 CNES, ETHZ, ICL, MPS, ISAE-Supaero, LPG, MFSC.
374 https://doi.org/10.18715/SEIS.INSIGHT.XB_2016
- 375 Knapmeyer-Endrun, B., Golombek, M. P., & Ohrnberger, M. (2017). Rayleigh Wave
376 Ellipticity Modeling and Inversion for Shallow Structure at the Proposed InSight
377 Landing Site in Elysium Planitia, Mars. *Space Science Reviews*, 211(1–4), 339–
378 382. <https://doi.org/10.1007/s11214-016-0300-1>
- 379 Larose, E., Khan, A., Nakamura, Y., & Campillo, M. (2005). Lunar subsurface
380 investigated from correlation of seismic noise. *Geophysical Research Letters*,
381 32(16), 1–4. <https://doi.org/10.1029/2005GL023518>
- 382 Lepore, S., Markowicz, K., & Grad, M. (2016). Impact of wind on ambient noise
383 recorded by seismic array in northern Poland. *Geophysical Journal International*,
384 205(3), 1406–1413. <https://doi.org/10.1093/gji/ggw093>
- 385 Lognonné, P., Banerdt, W. B., Giardini, D., Pike, W. T., Christensen, U., Laudet, P., et
386 al. (2019). SEIS: Insight’s Seismic Experiment for Internal Structure of Mars.
387 *Space Science Reviews* (Vol. 215). <https://doi.org/10.1007/s11214-018-0574-6>
- 388 Lognonné, P., Banerdt, W. B., Pike, W. T., Giardini, D., Christensen, U., Garcia, R. F., et
389 al. (2020). Constraints on the shallow elastic and anelastic structure of Mars from
390 InSight seismic data. *Nature Geoscience*, 13(3), 213–220.
391 <https://doi.org/10.1038/s41561-020-0536-y>
- 392 Longuet-Higgins, M., 1950. A theory of the origin of microseisms, *Phil. Trans. R. Soc.*,
393 243, 1–35.
- 394 Minato, S., Tsuji, T., Ohmi, S., & Matsuoka, T. (2012). Monitoring seismic velocity
395 change caused by the 2011 Tohoku-oki earthquake using ambient noise records.
396 *Geophysical Research Letters*, 39(9), 1–6. <https://doi.org/10.1029/2012GL051405>
- 397 Murdoch, N., Mimoun, D., Garcia, R. F., Rapin, W., Kawamura, T., Lognonné, P., et al.
398 (2017). Evaluating the Wind-Induced Mechanical Noise on the InSight
399 Seismometers. *Space Science Reviews*, 211, 429–455.
- 400 Nakamura, Y. & Anderson, D. L. (1979). Martian wind activity detected by a
401 seismometer at Viking Lander 2 site. *Geophys. Res. Lett.*, 6: 499–502.
402 <https://doi.org/10.1029/GL006i006p00499>

- 403 Nakata, N., Gualtieri, L., & Fichtner, A. (2019). *Seismic Ambient Noise*, Cambridge
404 University Press, Cambridge, UK.
- 405 Nimiya, H., Ikeda, T., & Tsuji, T. (2017). Spatial and temporal seismic velocity changes
406 on Kyushu Island during the 2016 Kumamoto earthquake. *Science Advances*,
407 3(11). <https://doi.org/10.1126/sciadv.1700813>
- 408 Nishida, K., Kawakatsu, H., & Obara, K. (2008). Three-dimensional crustal S wave
409 velocity structure in Japan using microseismic data recorded by Hi-net tiltmeters.
410 *Journal of Geophysical Research: Solid Earth*, 113(10), 1–22.
411 <https://doi.org/10.1029/2007JB005395>
- 412 Peterson, J. (1993). Observations and modeling of seismic background noise, *U.S. Geol.*
413 *Surv. Open File Report*, 93-322. <https://doi.org/10.3133/ofr93322>
- 414 Roux, P., Sabra, K. G., Kuperman, W. A., & Roux, A. (2005). Ambient noise cross
415 correlation in free space: Theoretical approach. *The Journal of the Acoustical*
416 *Society of America*, 117(1), 79–84. <https://doi.org/10.1121/1.1830673>
- 417 Saito, M., 1988. DISPER80: a subroutine package for calculation of seismic normal -
418 mode solution. in *Seismological Algorithm*, ed. Doorbos, D.J., Academic Press,
419 San Diego , CA .
- 420 Smith, M. D., Wolff, M. J., Spanovich, N., Ghosh, A., Banfield, D., Christensen, P. R.,
421 et al. (2006). One Martian year of atmospheric observations using MER Mini -
422 TES. *Journal of Geophysical Research*, 111, E12S13.
423 <https://doi.org/10.1029/2006JE002770>
- 424 Spiga, A., Banfield, D., Teanby, N. A., Forget, F., Lucas, A., Kenda, B., et al. (2018).
425 Atmospheric Science with InSight. *Space Science Reviews* (Vol. 214). Springer
426 Nature B.V. <https://doi.org/10.1007/s11214-018-0543-0>
- 427 Takagi, R., Nishida, K., Maeda, T., & Obara, K. (2018). Ambient seismic noise
428 wavefield in Japan characterized by polarization analysis of Hi-net records.
429 *Geophysical Journal International*, 215(3), 1682–1699.
430 <https://doi.org/10.1093/gji/ggy334>
- 431 Tanimoto, T., Eitzel, M. V., & Yano, T. (2008). The noise cross-correlation approach
432 for Apollo 17 LSPE data: Diurnal change in seismic parameters in shallow lunar
433 crust. *Journal of Geophysical Research E: Planets*, 113(8), 1–12.
434 <https://doi.org/10.1029/2007JE003016>
- 435 Tokimatsu, K. (1997). Geotechnical site characterization using surface waves. in
436 *Earthquake Geotechnical Engineering*, K. Ishihara (Editor), A. Balkema,
437 Rotterdam, The Netherlands, 1333–1368.
- 438 Wapenaar, K., & Fokkema, J. (2006). Green's function representations for seismic

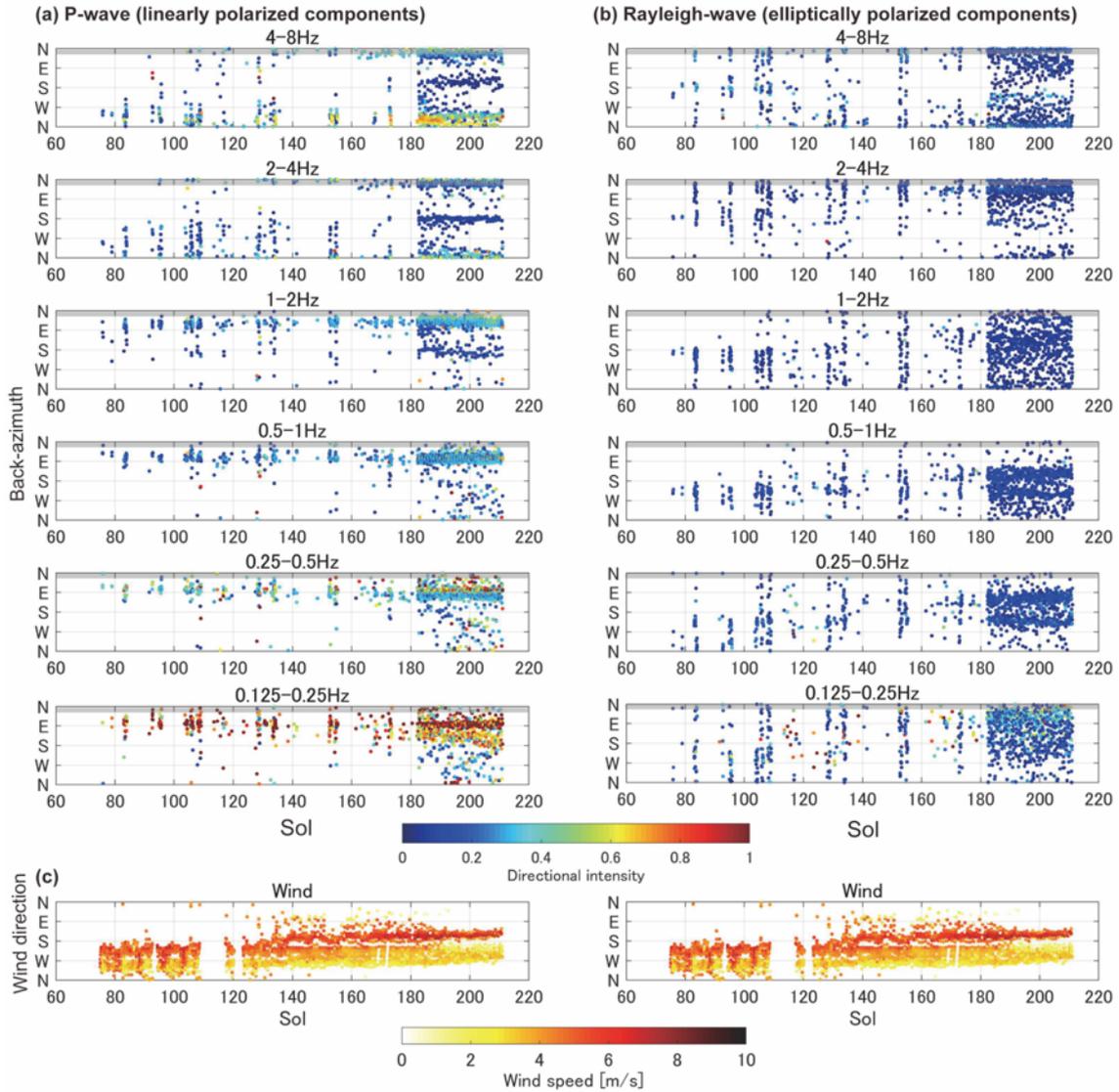
- 439 interferometry. *Geophysics*, 71(4). <https://doi.org/10.1190/1.2213955>
- 440 Warner, N. H., Golombek, M. P., Sweeney, J., Pivarunas, A. (2016). Regolith thickness
441 estimates from the size frequency distribution of rocky ejecta craters in
442 southwestern Elysium Planitia, Mars, in *47th Lunar and Planetary Science*
443 *Conference*, Abstract 2231.
- 444 Weaver, R. L., & Lobkis, O. I. (2004). Diffuse fields in open systems and the
445 emergence of the Green's function (L). *The Journal of the Acoustical Society of*
446 *America*, 116(5), 2731–2734. <https://doi.org/10.1121/1.1810232>
- 447 Witze, A. (2019). “Marsquakes” reveal red planet’s hidden geology. *Nature*, 576(7787),
448 348. <https://doi.org/10.1038/d41586-019-03796-7>
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452 Fig. 1. Temporal variation of power spectral density in the vertical and two horizontal
 453 components from Sols 194 to 197. The bottom figure shows the temporal variation of
 454 wind speed and air temperature.

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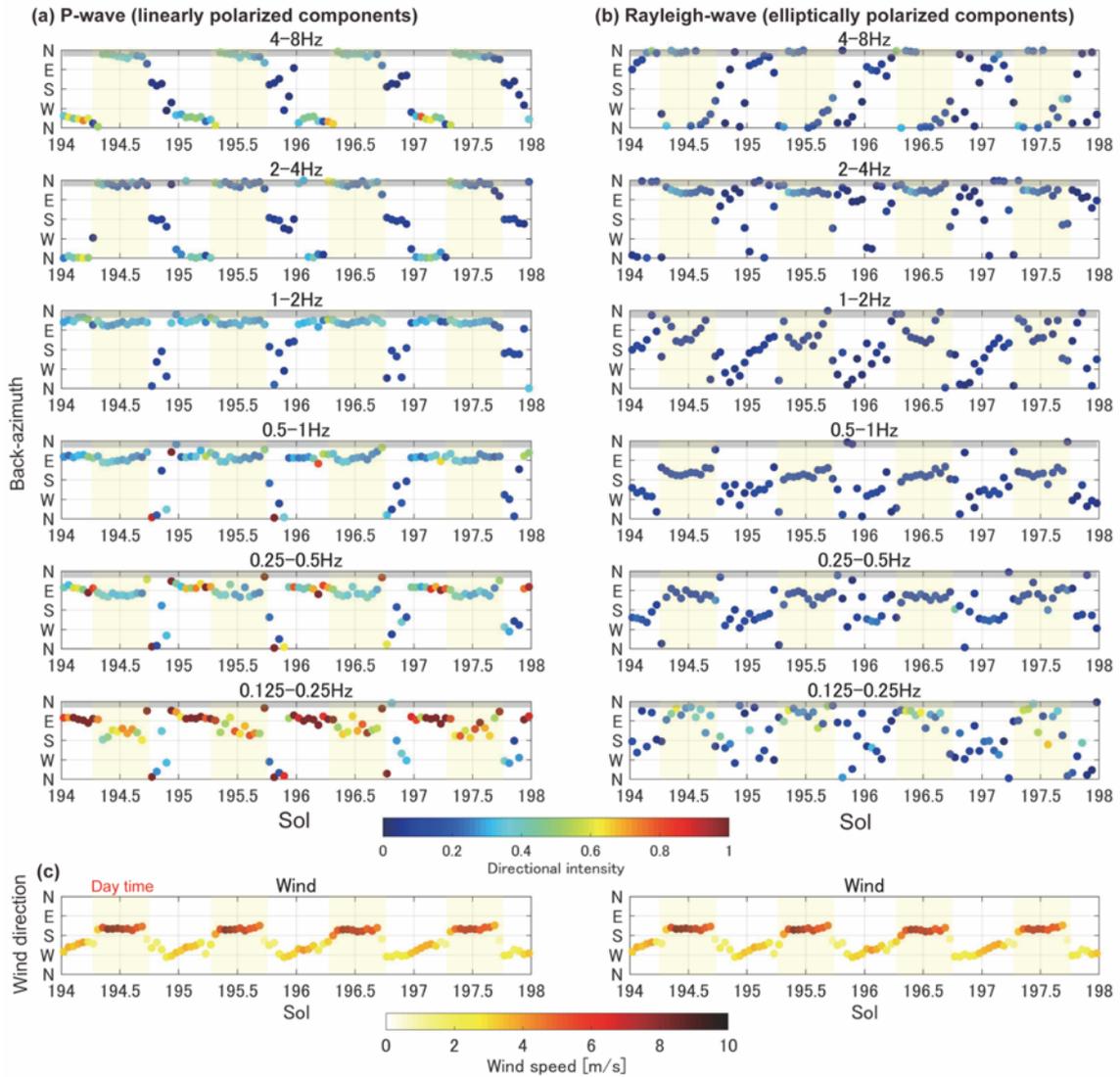
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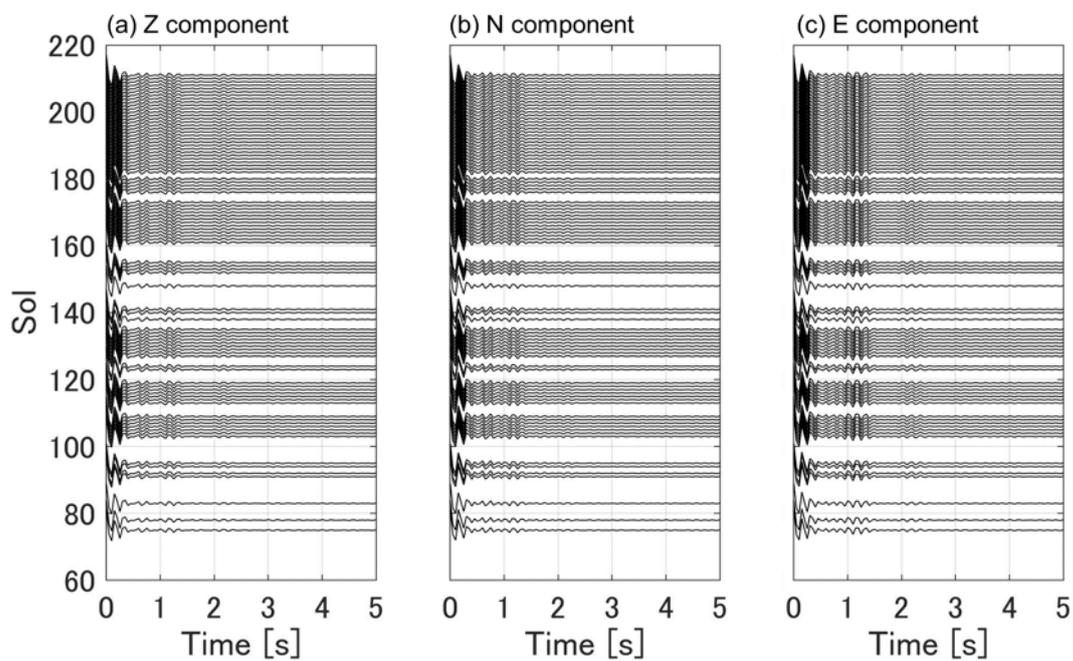
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Fig. 2. Temporal variation of dominant back-azimuths and directional intensity of (a) P-waves (linearly polarized components) and (b) Rayleigh waves (elliptically polarized components) in six single-octave frequency bands between Sols 75 and 211. (c) The wind speed and direction during the same period. The gray bar at the top of each back-azimuth plot indicates the direction of the tether connection between the seismometer and the lander (i.e., lander direction; Lognonné et al., 2020).



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465 Fig. 3. Temporal variations from Sols 194 to 197 in the dominant back-azimuths and
 466 directional intensity of (a) P-waves (linearly polarized components) and (b) Rayleigh
 467 waves (elliptically polarized components). (c) The wind speed and direction during the
 468 same period. Yellow shaded areas indicate daytime. The gray bar at the top of each back-
 469 azimuth plot indicates the direction of the tether connection between the seismometer and
 470 the lander (i.e., lander direction; Lognonné et al., 2020).
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473 Fig. 4. Temporal variation of autocorrelation functions of components from Sols 75 to
 474 211: (a) Vertical component; (b) NS component; (c) EW component. The vertical
 475 component could be similar to P-wave reflectivity whereas the NS and EW components
 476 could be S-wave reflectivity.

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